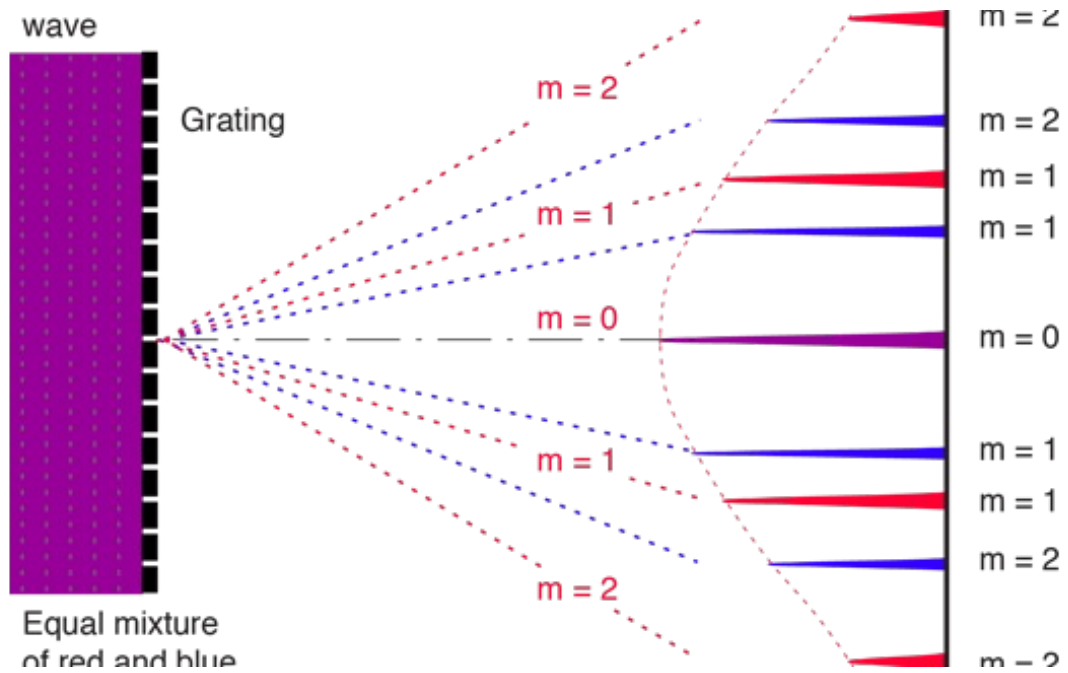
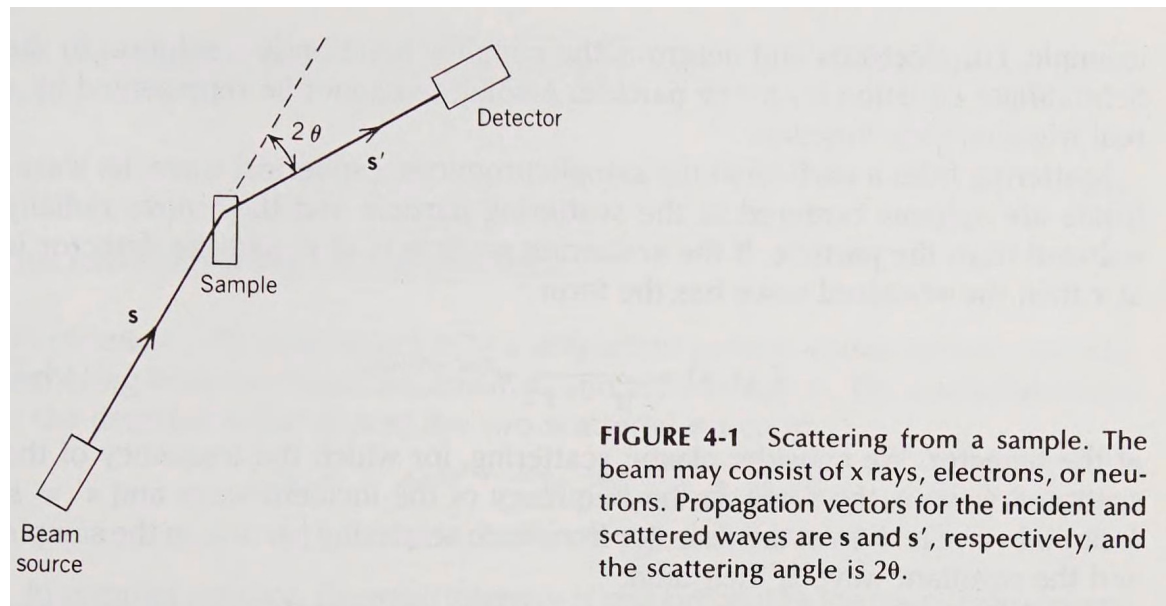
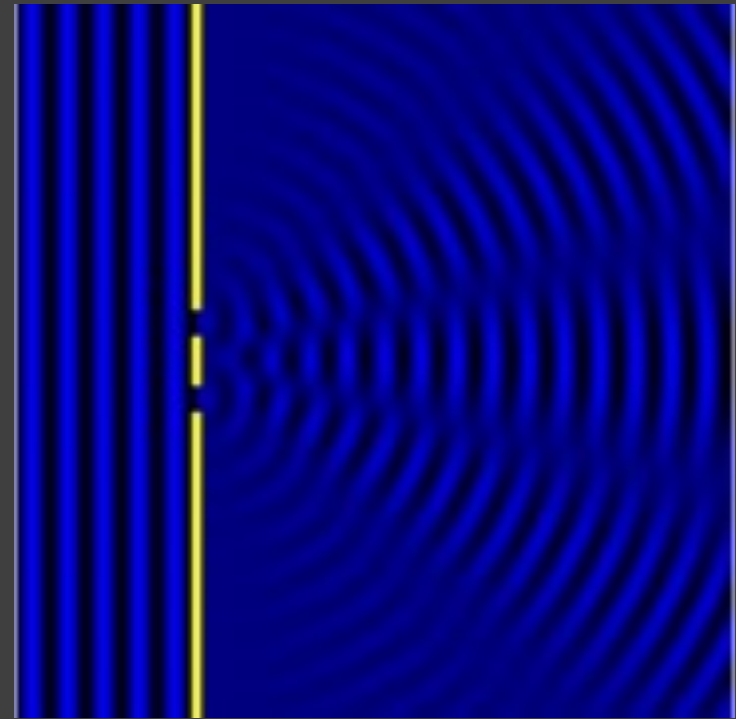
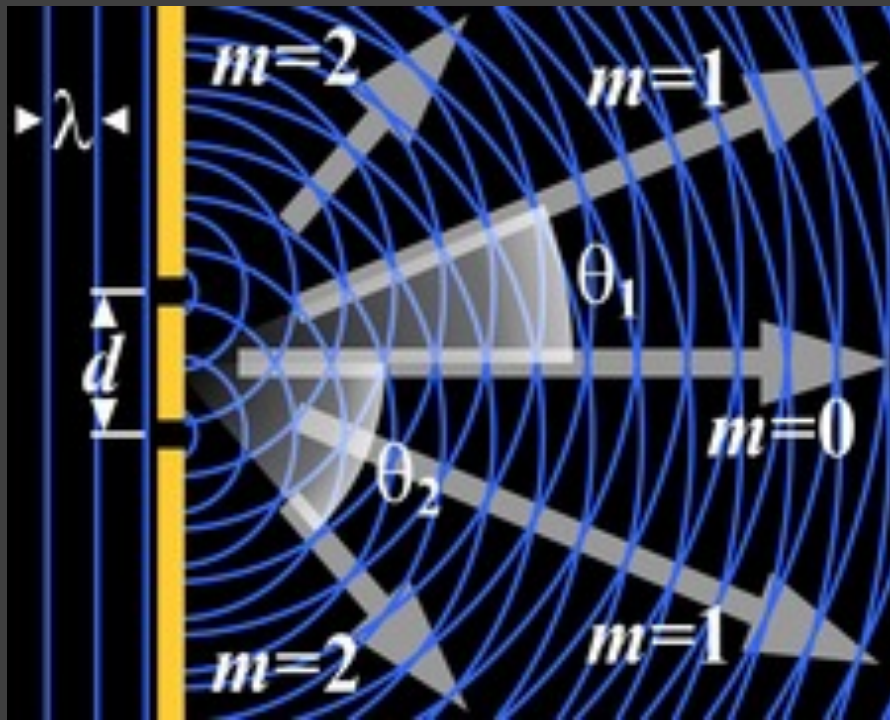


