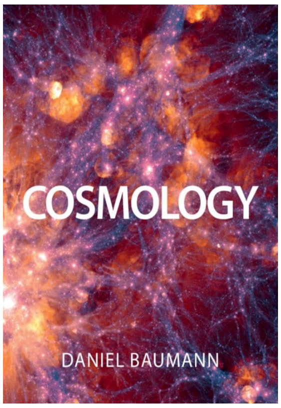
Universo Primitivo
2024-2025 (1º Semestre)

Mestrado em Física - Astronomia

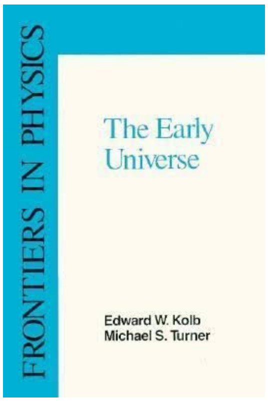
Chapter 3

3. Thermodynamics in an expanding universe
    - Natural Units;
    - Classification and properties of elementary particles;
    - Thermal evolution at equilibrium:
        - Density of states and macroscopic properties
        - Number density, energy density and pressure
            - Ultra-relativistic limit
            - Non-relativistic limit
    - Effective number of degrees of freedom
        - Internal degrees of freedom of particles according to the standard model of particle physics
        - Evolution of relativistic degrees of freedom
    - Entropy at equilibrium
        - Effective number of degrees of freedom in entropy;
        - Entropy conservation and its consequences;
        - Entropy and Temperature – time scaling for relativistic particles
    - Key events in the thermal history of the Universe

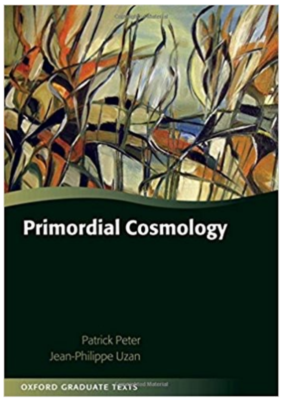
References



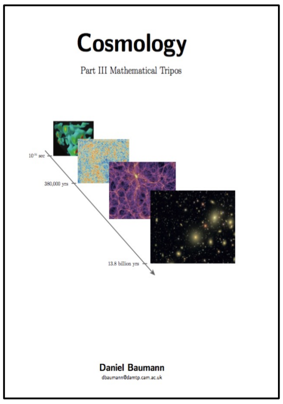
Ch. 3



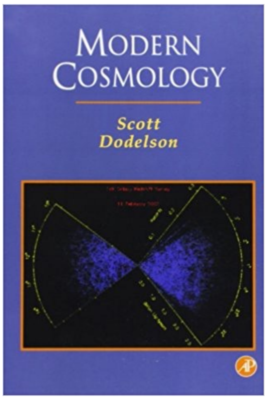
Ch. 3



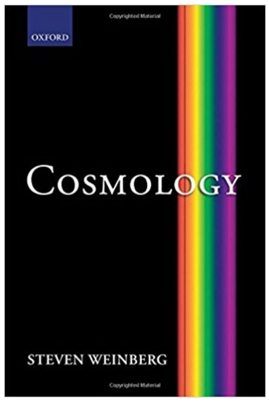
Ch. 4; Sec. 2.3



Ch. 3



Ch. 2



Ch. 2

Natural Units

In Particle Physics and Cosmology, the expression “**natural units**” usually refers to setting the following fundamental constants equal to unity:

$$c = k_B = \hbar = 1$$

These are the speed of light, the Boltzmann constant and the Planck constant ($\hbar = h/2\pi$). Note that setting $\hbar = 1$, means that $h = 2\pi$.

