

# ***Percolation theory***

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# Renormalization

*general concepts*

$$M_\infty \sim l^{d_f}$$

rescaling transformation	$\frac{\xi}{b}$
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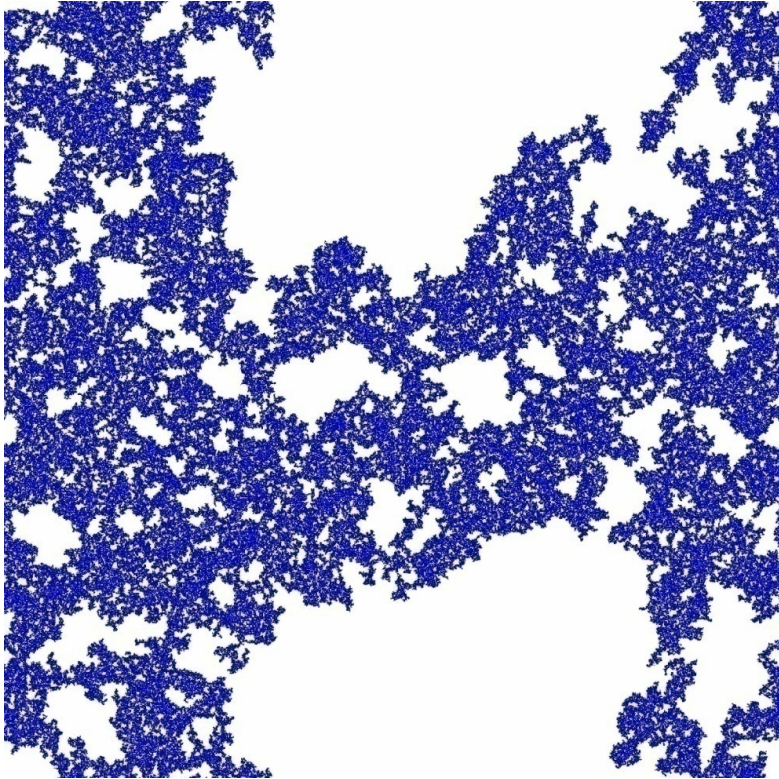
$\xi = \frac{\xi}{b}$
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*Fixed point*

$\xi = 0$

$\xi = \infty$

$p = \{0, p_c, 1\}$
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# Renormalization

## *rescaling*

*rescaling  
transformation*

$$\xi(p) = c|p - p_c|^{-\nu} \xrightarrow{\xi \mapsto \xi/b} \frac{\xi}{b} = c|T_b(p) - p_c|^{-\nu}$$

$$\frac{1}{b} = \left( \frac{|T_b(p) - p_c|}{|p - p_c|} \right)^{-\nu}$$

$$\nu = \frac{\ln b}{\ln \left( \frac{|T_b(p) - p_c|}{|p - p_c|} \right)}$$

$$T_b(p_c) = p_c$$

$$\nu = \frac{\ln b}{\ln \left( \left| \frac{dT_b}{dp} \right|_{p_c} \right)}$$

$T_b(p)$  requires *full information*  
about the system.

# Renormalization

real-space renormalization  
group transformation  $R_b$

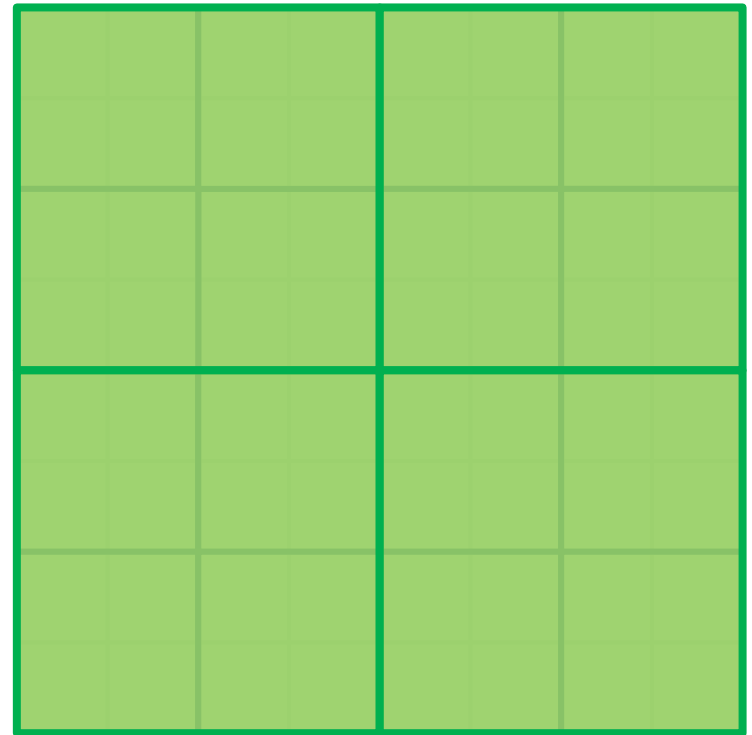
## *real-space renormalization*

$$R_b(p^*) = p^* \leftarrow \text{fixed point}$$

$$\nu \approx \frac{\ln b}{\ln \left( \frac{|R_b(p) - p_c|}{|p - p_c|} \right)}$$

$$\nu \approx \frac{\ln b}{\ln \left( \left| \frac{dR_b}{dp} \right|_{p_c} \right)}$$