4. Diffraction



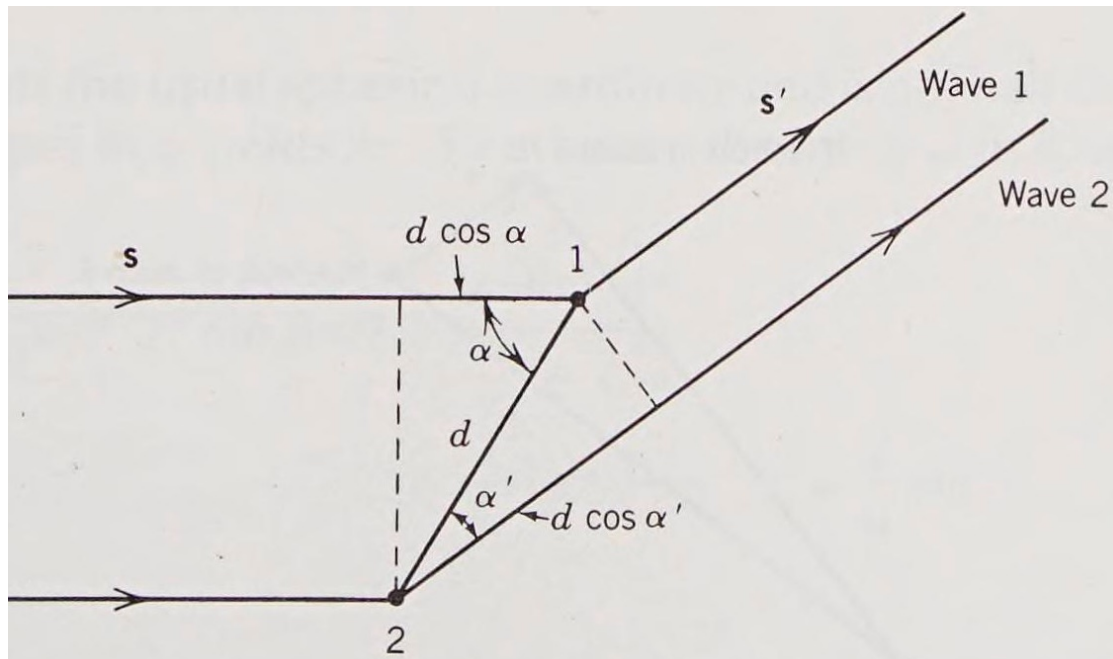
Scattering from a sample

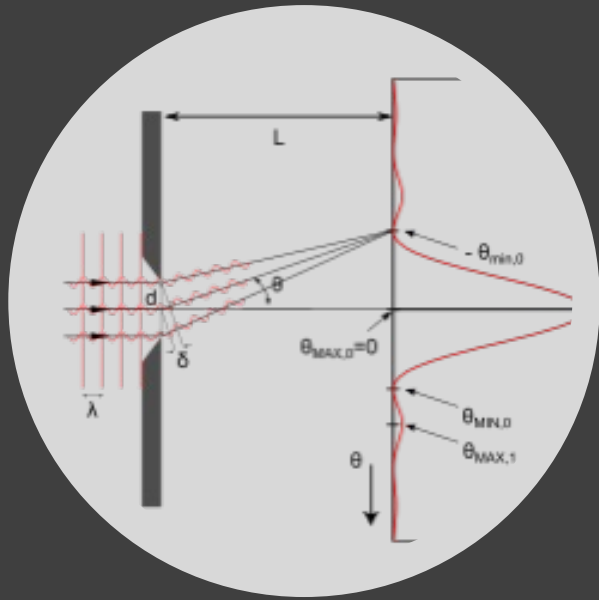




Two slits

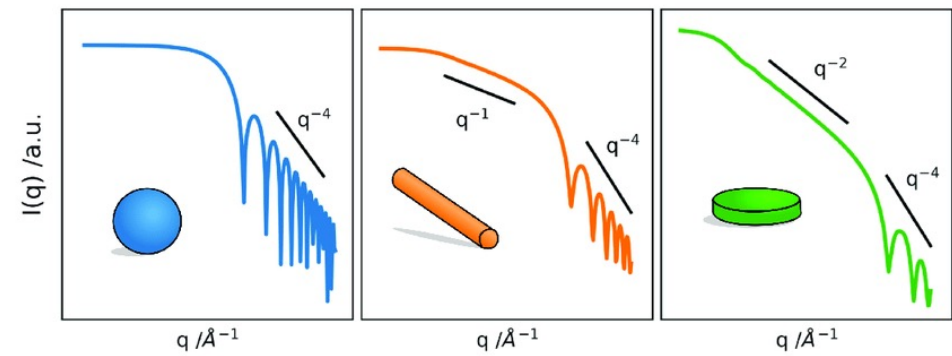
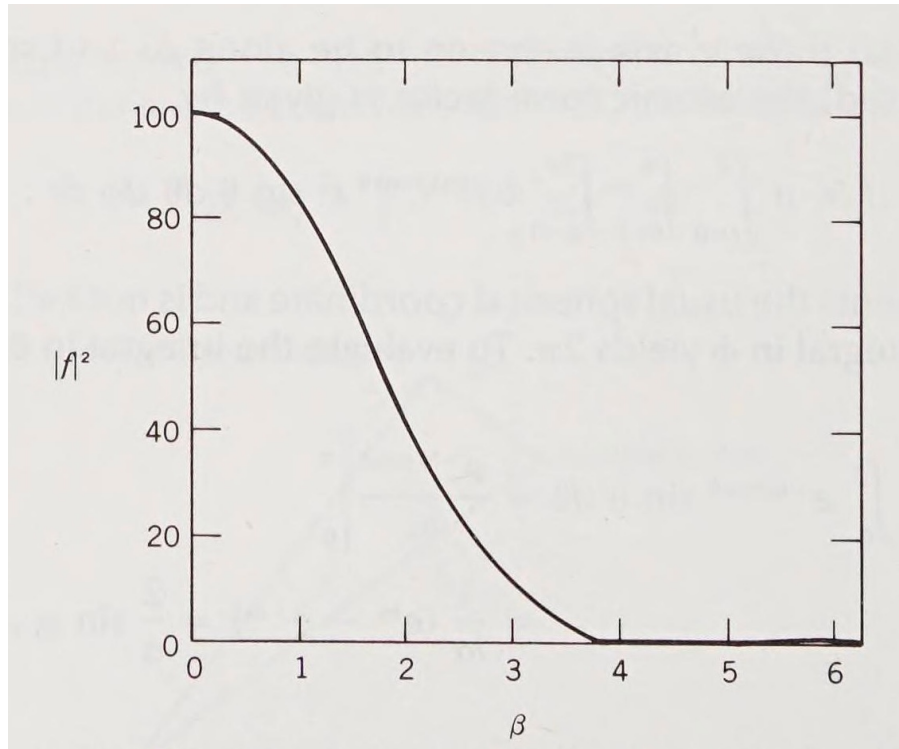
Scattering from two particles