As a consequence, the following fundamental properties (**time; length, temperature and mass**) can be written in **units of energy** (usually expressed in GeV, MeV, keV):

$$\begin{aligned} 1\ \mathrm{s} &= 1.5 \times 10^{24}\ \mathrm{GeV}^{-1}, \\ 1\ \mathrm{m} &= 5 \times 10^{15}\ \mathrm{GeV}^{-1}, \\ 1\ \mathrm{K} &= 8.6 \times 10^{-14}\ \mathrm{GeV} = 8.6 \times 10^{-5}\ \mathrm{eV}, \\ 1\ \mathrm{kg} &= 5.6 \times 10^{26}\ \mathrm{GeV}. \end{aligned}$$

where $1\mathrm{eV} = 1.6 \times 10^{-19}\mathrm{J}$ $\Rightarrow$ $1\mathrm{J} = 6.2 \times 10^{9}\mathrm{GeV}$. $1\ \mathrm{J} = 1\ \mathrm{kg\ m^2 s^{-2}}$

Natural Units

To prove these, use the definitions of the following constants in the IS system and the definition of electron volt in Jules.

$c = 3 \times 10^8$ m s^{-1},	velocidade da luz no vácuo;
$G = 6.67 \times 10^{-11}$ m^3 kg^{-1}s^{-2},	constante gravitacional;
$h = 6.6 \times 10^{-34}$ J s,	constante de Planck;
$e = 1.6 \times 10^{-19}$ C,	carga elementar;
$k_B = 1.38 \times 10^{-23}$ J K^{-1},	constante de Boltzmann.

Example: of the mass of known particles in MeV:

Espécie	Símbolo	Massa (MeV)	Carga (e)
Protão	p	938.3	+1
Neutrão	n	939.6	0
Electrão	e^-	0.511	-1
Neutrinos	ν_e, ν_μ, ν_τ	?	0
Fotão	γ	0	0
Matéria Escura	–	?	0?
Energia Escura	–	?	?

Classification of elementary particles

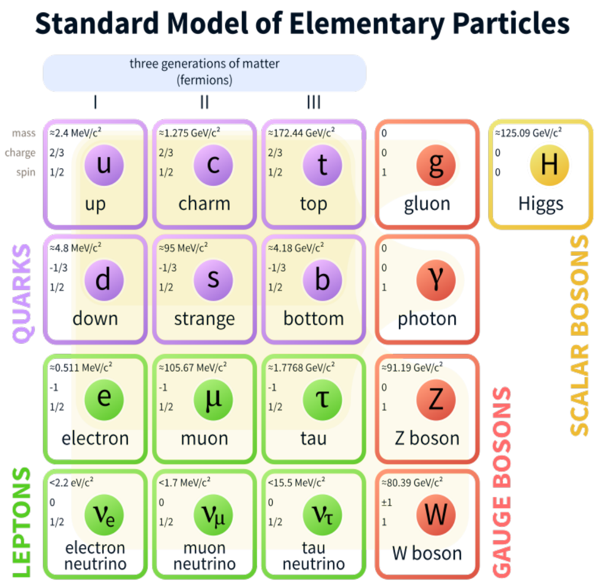
The Standard Model of Particle Physics (SMPF) predicts various families of particles some of them are **fundamental** and other “composite” particles.

Fundamental particles are not known to have internal structure. **Composite particles** have internal structure (i.e. are made of other particles).

All particles of the SMPF can be classified in the following way:

Name		Spin	Examples
Hadrons	Baryons = qqq	$n + \frac{1}{2}$	$p^+, n^0, \Delta, \Lambda, \Sigma, \Omega, \Xi \cdots$
	Mesons = $q\bar{q}$	n	$\pi^{0,\pm}, K^{0,\pm}, J/\psi, D^0, B^0, \eta, \cdots$
Leptons		$\frac{1}{2}$	$e^-, \nu_e, \mu^-, \nu_\mu, \tau^-, \nu_\tau.$
Gauge fields		1	$\gamma, Z^0, W^\pm, g^a.$

Classification of elementary particles



Gauge Fields (exchange Bosons):

Are fundamental particles that mediate interactions:

- Photon γ – electromagnetic;
- 8 gluons g – strong interaction
- Z and $W^{\pm}$ – weak interaction
- *Graviton? ($h_{\mu\nu}$)* – gravitational interaction (quantum gravity)

Leptons:

Are fundamental particles that interact via the **electromagnetic** and **weak** forces.