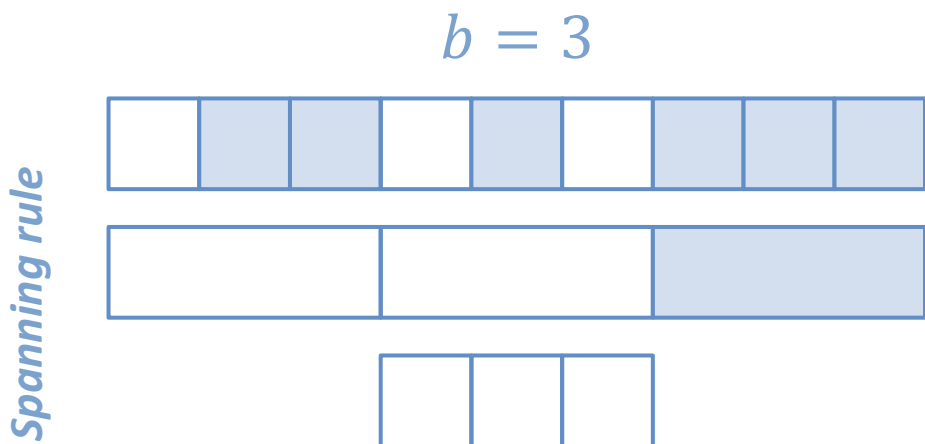
$b = 2$



# Renormalization

## 1D

1. *Divide* into blocks of size  $b$
2. *Replace* each block by a single one
3. *Rescale* all lengths by a factor of  $b$



$$R_b(p) = p^b$$

$\swarrow$   $\searrow$

$$p^* = 0 \qquad p^* = 1$$

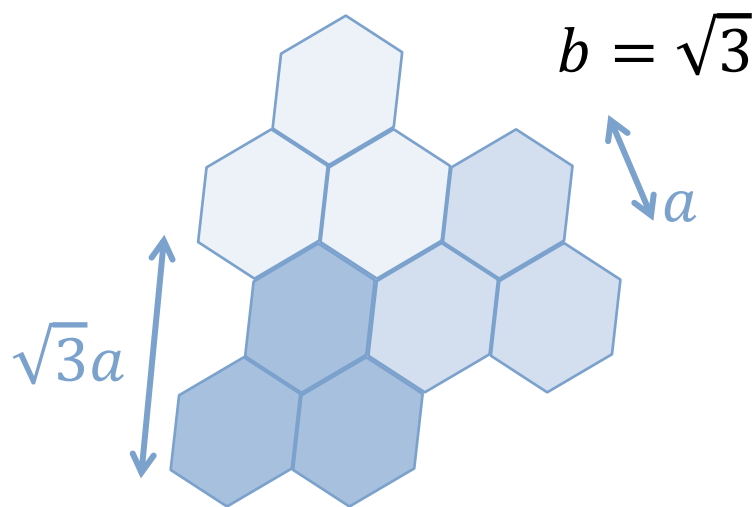
$$\xi(R_b(p)) = -\frac{1}{\ln R_b(p)} = -\frac{1}{\ln p^b} = \frac{\xi(p)}{b}$$

$$\left. \frac{dR_b}{dp} \right|_{p^*=1} = bp^{b-1} \Big|_{p^*=1} = b$$

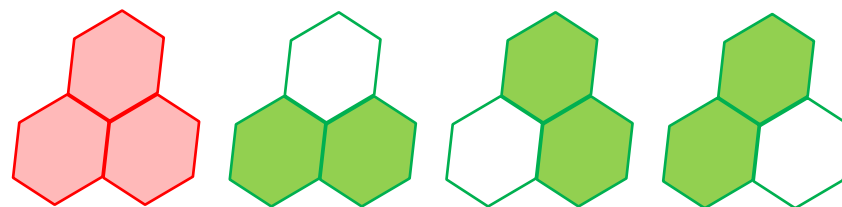
$$v = \frac{\ln b}{\ln \left( \left| \frac{dR_b}{dp} \right|_{p_c} \right)} = 1$$