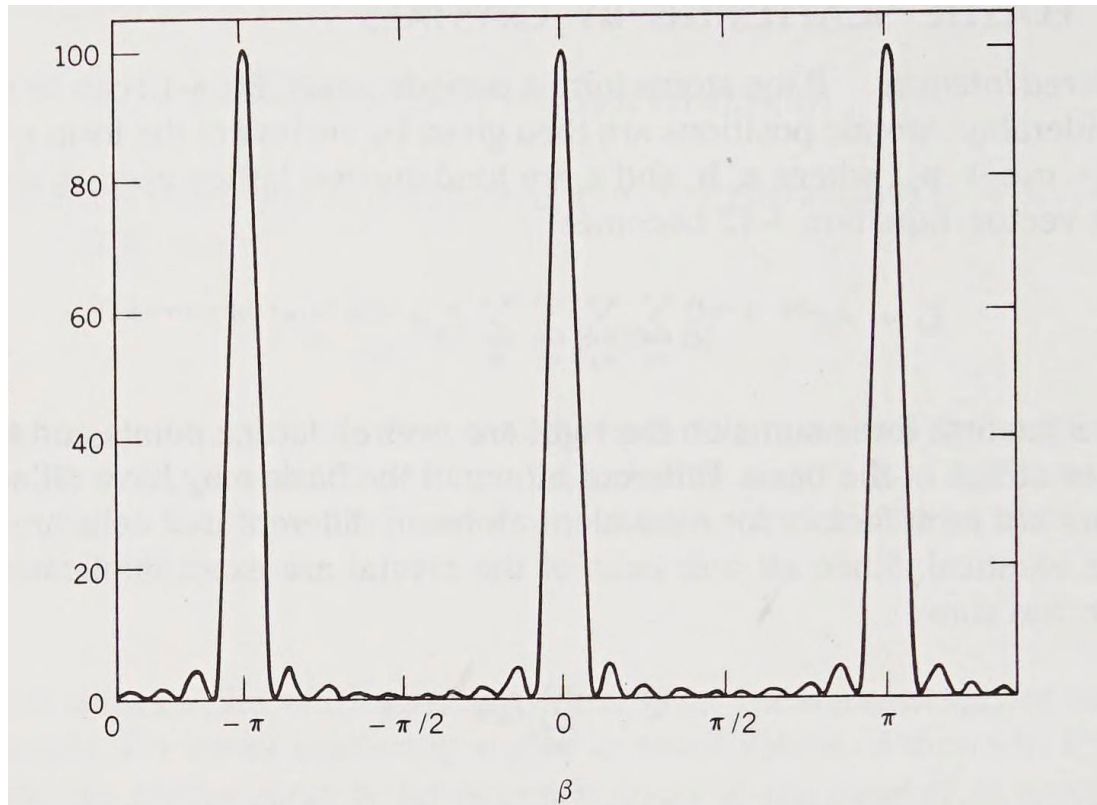




Diffraction from an aperture

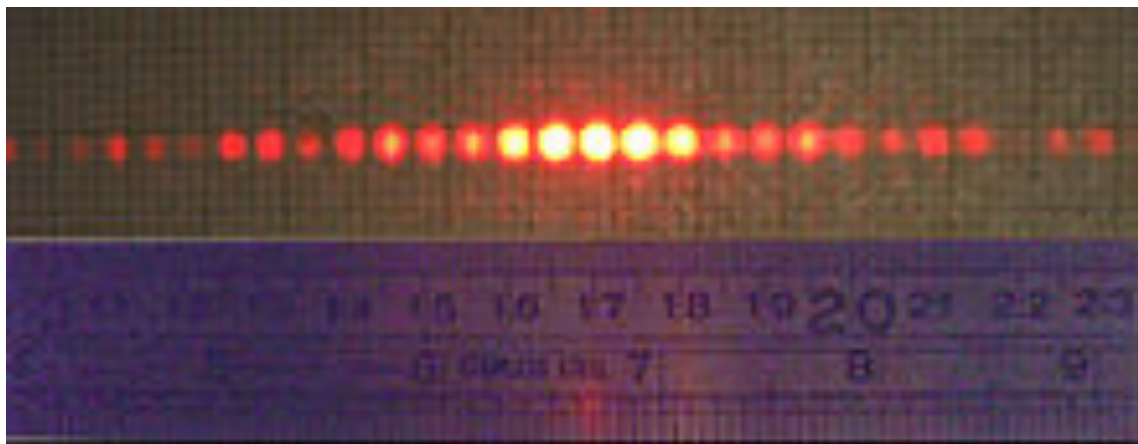
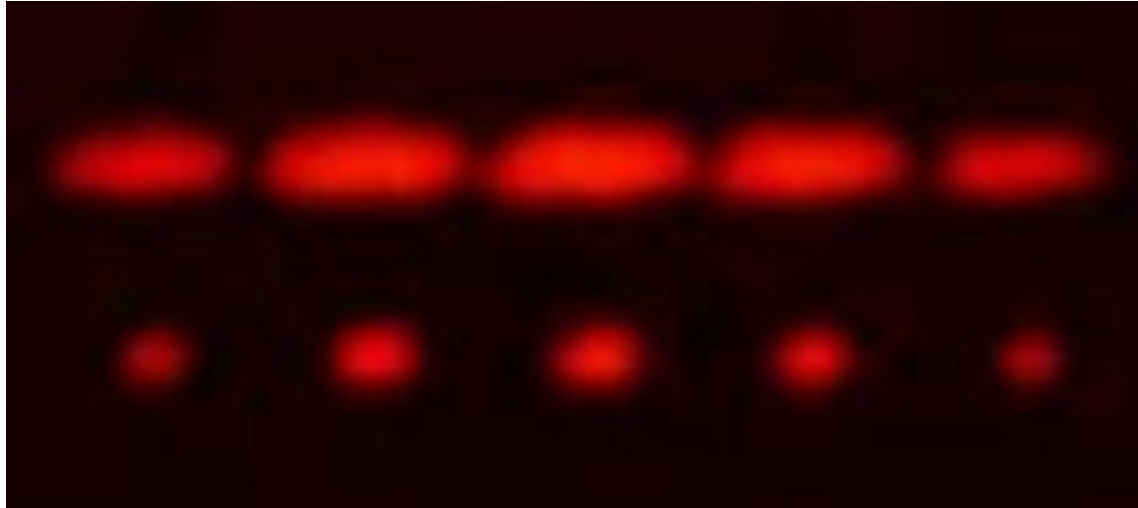
Form factor





$$\beta = \frac{1}{2} \Delta \mathbf{s} \cdot \mathbf{a} \text{ and } g(\beta) = \frac{\sin^2(N\beta)}{\sin^2(\beta)}$$