- Come in doublets with respect to the weak force
- Only distinguishable by the mass
- Stable doublet: is the electron/electron neutrino

Hadrons:

Have internal structure and interact via **all types of forces**.

- are made of **quarks**, confined in sets of 2 (Mesons) or 3 (Baryons) particles: up, down; charm, strange; top; bottom (u, d, c, s, t, b)

Scalar Higgs Boson

- Higgs Field: The Higgs mechanism is a process describing the **Electroweak symmetry breaking** and the **generation of the mass** of all fermions and massive bosons.

Thermal evolution at equilibrium

Fundamental assumptions about the primordial universe:

- All **fluid species** are assumed to **behave as ideal fluids**.

- **Thermal equilibrium** of a fluid species **may be established** whenever the particles' interaction rate, $\Gamma(t)$, (expressed as *the number of interaction events per unit of time*) is larger than the expansion rate of the Universe, $H(t) = \dot{a}/a$:

$$\Gamma(t) \gg H(t)$$

- The way to describe a fluid component is through its **distribution function** $f(\boldsymbol{x}, \boldsymbol{p}, E, t)$. It gives the mean **number density of particle states in the position, $\boldsymbol{x} \pm \boldsymbol{dx}$, with momentum, $\boldsymbol{p} \pm \boldsymbol{dp}$**.

- In classical mechanics f is defined as the number of particles per **phase space** volume: $dN = f(\boldsymbol{x}, \boldsymbol{p}, E, t)\, d^3x\, d^3p$

- If space is **homogeneous,** the distribution function must be *independent of* $\boldsymbol{x}$. Moreover, assuming **isotropy**, f must be a function of $p = |\boldsymbol{p}|$, so $f = f(p, E, t)$.

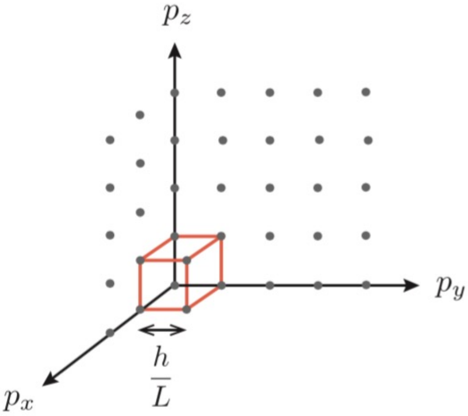
Thermal evolution at equilibrium

From quantum states to microscopic properties:

Under the assumptions of *homogeneity* and *isotropy*, the number of particles $dN = f(\boldsymbol{x}, \boldsymbol{p}, E, t)\, d^3x\, d^3p$ does not depend on $\boldsymbol{x}$ and is only a function of $p = |\boldsymbol{p}|$.

The number density of particles states is defined as:

$$n = \frac{g}{(2\pi)^3} \int \mathrm{d}^3p\, f(p)$$



Likewise, one can obtain the energy density of particles in real space by weighting each momentum eigen-state by its energy, $E = \sqrt{m^2 + p^2}$, and therefore:

$$\rho = \frac{g}{(2\pi)^3} \int \mathrm{d}^3p\, f(p)\, E(p)$$

The computation of the pressure of particles results in a similar way (This can be derived using statistical mechanics assuming a gas of weakly interacting particles, see slides 14-15).

$$P = \frac{g}{(2\pi)^3} \int \mathrm{d}^3p\, f(p)\, \frac{p^2}{3E}$$

g is called the "**internal degrees of freedom**" of each momentum state

Thermal evolution at equilibrium

The phase-space of "free" particles in Quantum physics:

The Uncertainty Principle (U.P.) introduces a fundamental uncertainty in the $\{\vec{x}, \vec{p}\}$ phase space. For example, taking the U.P. in its form of 1927, and considering only the x and p_x components of position and momentum of the phase space one has:

$$\Delta x\, \Delta p_x \gtrsim h \iff \Delta x\, \Delta p_x \gtrsim 2\pi\hbar$$

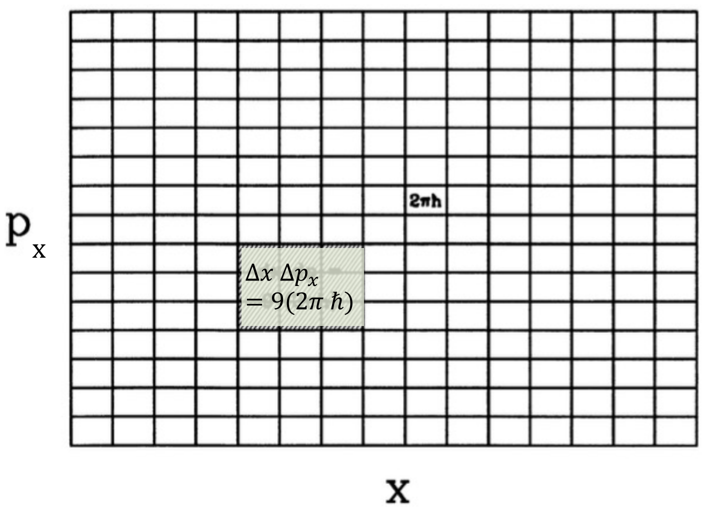


Figure 2.4. Phase space of position and momentum in one dimension. Volume of each cell is $2\pi\hbar$, the smallest region into which a particle can be confined because of Heisenberg's principle. Shaded region has infinitesmal volume $dxdp$. This covers nine cells. To count the appropriate number of cells, therefore, the phase space integral must be $\int dxdp/(2\pi\hbar)$.

This can be seen as setting fundamental uncertainty **"confinement" regions** in the $\{x, p_x\}$ plane of the phase space. Each region, representing a possible $\boldsymbol{\{x, p_x\}}$ **– state,** has a minimum area set by the (1-Dim) U.P.

$$\Delta x\, \Delta p = 2\pi\hbar$$

Taking the U.P. in all 3-Dim of space and momentum $\{\boldsymbol{x}, \boldsymbol{p}\}$, the **confinement regions** will extend over a 6-dim volume given by ($\boldsymbol{\Delta x \Delta p} \equiv \Delta x \Delta y \Delta z \Delta p_x \Delta p_y \Delta p_z$):

$$\boldsymbol{\Delta x\, \Delta p} = (2\pi\hbar)^3$$

So, the **total number of confinement regions (states)** in the phase space can be estimated as (the 2nd equality sets $\hbar$=1):

$$\int \frac{\boldsymbol{dxdp}}{(2\pi\hbar)^3} = \frac{\mathbf{1}}{(2\pi)^3} \int \boldsymbol{dxdp}$$

Phase space density

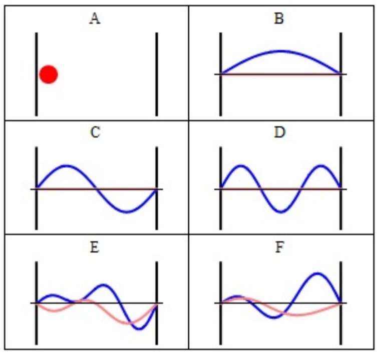
Thermal evolution at equilibrium

The phase-space of "free" particles in Quantum physics:

In quantum mechanics the **momentum operator ($\hat{p} = i\hbar\nabla$) eigenstates** of a free particle inside a box of volume, $V = L^3$, has a **discrete spectrum of momentum / energy eigenstates,** described by the (time-independent) Schrödinger equation:

$$\frac{p^2}{2m}\psi = -\frac{\hbar^2\nabla^2}{2m}\psi = E\psi \Leftrightarrow \nabla^2\psi = -k^2\psi$$

where, $k^2 = 2mE/\hbar^2$ and $p = \hbar k$.