# Renormalization

## 2D triangular lattice



Spanning rule



$$R_b(p) = p^3 + 3p^2(1 - p)$$

$$p^* = 0$$

$$p^* = 1/2$$

$$p^* = 1$$

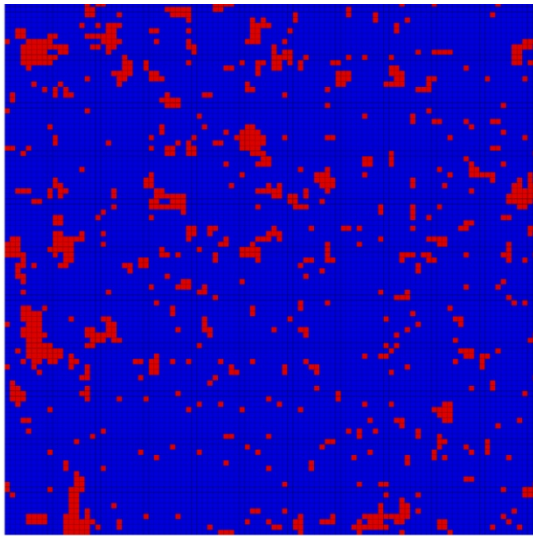
$$\nu = \frac{\ln b}{\ln \left( \left| \frac{dR_b}{dp} \right|_{p^*} \right)} = \frac{\ln \sqrt{3}}{\ln(3/2)} \approx 1.355$$

$$\nu = \frac{4}{3}$$

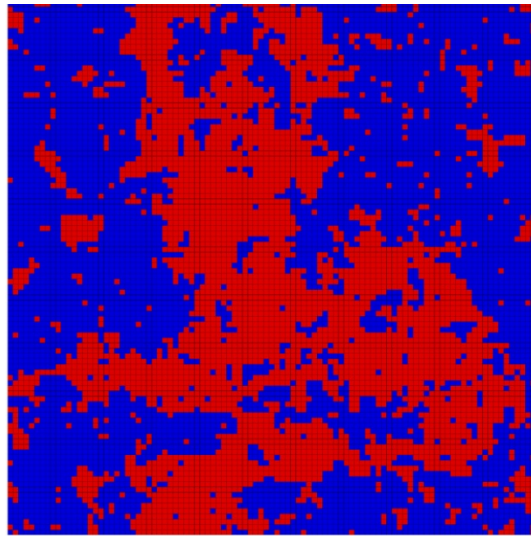
# Percolation and magnetic models

## *Ising model (2-state Potts model)*

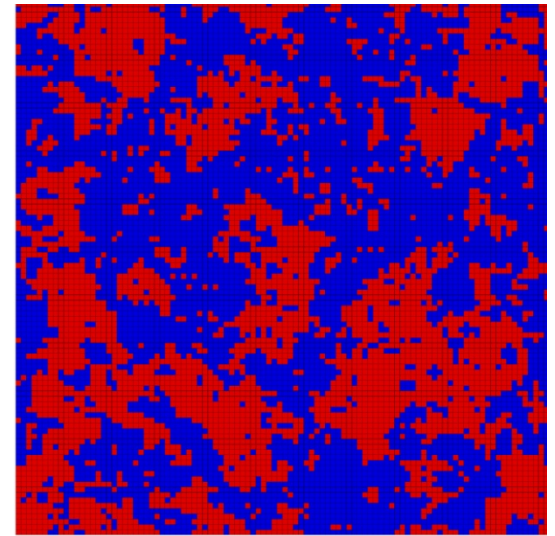
$$\mathcal{H} = J \sum_b \varepsilon_b \quad \text{with } \varepsilon_b = \begin{cases} 0, & \text{if endpoints are in the same state} \\ 1, & \text{if endpoints are in different states} \end{cases}$$