Scattering
function for
 $N=10$



Diffraction patterns
for two, five and
one hundred and
fifty slits

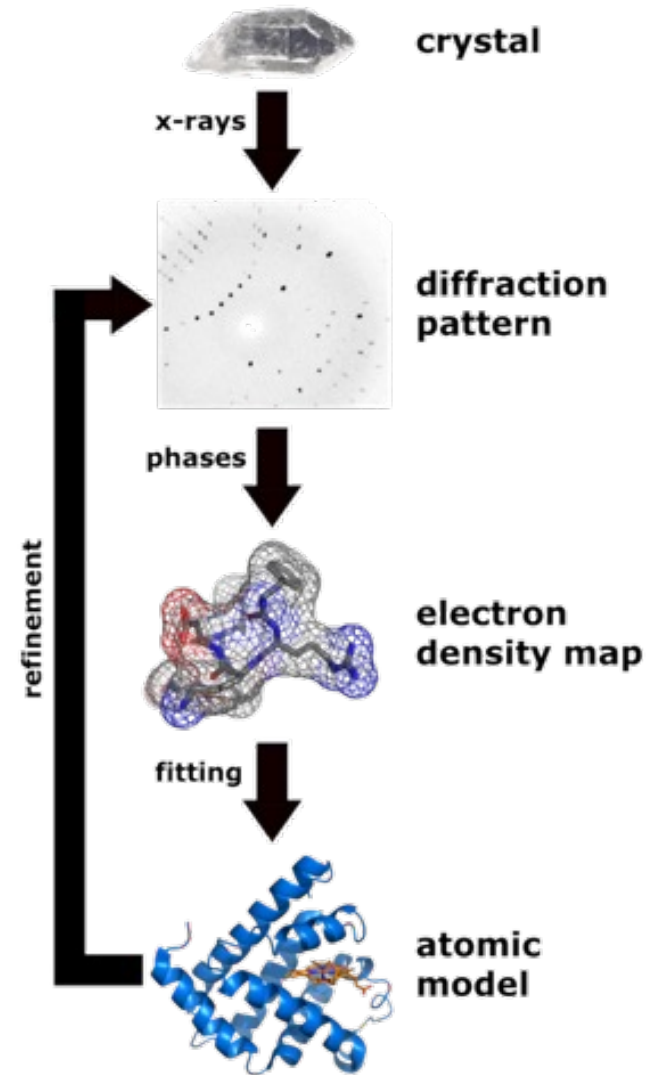
Laue scattering condition

$$\Delta \mathbf{s} = h\mathbf{A} + k\mathbf{B} + l\mathbf{C}.$$

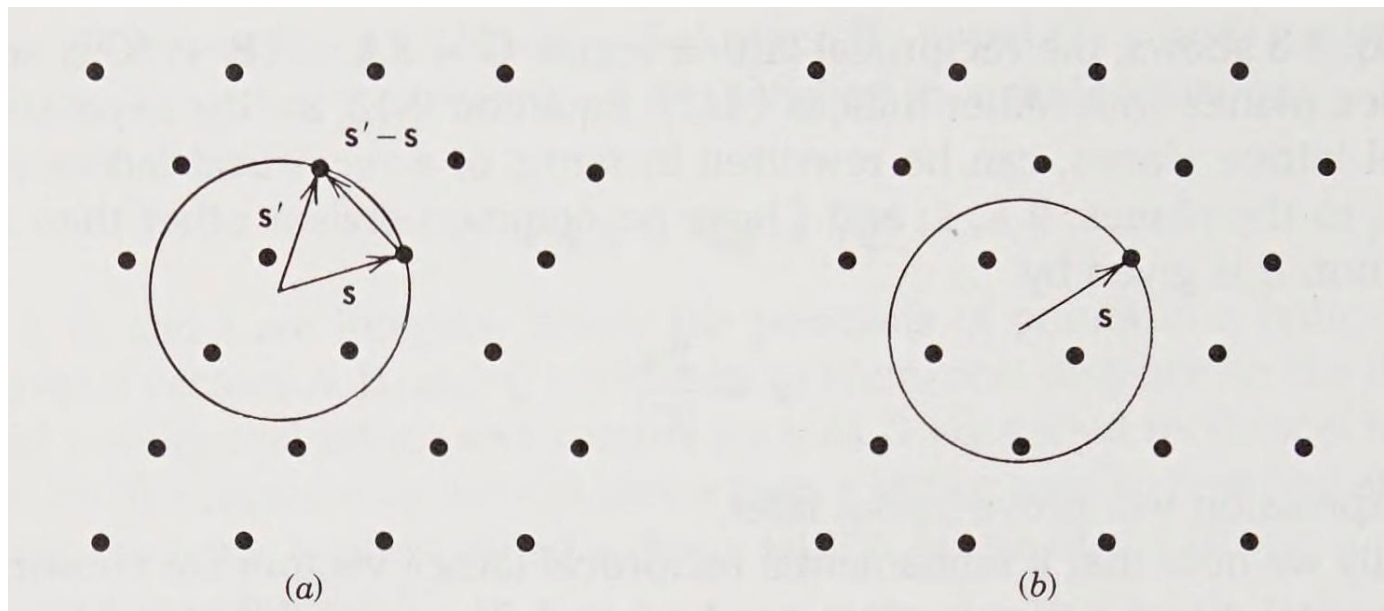


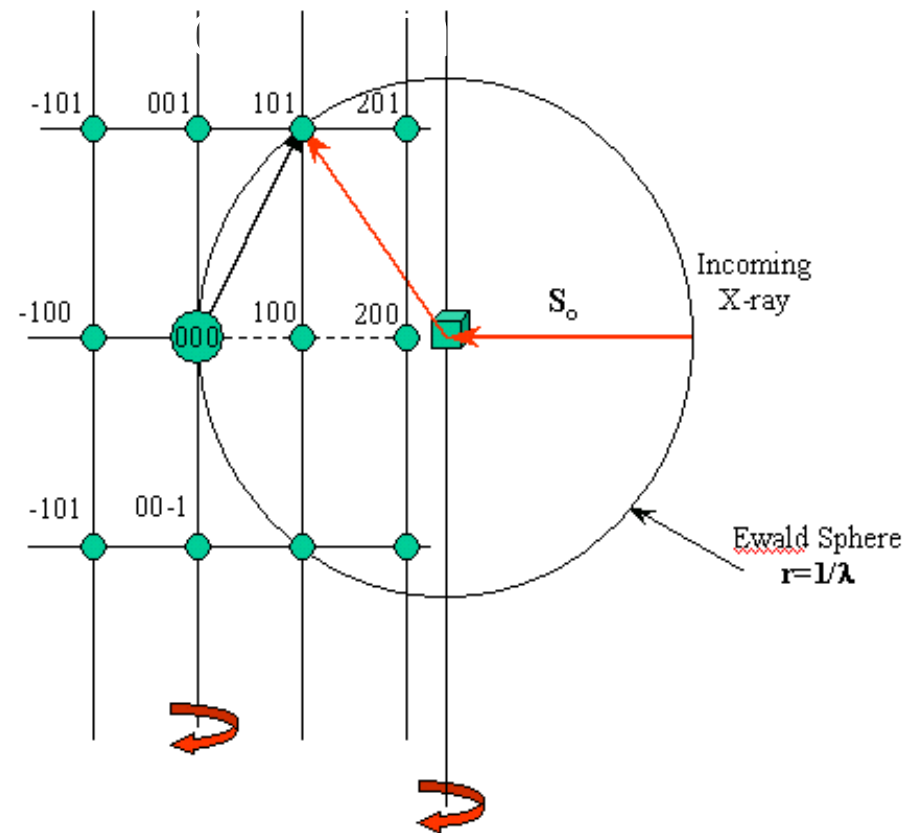
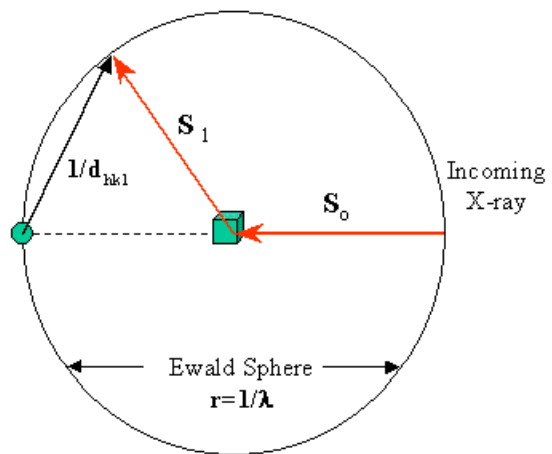
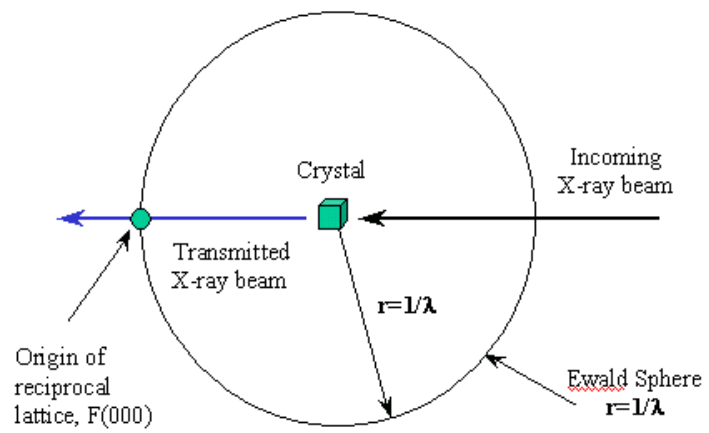
NOBEL PRIZE IN
PHYSICS, 1914

Max von Laue

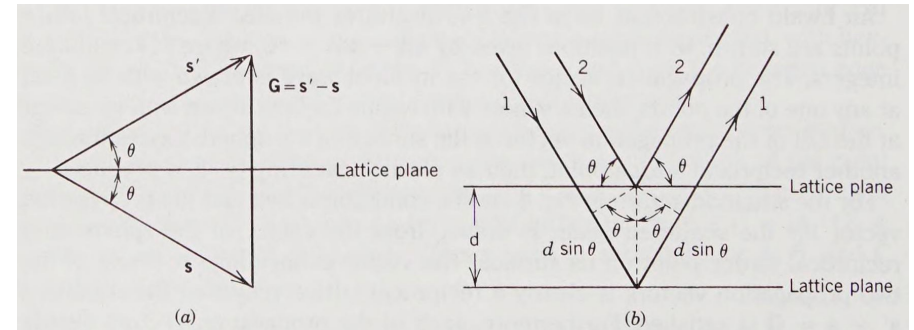
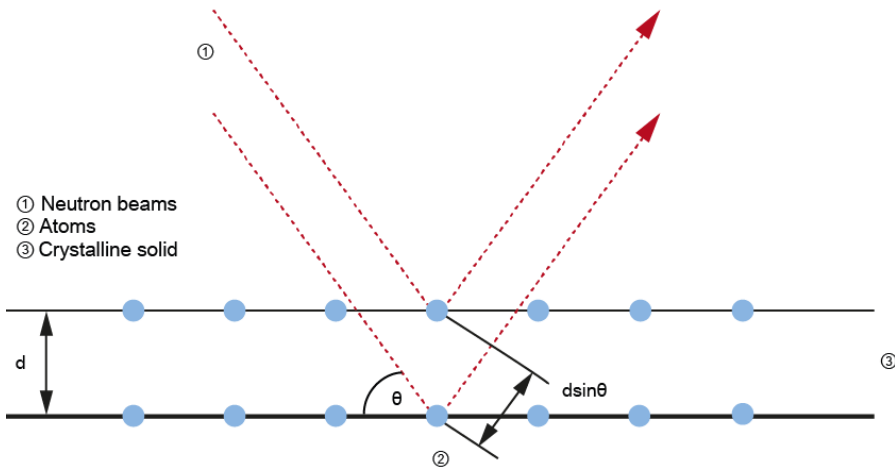