The **1D** solution for the boundary condition $\psi(0) = \psi(L) = 0$ is of the form $\psi(x) = A\sin(k_n x)$**,** where:

$$k_n = n\pi/L, \qquad \text{with}\, n > 0$$

The energy of each mode *n* is:

$$E_n = \frac{p_n^2}{2m} = \frac{\hbar^2 k_n^2}{2m} = \frac{\hbar^2}{2m}\frac{\pi^2}{L^2}n^2$$

In **3D**, the possible energy and momentum states are ($\vec{N} = (n_x, n_y, n_z)$):

$$\vec{p}_{\vec{N}} = \hbar\vec{k}_{\vec{N}} = \frac{\hbar\pi}{L}(n_x, n_y, n_z)$$

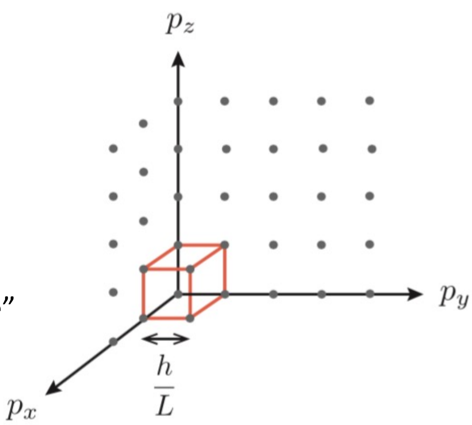
$$E_{\vec{N}} = \frac{p_{\vec{N}}^2}{2m} = \frac{\hbar^2 k_{\vec{N}}^2}{2m} = \frac{\hbar^2}{2m}\frac{\pi^2}{L^2}(n_x^2 + n_y^2 + n_z^2)$$

Thermal evolution at equilibrium

The phase-space of a species in Quantum physics:

Therefore, the allowed **momentum eigenstates** in one octant of the $\vec{N} = (n_x, n_y, n_z)$ space is:

$$\vec{p}_{\vec{N}} = \frac{\pi\hbar}{L}\vec{N} \quad \Leftrightarrow \quad \vec{N} = \frac{L}{\pi\hbar}\vec{p}_{\vec{N}}$$



So, the volume elements $d^3N \equiv d\vec{N}$ and $d^3p \equiv d\vec{p}_{\vec{N}}$, in both $\vec{N}$ and $\vec{p}$ spaces are related by:

$$d^3N = \left(\frac{L}{\pi\hbar}\right)^3 d^3p$$

One must keep in mind that all possible $\vec{N}$ are in the "positive" octant. To compute the total number of possible states one can either integrate over the "positive" octant of $\vec{p}$ or the whole momentum space and divide by 8:

$$N = \frac{1}{8}\int\left(\frac{L}{\pi\hbar}\right)^3 d^3p = \int\frac{V}{(2\pi\hbar)^3}d^3p = \frac{V}{(2\pi)^3}\int d^3p$$

where **we have set $\hbar = 1$ in the last equality**. So, the number density of possible states is:

$$\frac{N}{V} = \frac{1}{(2\pi)^3}\int d^3p$$

Phase space density in natural units

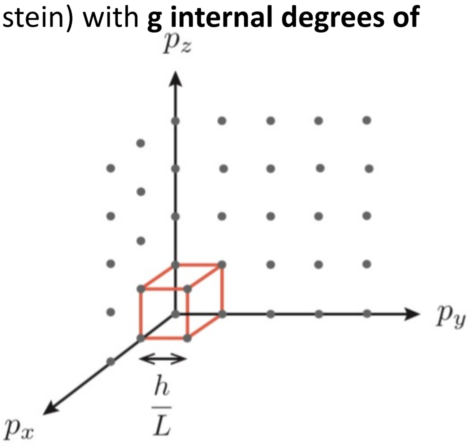
Note that the integral is done over the whole $\vec{p}$ space. Note also that $L^3 = V = \int d^3x$.