$$T < T_c$$



$$T = T_c$$

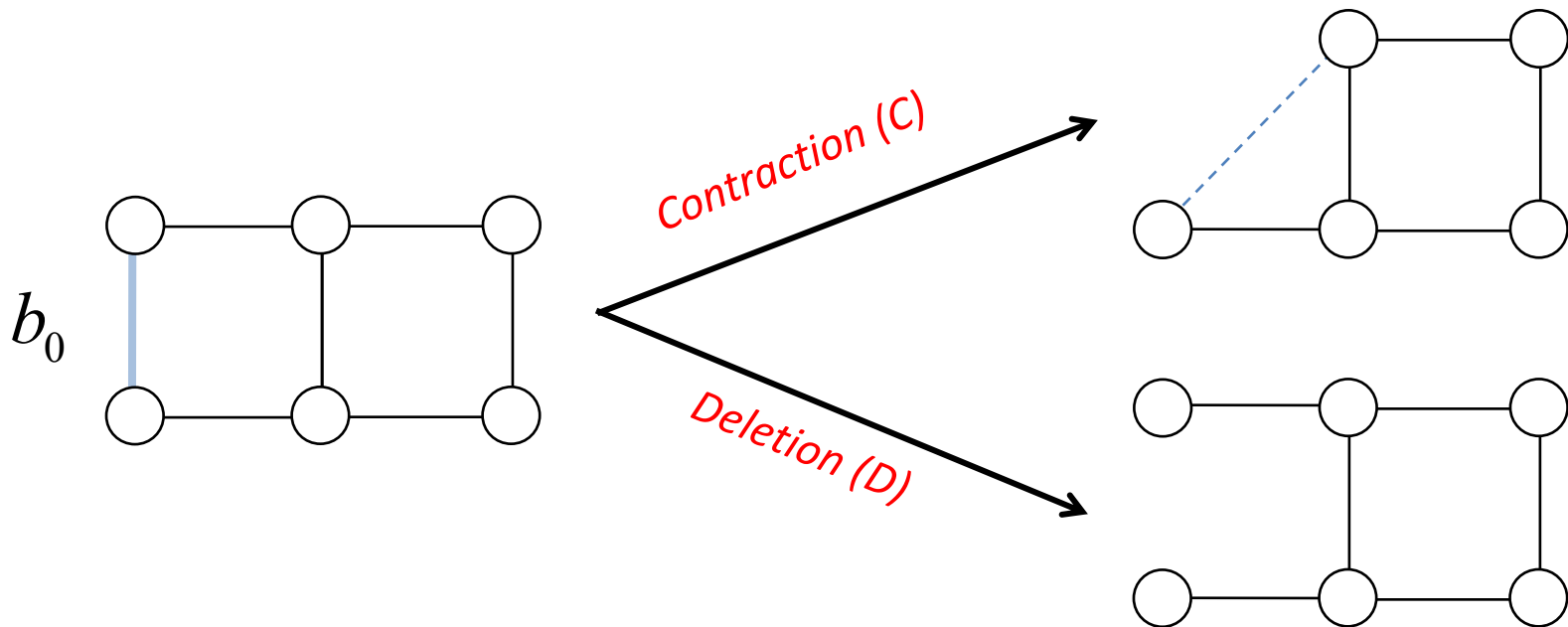


$$T > T_c$$

# Percolation and magnetic models

## *theorem of Kasteleyn-Fortuin*

$$\mathcal{H} = J \sum_b \varepsilon_b \quad \text{with } \varepsilon_b = \begin{cases} 0, & \text{if endpoints are in the same state} \\ 1, & \text{if endpoints are in different states} \end{cases}$$





# Percolation and magnetic models

## *theorem of Kasteleyn-Fortuin*

$$\mathcal{H} = J \sum_b \varepsilon_b$$

$$Z = \sum_X e^{-\beta \mathcal{H}(X)} = \sum_X e^{-\beta J \sum_b \varepsilon_b} = \sum_X \prod_b e^{-\beta J \varepsilon_b}$$

Partition  
Function

Consider bond  $b_0$  with endpoints  $i$  and  $j$ :

$$\begin{aligned} Z &= \sum_X e^{-\beta J \varepsilon_{b_0}} \prod_{b \neq b_0} e^{-\beta J \varepsilon_b} = \sum_{\substack{X: \\ \sigma_i = \sigma_j}} \prod_{b \neq b_0} e^{-\beta J \varepsilon_b} + e^{-\beta J} \sum_{\substack{X: \\ \sigma_i \neq \sigma_j}} \prod_{b \neq b_0} e^{-\beta J \varepsilon_b} \\ &= Z_C + e^{-\beta J} (Z_D - Z_C) = (1 - e^{-\beta J}) Z_C + e^{-\beta J} Z_D = p Z_C + (1 - p) Z_D \end{aligned}$$

$Z_C$  and  $Z_D$  are the partition functions of the graphs contracted and deleted at  $b_0$

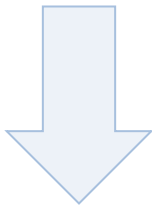
$$p \equiv 1 - e^{-\beta J}$$

# Percolation and magnetic models

## *theorem of Kasteleyn-Fortuin*

For bond  $b_1$ :  $Z = pZ_{C_{b_1}} + (1-p)Z_{D_{b_1}}$ , and doing the same for  $b_2$ :

$$Z = p^2 Z_{C_{b_1}C_{b_2}} + p(1-p)Z_{C_{b_1}D_{b_2}} + (1-p)pZ_{D_{b_1}C_{b_2}} + (1-p)^2 Z_{D_{b_1}D_{b_2}}$$



Doing the same for all edges, one obtains a set of separated points/clusters (contracted graphs). Each can be in  $q$  different states.

$$Z = \sum_{\text{Configurations of bond percolation}} q^{\text{\# of clusters}} p^c (1-p)^d$$

$c$  and  $d$  are the number of contracted and deleted bonds.

$$p \equiv 1 - e^{-\beta J}$$