Ewald construction



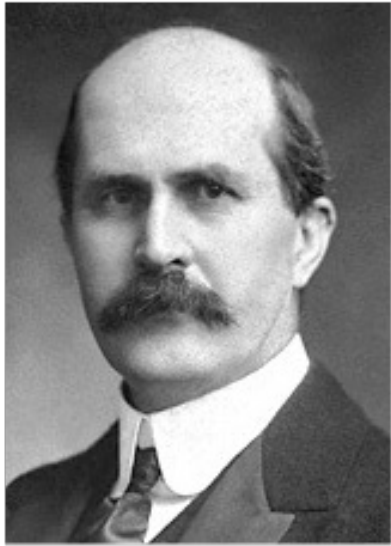


Bragg diffraction

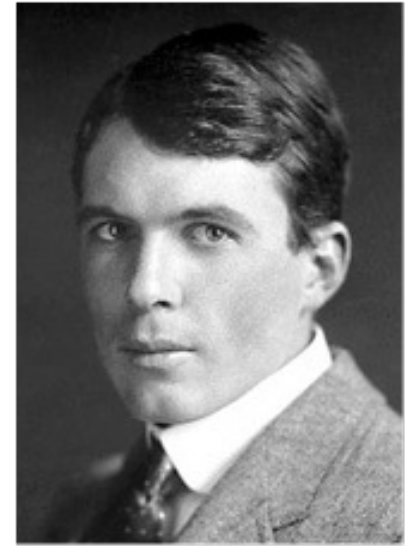
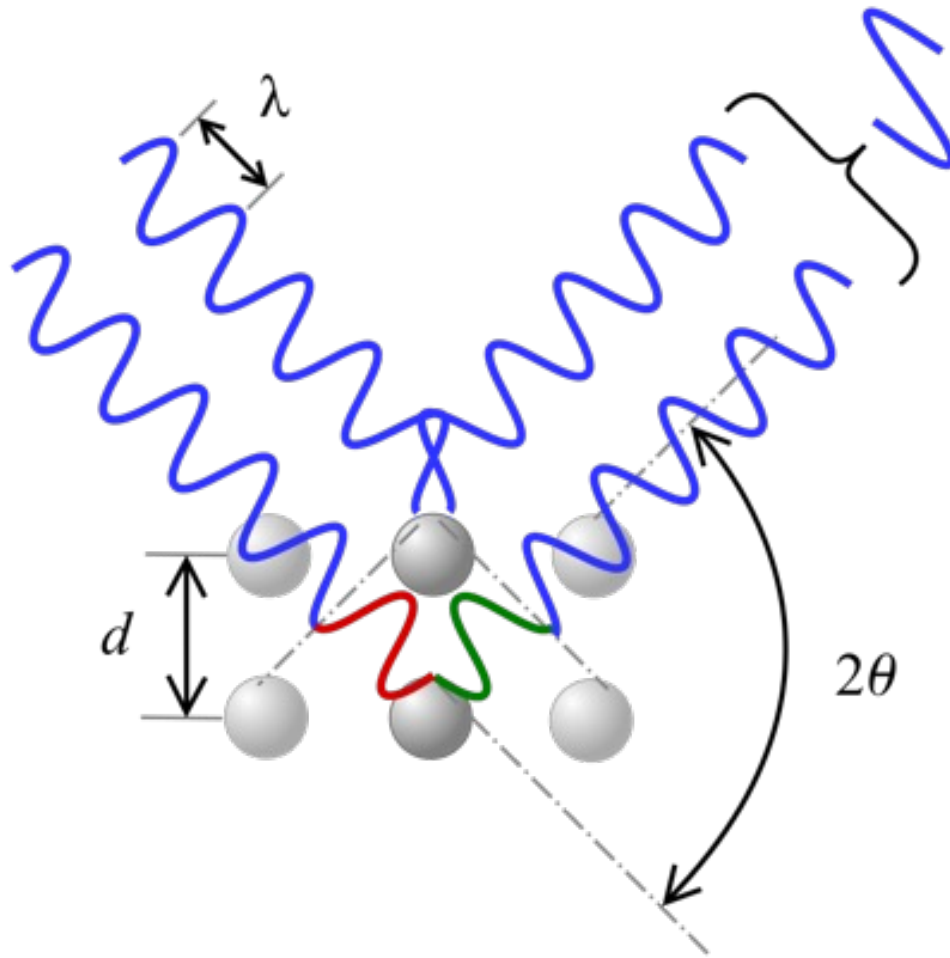


$$2d \sin \theta = n\lambda$$

Bragg scattering



WH BRAGG



WL BRAGG

NOBEL PRIZE IN PHYSICS, 1915

Structure factor

$$F = \sum_m f_m e^{-2\pi i(u_m k + v_m k + w_m l)}.$$

Structure factor

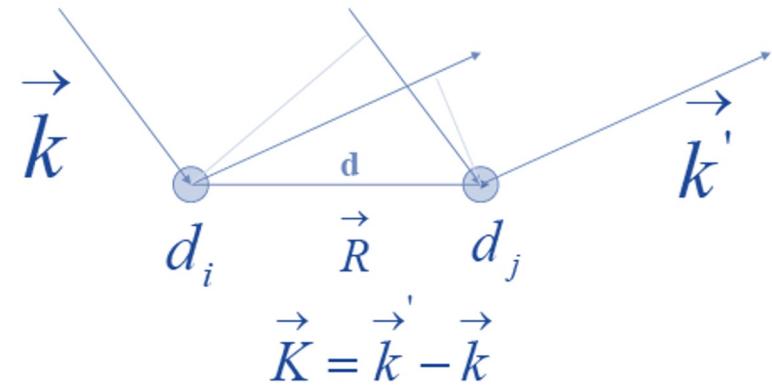
The amplitude of the rays scattered at d_1, d_2, d_3, \dots are in the ratios :

$$e^{-iK \cdot d_j}$$

The net ray scattered by the entire cell:

$$S_k = \sum_{j=1}^n e^{-iK \cdot d_j}$$

$$I_{(hkl)} \propto |S_k|^2$$

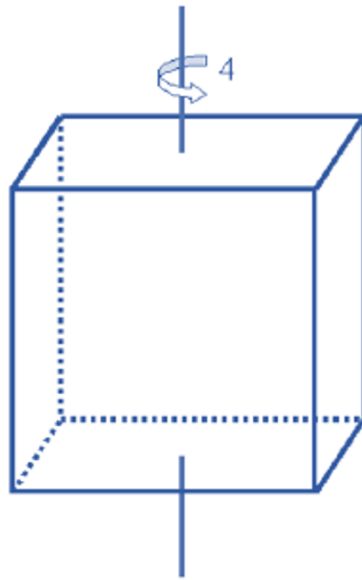


Phase difference: $K \cdot (d_i - d_j)$

The amplitude of the two rays differ: $e^{iK \cdot (d_i - d_j)}$

Structure factor SC

For simple cubic: (0,0,0)



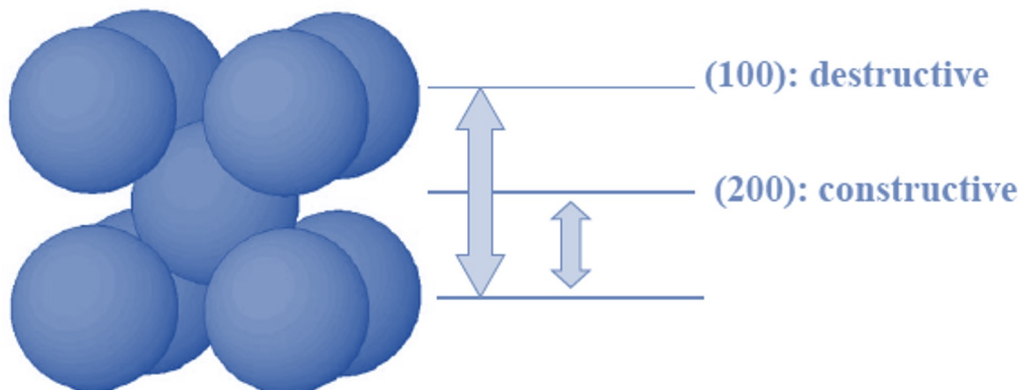
$$S_k = e^{-iK \cdot 0} = 1$$

Structure factor BCC

For BCC: $(0,0,0), (1/2, 1/2, 1/2)$ Two point basis

$S=2$, when $h+k+l$ even

$S=0$, when $h+k+l$ odd, systematical absence



For BCC: $(0,0,0), (1/2, 1/2, 1/2)$ Two point basis

The diagram shows a BCC unit cell with blue spheres representing atoms. The two-point basis is highlighted: one sphere at the origin $(0,0,0)$ and another sphere at the body center $(1/2, 1/2, 1/2)$.