Thermal evolution at equilibrium

The phase-space of a species in Quantum physics:

So, the **number density of momentum states** for a particle species with a **state occupation distribution function** $f(\boldsymbol{p}, E, t)$ (either Fermi-Dirac or Bose-Einstein) with **g internal degrees of freedom** is:

$$n = \frac{N}{V} = \frac{g}{(2\pi)^3} \int f(\boldsymbol{p}, E, t)\, d^3p$$



From this, one can obtain the energy density of that particle species just by weighting each momentum eigen-state by its energy, $E = \sqrt{m^2 + p^2}$:

$$\rho = \frac{g}{(2\pi)^3} \int f(\boldsymbol{p}, E, t)\, E(p)\, d^3p$$

Note that m is the particle's rest mass.
For the pressure one obtains (see the following 2 slides) a similar expression:

$$P = \frac{g}{(2\pi)^3} \int f(\boldsymbol{p}, E, t) \frac{p^2}{3E(p)}\, d^3p$$

The internal degrees of freedom, g, accounts for (quantum) particle properties that do not impact on their momentum eigen-states.

Thermal evolution at equilibrium

Derivation of (done in class),

$$P = \frac{g}{(2\pi)^3} \int \mathrm{d}^3p\, f(p)\, \frac{p^2}{3E}$$

(from Baumann lectures Chap. 3.2)

Let's assume a gas of weakly interacting particles in statistical mechanics.

Consider the area element dA, in the figure on the left. Particles move with $E(|v|)$.

The number of particles in the shaded volume $dV = |v|dt\, dA_s = |v|dt\, d\Omega R^2$ is:

$$\mathrm{d}N = \frac{g}{(2\pi)^3} f(E) \times R^2 |\boldsymbol{v}| \mathrm{d}t\, \mathrm{d}\Omega$$

Not all particles in dV will hit dA.
*Only a fraction of this particles, with $\hat{v}.\hat{\boldsymbol{n}} = cos(\theta)$, i.e. with the direction, $\boldsymbol{v}$, will hit dA. So, **assuming isotropy**, the number of particles arriving on dA is:*

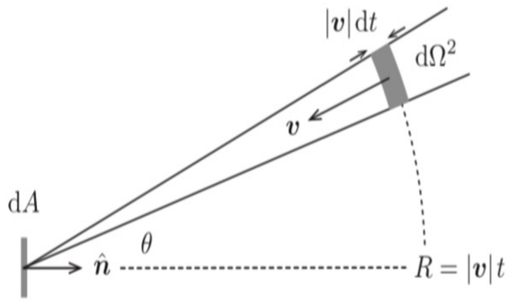


Figure 3.3: Pressure in a weakly interacting gas of particles

$$\mathrm{d}N_A = \frac{|\hat{\boldsymbol{v}} \cdot \hat{\boldsymbol{n}}|\, \mathrm{d}A}{4\pi R^2} \times \mathrm{d}N = \frac{g}{(2\pi)^3} f(E) \times \frac{|\boldsymbol{v} \cdot \hat{\boldsymbol{n}}|}{4\pi}\, \mathrm{d}A\, \mathrm{d}t\, \mathrm{d}\Omega$$

Thermal evolution at equilibrium

(Derivation continuation…)

$$\mathrm{d}N_A = \frac{|\hat{\boldsymbol{v}}\cdot\hat{\boldsymbol{n}}|\,\mathrm{d}A}{4\pi R^2}\times \mathrm{d}N = \frac{g}{(2\pi)^3}f(E)\times\frac{|\boldsymbol{v}\cdot\hat{\boldsymbol{n}}|}{4\pi}\,\mathrm{d}A\,\mathrm{d}t\,\mathrm{d}\Omega$$

*If these dN_A particles **collide elastically** at dA, each particle transfers a momentum $2|\boldsymbol{p}.\hat{n}|$ (because the particle is assumed to collide elastically and is reflected with the same angle of impact).*

