Kelvin's circulation theorem

- An ideal fluid that is vorticity free at a given instant is vorticity free at all times.
- Demonstration: see Faber 120-122
- In three dimensions the conservation of vorticity (which corresponds to the conservation of angular momentum in mechanics) takes a somewhat subtle form.
- The circulation of a velocity field is defined to be

$$K\{t\} = \oint u\{x,t\} \cdot \mathrm{d}l,$$



where the line is a closed loop which moves with the fluid.

Circulation and vorticity

• By Stokes' theorem

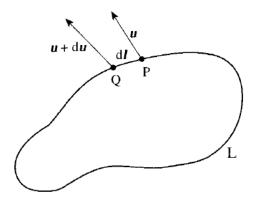
$$K = \oint_{C(t)} \mathbf{u} \cdot d\mathbf{l} = \int_{S(t)} (\underbrace{\mathbf{\nabla} \times \mathbf{u}}_{\Box}) \cdot \mathbf{n} dS = \int_{S(t)} \Omega \cdot \mathbf{n} dS,$$

where S(t) is a surface whose edges connect with C(t).

K is zero for all loops if Ω is zero in the domain!

Kelvin's theorem asserts that

$$\frac{DK}{Dt} = 0.$$



Demonstration

The loop moves with the flow and thus

$$\frac{DK}{Dt} = \oint \frac{D\vec{n}}{Dt} \cdot \vec{d\vec{r}} + \vec{n} \cdot \frac{D(\vec{d\vec{r}})}{Dt}$$

The second term is the relative velocity of two nearby points on the loop and can be written as $(\partial u/\partial l)dl$,

×

72

$$\frac{D(d\vec{r})}{Dt} = d\vec{r} \cdot \nabla \vec{u}$$

Hence

$$\oint \vec{u} \cdot D(d\vec{\ell}) = \oint \vec{u} \cdot \left[(d\vec{\ell} \cdot \nabla) \cdot \vec{u} \right] = \frac{1}{2} \oint \nabla(\vec{u}^2) \cdot d\vec{\ell}$$

$$\int \int (\nabla \times \nabla \cdot \vec{u}^2) \cdot d\vec{S} = 0$$

$$= \int_{2}^{1} \int_{2}^{1} (\nabla \times \nabla \cdot \vec{u}^2) \cdot d\vec{S} = 0$$
18

150

If the fluid is incompressible, using Euler:

$$\frac{Du}{Dt} = -\nabla \left(\frac{p}{\rho} + gz\right),$$

$$\int_{C} \frac{Du}{Dt} \cdot d\tau = -\oint_{C} \nabla \left(\frac{p}{\rho} + gz\right) \cdot d\tau$$

$$\int_{C} \frac{Stokes}{p} - \int_{C} \left[\nabla \times \nabla \left(\frac{p}{\rho} + gz\right)\right] \cdot ds = 0$$

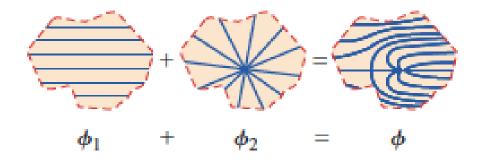
Therefore:

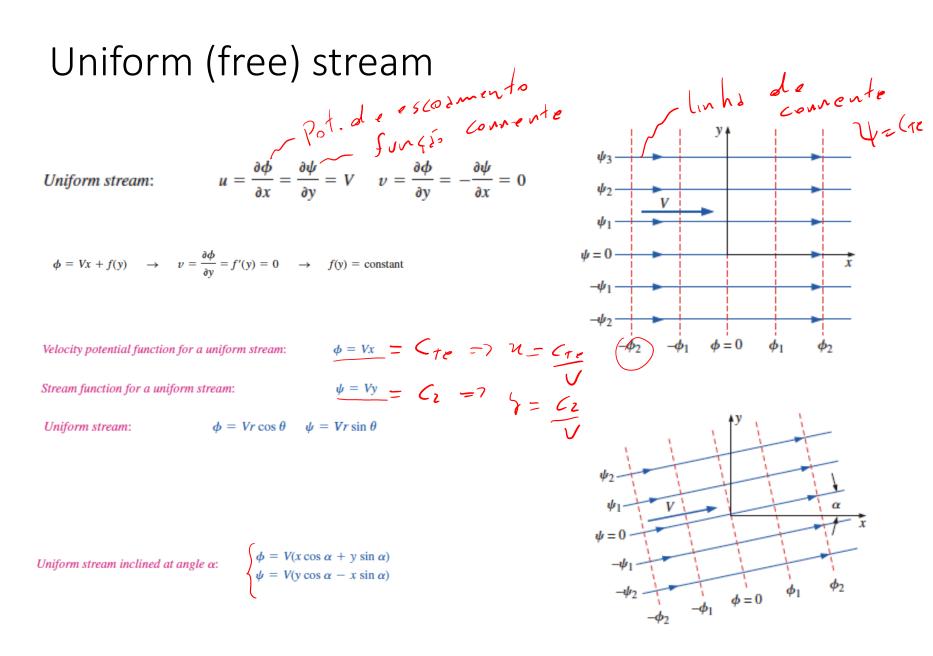
$$\frac{DK}{Dt} = 0 \qquad - \int (v do i de) (1 - v do i de) (1 - v do i de) (1 - v do de) (1 - v$$

Superposition

 $\vec{V} = \nabla \phi$

- Since the Laplace equation is a linear homogeneous differential equation, the linear combination of two or more solutions of the equation must also be a solution.
- For example, if ϕ_1 and ϕ_2 are each solutions of the Laplace equation, then $A \phi_1 + B \phi_2$ are also solutions, where A and B are arbitrary constants. $\nabla^2 \phi_r = \nabla^2 (A \phi_A + B \phi_z) = A \nabla^2 \phi_z + B \nabla^2 \phi_z$ • By extension, you may combine several solutions of the Laplace
- equation, and the combination is guaranteed to also be a solution.





Line source or sink

Let the volume flow rate per unit depth, be the line source strength, m

$$\frac{\dot{V}}{L} = 2\pi r u_r \qquad u_r = \frac{\dot{V}/L}{2\pi r}$$

The components of the velocity are

Line source:

$$u_r = \frac{\partial \phi}{\partial r} = \frac{1}{r} \frac{\partial \psi}{\partial \theta} = \frac{\dot{V}/L}{2\pi r} \qquad u_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = -\frac{\partial \psi}{\partial r} = 0$$

$$\frac{\partial \psi}{\partial r} = -u_{\theta} = 0 \quad \rightarrow \quad \psi = f(\theta) \quad \rightarrow \quad \frac{\partial \psi}{\partial \theta} = f'(\theta) = ru_{r} = \frac{\dot{V}/L}{2\pi}$$

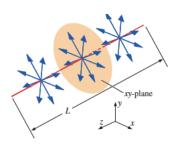
With solution

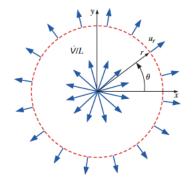
$$f(\theta) = \frac{\dot{V}/L}{2\pi}\theta + \text{constant}$$

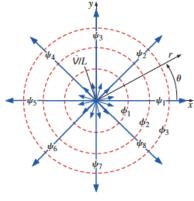
Ф

Line source at the origin:

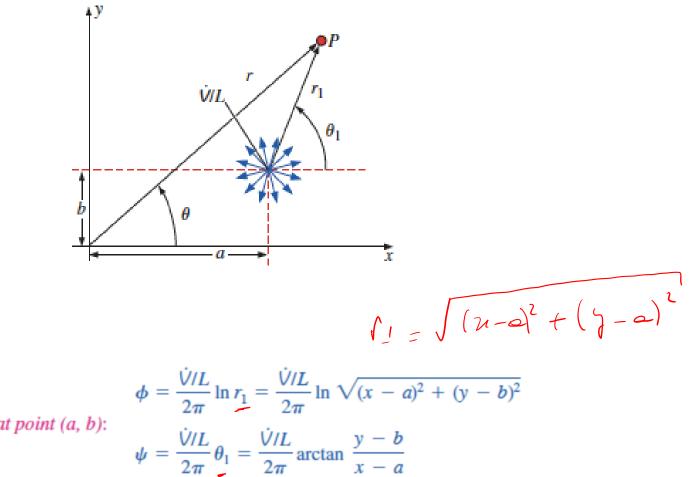
$$= \frac{\dot{V}/L}{2\pi} \ln r \quad \text{and} \quad \psi = \frac{\dot{V}/L}{2\pi} \theta$$







Line source or sink at an arbitrary point



Line source at point (a, b):

Superposition of a source and sink of equal strength

Line source at (-a, 0): $\psi_1 = \frac{\dot{V}/L}{2\pi} \theta_1$ where $\theta_1 = \arctan \frac{y}{x+a}$ Similarly for the sink, Line sink at (a, 0): $\psi_2 = \frac{-\dot{V}/L}{2\pi} \theta_2$ where $\theta_2 = \arctan \frac{y}{x-a}$ Composite stream function: $\psi = \psi_1 + \psi_2 = \frac{\dot{V}/L}{2\pi} (\theta_1 - \theta_2)$

Final result, Cartesian coordinates:
$$\psi = \frac{-\dot{V}/L}{2\pi} \arctan \frac{2ay}{x^2 + y^2 - a^2}$$

Final result, cylindrical coordinates:
$$\psi = \frac{-\dot{V}/L}{2\pi} \arctan \frac{2ar \sin \theta}{r^2 - a^2}$$

Using

$$rctan(u)\pm rctan(v)=rctaniggl(rac{u\pm v}{1\mp uv}iggr)\pmod{\pi},\quad uv
eq 1\,.$$

24

Line vortex

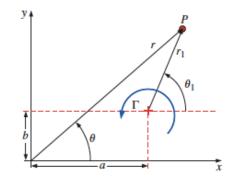
The radial component of the velocity is zero and

Line vortex:
$$u_r = \frac{\partial \phi}{\partial r} = \frac{1}{r} \frac{\partial \psi}{\partial \theta} = 0$$
 $u_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = -\frac{\partial \psi}{\partial r} = \frac{\Gamma}{2\pi r}$

where $\Gamma = 2\pi r u_{\theta}$, is the circulation, around a loop of radius r.



x



θ

Line vortex at the origin: $\phi = \frac{\Gamma}{2\pi} \theta \quad \psi = -\frac{\Gamma}{2\pi} \ln r$ $\phi = \frac{\Gamma}{2\pi} \theta_1 = \frac{\Gamma}{2\pi} \arctan \frac{y - b}{x - a}$ Line vortex at point (a, b): $\psi = -\frac{\Gamma}{2\pi} \ln r_1 = -\frac{\Gamma}{2\pi} \ln \sqrt{(x - a)^2 + (y - b)^2}$

Superposition of a line sink and a line vortex at the origin

The stream function is

Superposition:

with streamlines

Streamlines:

 $\psi = \frac{\dot{V}/L}{2\pi}\theta - \frac{\Gamma}{2\pi}\ln r$ y, m 0.3 $r = \exp\left(\frac{(\dot{V}/L)\theta - 2\pi\psi}{\Gamma}\right)$

Velocity components:

 $u_r = \frac{1}{r} \frac{\partial \psi}{\partial \theta} = \frac{\dot{V}/L}{2\pi r}$ $u_{\theta} = -\frac{\partial \psi}{\partial r} = \frac{\Gamma}{2\pi r}$

Note that velocity diverges at the origin, which is a singularity (unphysical).

0.8

0.9

2

0.6

0

x, m

0.5