$$S_k = \sum_{j=1}^2 e^{-iK \cdot d_j} = e^{-iK \cdot 0} + e^{-iK \cdot \frac{1}{2}(\vec{x} + \vec{y} + \vec{z})}$$

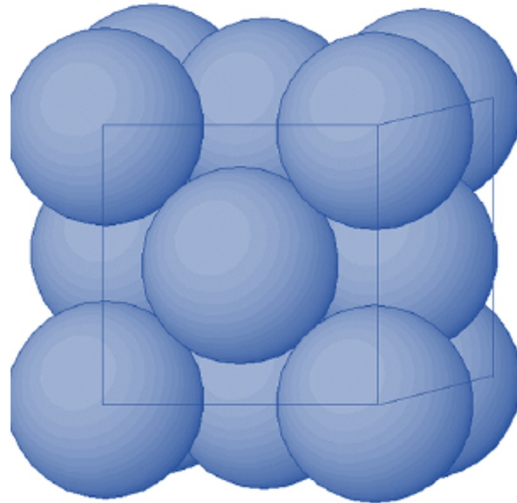
$$= 1 + e^{-i\pi(h+k+l)}$$

$$= 1 + (-1)^{h+k+l}$$

$S=2$, when $h+k+l$ even

$S=0$, when $h+k+l$ odd, systematical absence

Structure factor FCC



FCC

**S=4 when $h+k, k+l, h+l$
all even (h, k, l all even/odd)**

S=0, otherwise

$$S_k = 1 + e^{-i\pi(h+k)} + e^{-i\pi(h+l)} + e^{-i\pi(l+k)}$$

Rules for diffraction peaks of cubic lattices

Rules for Determining the Diffracting (hkl) Planes in Cubic Crystals				
Bravais lattice	Reflections present	Reflections absent	Observable diffraction peaks	
BCC FCC	$(h + k + l) = \text{even}$ (h, k, l) all odd or all even	$(h + k + l) = \text{odd}$ (h, k, l) not all odd or all even	$h^2 + k^2 + l^2$ Ratio	
Miller indices of the Diffracting Planes for BCC and FCC Lattices			SC: 1,2,3,4,5,6,8,9,10,11,12..	
Cubic planes (hkl)	$h^2 + k^2 + l^2$	Simple cubic Sum $\Sigma(h^2 + k^2 + l^2)$	FCC	BCC
{100}	$1^2 + 0^2 + 0^2$	1	...	110
{110}	$1^2 + 1^2 + 0^2$	2
{111}	$1^2 + 1^2 + 1^2$	3	111	...
{200}	$2^2 + 0^2 + 0^2$	4	200	200
{210}	$2^2 + 1^2 + 0^2$	5
{211}	$2^2 + 1^2 + 1^2$	6	...	211
{220}	$2^2 + 2^2 + 0^2$	8	220	220
{221}	$2^2 + 2^2 + 1^2$	9
{310}	$3^2 + 1^2 + 0^2$	10	...	310

BCC: 2,4,6,8,10, 12, 14....
(1,2,3,4,5,6,7...)

FCC: 3,4,8,11,12,16,24....

Observable diffraction peaks

$$h^2 + k^2 + l^2$$

Ratio

SC: 1,2,3,4,5,6,8,9,10,11,12..

BCC: 2,4,6,8,10, 12, 14....
(1,2,3,4,5,6,7...)

FCC: 3,4,8,11,12,16,24....

$$d_{hkl} = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$$

$$2d \sin \theta = n\lambda$$

$$\sin^2 \theta = \frac{\lambda^2}{4a^2} (h^2 + k^2 + l^2)$$

$$d_{hkl} = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$$

$$2d \sin \theta = n\lambda$$

Polyatomic structures

Polyatomic Structures

$$S_k = \sum_{j=1}^n f_j(k) e^{-i\mathbf{K} \cdot \vec{d}_j}$$

$$\mathbf{K} = h\mathbf{a}^* + k\mathbf{b}^* + l\mathbf{c}^*$$
$$\vec{d}_j = x\vec{a} + y\vec{b} + z\vec{c}$$

f_j : atomic scattering factor
 \propto No. of electrons

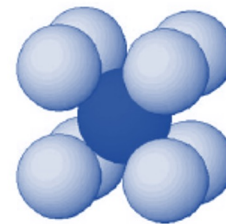
$$S_k = \sum_{j=1}^n f_j e^{-i\mathbf{K} \cdot \vec{d}_j} = \sum_{j=1}^n f_j e^{-2\pi i(hx+ky+lz)}$$

Simple Cubic Lattice

Caesium Chloride
(CsCl) is **primitive**
cubic

Cs (0,0,0)

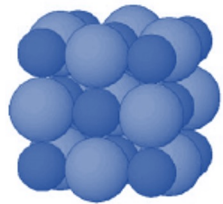
Cl (1/2,1/2,1/2)



$$S_k = f_{Cs} + f_{Cl} e^{-i\pi(h+k+l)}$$

What about CsI?

Polyatomic structures



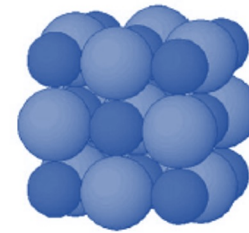
FCC Lattices

Sodium Chloride
(NaCl)

Na: (0,0,0)(0,1/2,1/2)(1/2,0,1/2)(1/2,1/2,0)

Cl: (1/2,1/2,1/2) (1/2,0,0)(0,1/2,0)(0,0,1/2)

Add (1/2,1/2,1/2)



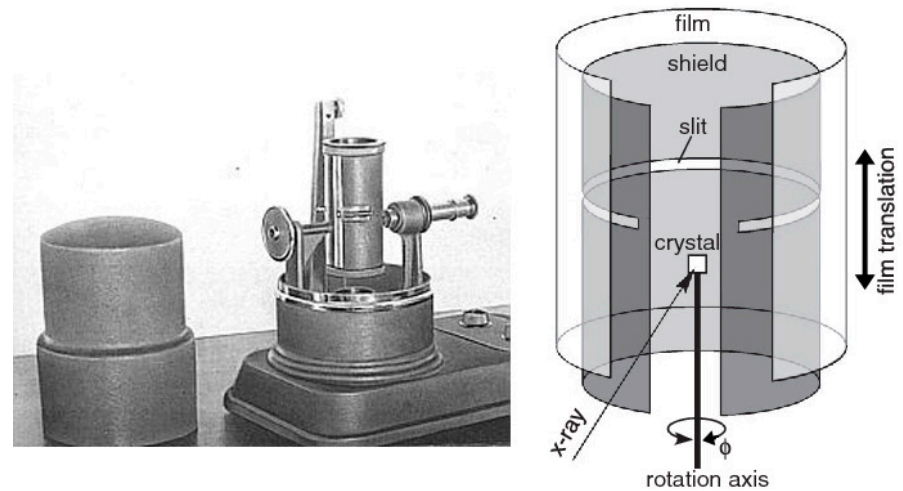
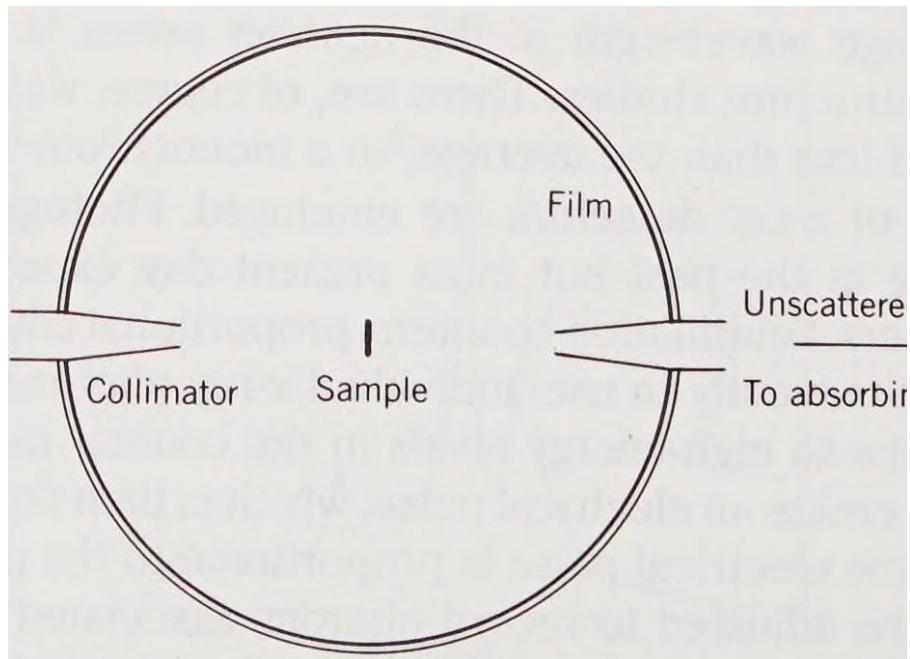
$$S_k = [f_{Na} + f_{Cl}e^{-i\pi(h+k+l)}][1 + e^{-i\pi(h+k)} + e^{-i\pi(h+l)} + e^{-i\pi(l+k)}]$$