So, the pressure dP (defined as force / area = momentum / time / area) by these particles at dA is:

$$\begin{aligned}\mathrm{d}P(|\boldsymbol{v}|) &= \int \frac{2|\boldsymbol{p}\cdot\hat{\boldsymbol{n}}|}{\mathrm{d}A\,\mathrm{d}t}\,\mathrm{d}N_A \\ &= \frac{g}{(2\pi)^3}f(E)\times\frac{p^2}{2\pi E}\int\cos^2\theta\,\sin\theta\,\mathrm{d}\theta\,\mathrm{d}\phi \\ &= \frac{g}{(2\pi)^3}\times f(E)\,\frac{p^2}{3E}\end{aligned}$$

where $|\boldsymbol{v}| = |\boldsymbol{p}|/E$ and the integration is made over the hemisphere of particles moving towards dA (i.e. with $\hat{\boldsymbol{v}}\cdot\hat{\boldsymbol{n}} \equiv -\cos\theta < 0$)

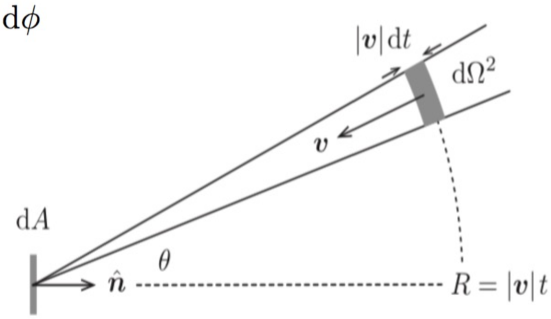


Figure 3.3: *Pressure in a weakly interacting gas of particles*

Thermal evolution at equilibrium

Local Kinetic equilibrium

If particles **exchange momentum and energy** in an **efficient way**, the system is said to be in **kinetic equilibrium**. If the system achieves a **maximum entropy state**, then particles are distributed according to the **Fermi-Dirac** or **Bose-Einstein** distribution functions:

$$f(p) = \frac{1}{e^{(E(p)-\mu)/T}\pm 1}$$

$+$ Fermions
$-$ Bosons

Where T is the temperature of the system and μ is the chemical potential defined as the change of energy with respect of the number of particles, at constant entropy, volume, and number other particle species.

$$\mu_i = \left(\frac{\partial U}{\partial N_i}\right)_{S,V,N_{j\neq i}} \quad \text{or} \quad \mu_i = -T\left(\frac{\partial S}{\partial N}\right)_{U,V,\,N_{j\neq i}}$$

At low temperature $T \ll E-\mu$ both distributions reduce to the "**Maxwell-Boltzmann"** distribution:

$$f(p) \approx e^{-(E(p)-\mu)/T}$$

Thermal evolution at equilibrium

Local Chemical equilibrium

- If a particle species, $\boldsymbol{i}$, is in **chemical equilibrium**, then $\boldsymbol{\mu_i}$ is related to the other species chemical potential. For example, if one has the following interaction (reaction) among species:

$$1 + 2 \leftrightarrow 3 + 4 \quad \text{then} \quad \mu_1 + \mu_2 = \mu_3 + \mu_4$$

- Photons have chemical potential equal to zero, i.e. $\boldsymbol{\mu_\gamma = 0}$, because **the number of photons is not conserved**. For example, in a double scattering interaction one has

$$e^- + \gamma \leftrightarrow e^- + \gamma + \gamma$$

- This implies that a particle, $\boldsymbol{X}$, and its antiparticle, $\boldsymbol{\bar{X}}$, $(X + \bar{X} \leftrightarrow \gamma + \gamma)$ have symmetric chemical potentials $\boldsymbol{\mu_X = -\mu_{\bar{X}}}$.