S=4 (f_{Na} + f_{Cl}) when h, k, l, all even;

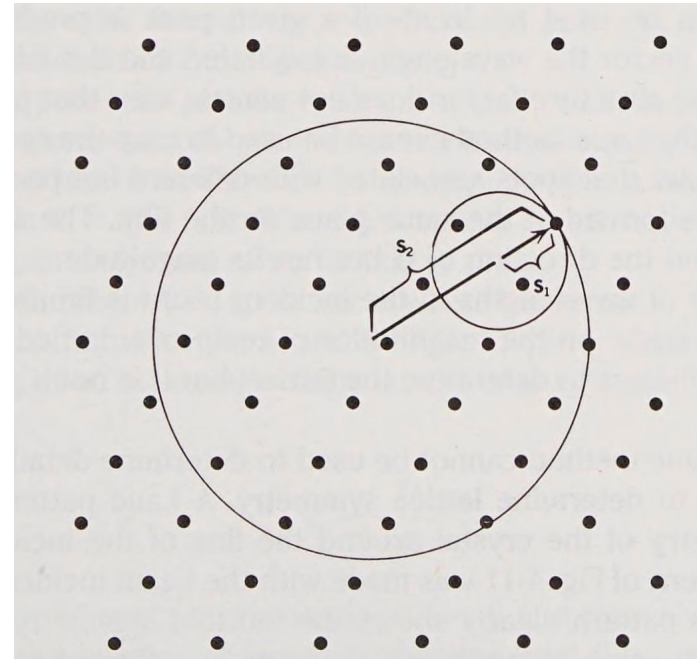
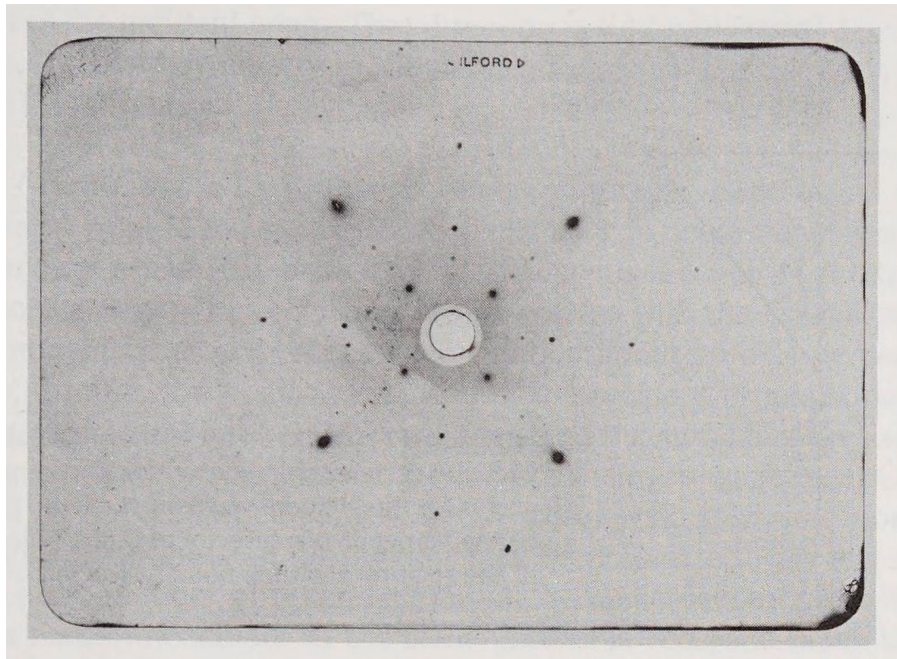
S=4(f_{Na}-f_{Cl}) when h, k, l all odd

S=0, otherwise

Laue method: cylindrical diffraction camera

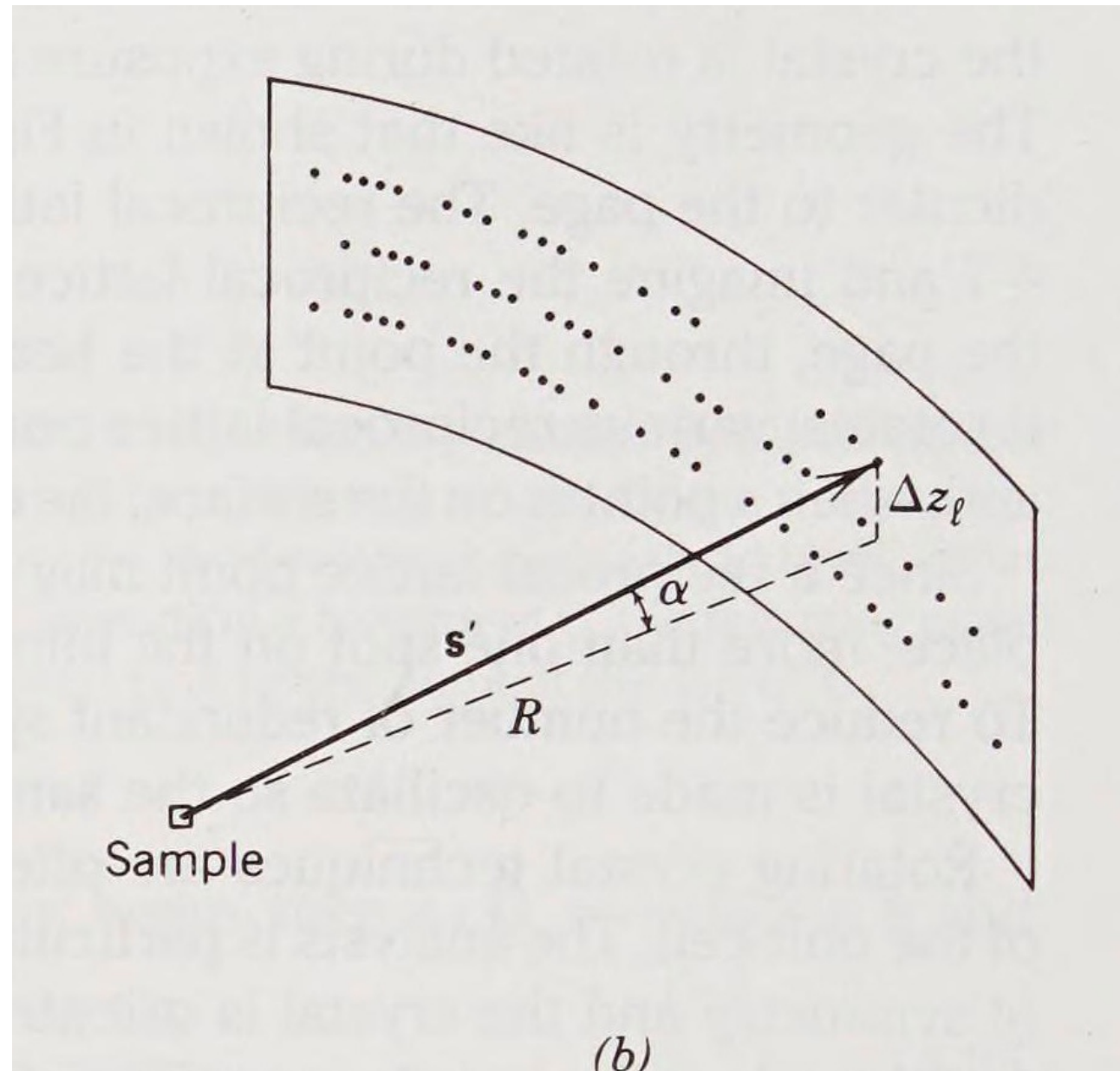


Laue pattern and Ewald construction

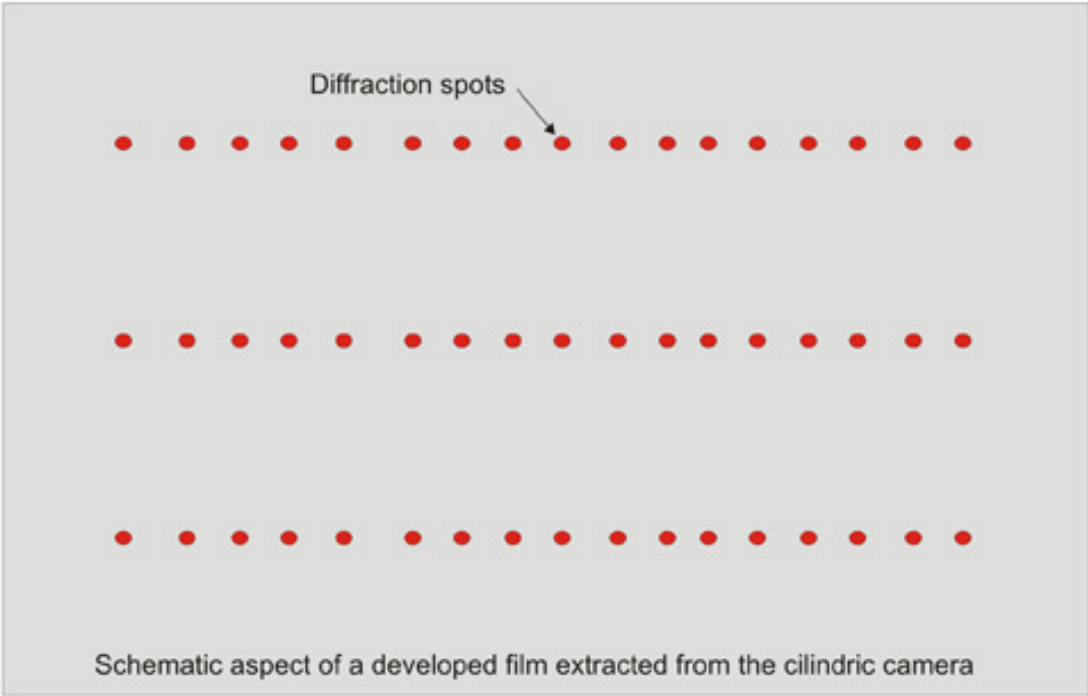
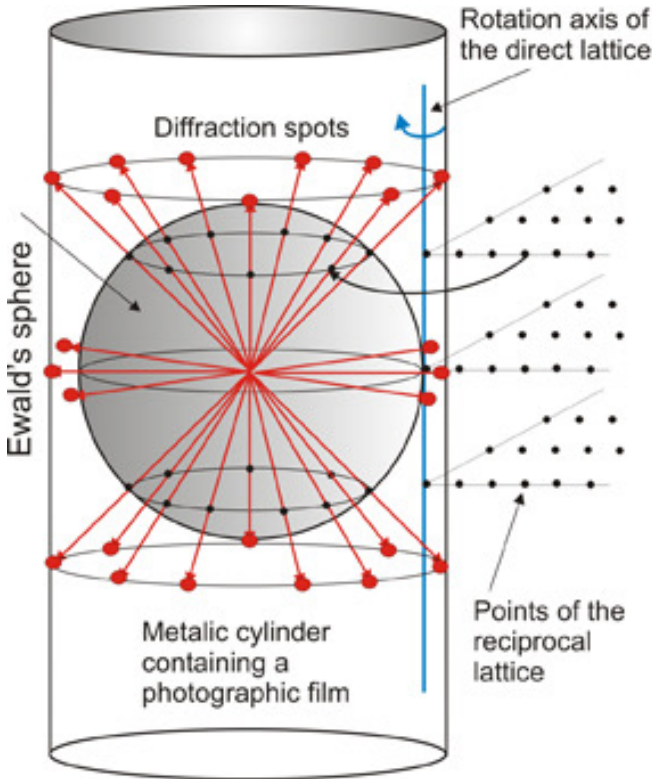


Rotating crystal

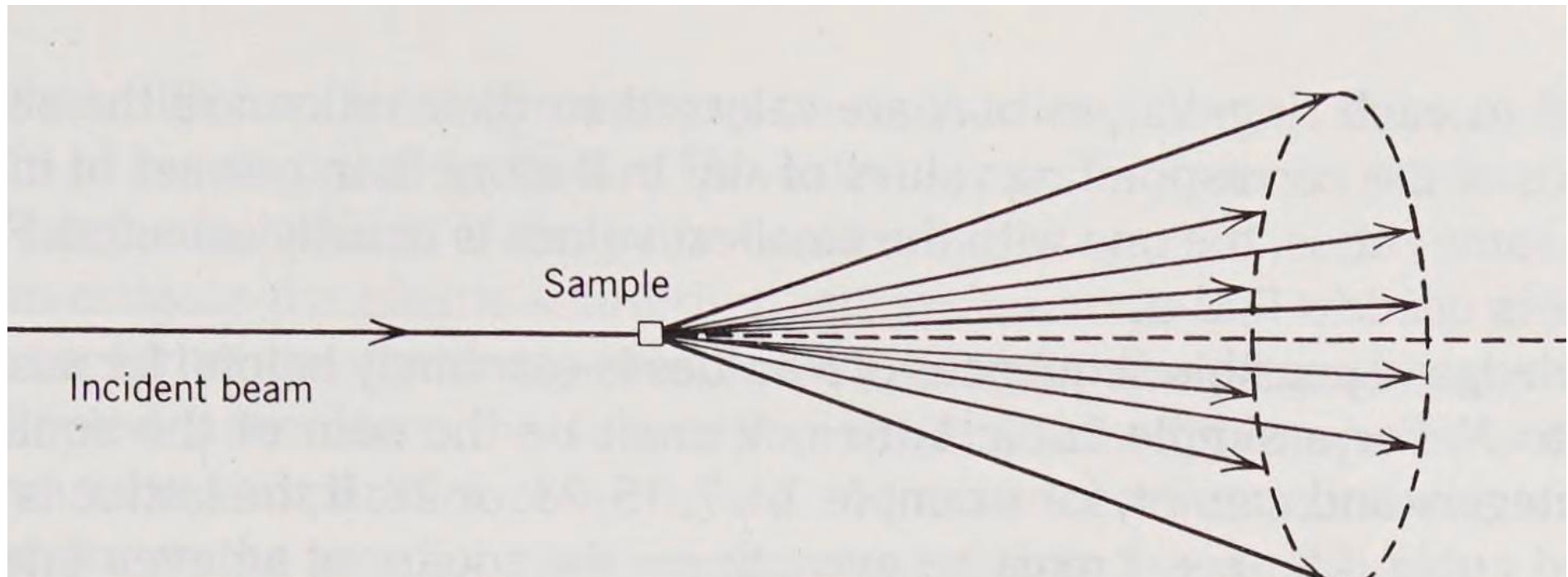
Diagram of intensity peaks



Rotating crystal



Powder sample



Powder pattern

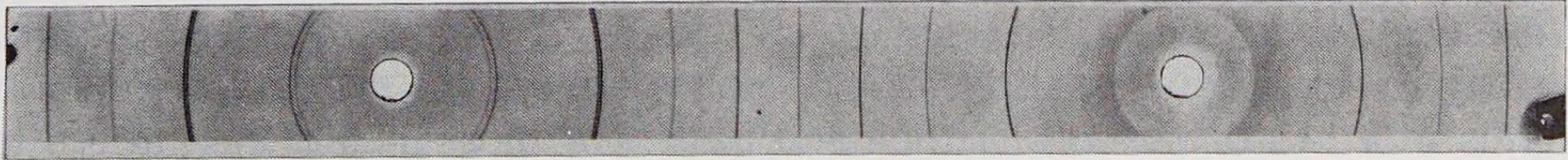


FIGURE 4-16 Powder pattern for tungsten. Each ring corresponds to a different scattering angle. The pattern consists of two sets of rings, for beams scattered in the back (on the left) and forward (on the right) directions, respectively. (From Walter Kiszewick, unpublished. Used with permission.)