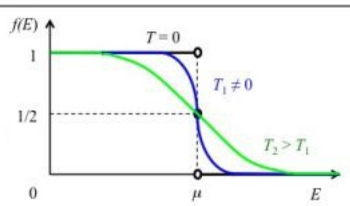
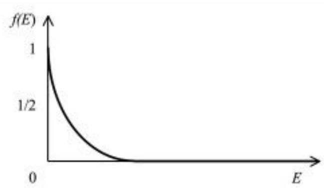
Local Thermal equilibrium

- **Thermal equilibrium** is achieved for **species which are both in kinetic and chemical equilibrium**. These species then share the same temperature, $T_i = T$.

Thermal evolution at equilibrium

Particle distribution functions

Quantum Statistics Summary

	Fermi-Dirac distribution	Bose-Einstein distribution
Function	$f(E) = \frac{1}{\exp[(E-\mu)/k_B T] + 1}$	$f(E) = \frac{1}{\exp[(E-\mu)/k_B T] - 1}$
Energy Dependence		
Quantum Particles	Undistinguishable particles obeying to the Pauli's Principle: **only one particle per state**	Undistinguishable particles not subject to the Pauli's Principle: **many particles can occupy one state**
Spins	semi-integer spins	integer spins
Properties	At temperature of 0 K, each energy level is occupied by two Fermi particles with opposite spins. Examples: electron, proton, neutron…	At very low temperature, large numbers of Bosons fall into lowest energy state. Examples: photon, gluon, mesons…

Thermal evolution at equilibrium:

Assuming homogeneity and isotropy the integrals in the expression of n, ρ, and P (expressions in slides 9 or 13) can be easily computed using spherical coordinates in the momentum space. The integrations of the angular part of the momentum space give 4π, so (note that $E = \sqrt{m^2+p^2}$ and $f(p,E,t) = f(p,t)$):

$$n = \frac{g}{(2\pi)^3}\int f(p,E,t)\, d^3p \quad \Leftrightarrow \quad \boxed{n = \frac{g}{2\pi^2}\int_0^\infty f(p,t)\, p^2\, dp}$$

$$\rho = \frac{g}{(2\pi)^3}\int f(p,E,t)\, E(p)\, d^3p \quad \Leftrightarrow \quad \boxed{\rho = \frac{g}{2\pi^2}\int_0^\infty f(p,t)\, p^2 E(p)\, dp}$$

$$P = \frac{g}{(2\pi)^3}\int f(p,E,t)\frac{p^2}{3E(p)}\, d^3p \quad \Leftrightarrow \quad \boxed{P = \frac{g}{2\pi^2}\int_0^\infty f(p,t)\frac{p^4}{3E(p)}\, dp}$$

- In general, these expressions need to be solved numerically.
- However, for some cases of interest it is possible to derive analytical solutions.
- These are the cases of **ultra-relativistic particles ($m \ll T$)** and **non-relativistic** ($m \gg T$) with vanishing chemical potential ($\mu = 0$)

Thermal evolution at equilibrium:

Whenever the **chemical potential is zero** (photons) or **negligible** the number and energy densities are ($E = \sqrt{m^2+p^2}$):

$$n = \frac{g}{2\pi^2}\int_0^\infty \mathrm{d}p\, \frac{p^2}{\exp\left[\sqrt{p^2+m^2/T}\right]\pm 1}$$

$$\rho = \frac{g}{2\pi^2}\int_0^\infty \mathrm{d}p\, \frac{p^2\sqrt{p^2+m^2}}{\exp\left[\sqrt{p^2+m^2/T}\right]\pm 1}$$

Defining $\boldsymbol{x \equiv m/T}$ and $\boldsymbol{\xi \equiv p/T}$ these integrals can be written as

$$n = \frac{g}{2\pi^2}T^3\, I_\pm(x)\ , \qquad I_\pm(x) \equiv \int_0^\infty \mathrm{d}\xi\, \frac{\xi^2}{\exp\left[\sqrt{\xi^2+x^2}\right]\pm 1}$$

$$\rho = \frac{g}{2\pi^2}T^4\, J_\pm(x)\ , \qquad J_\pm(x) \equiv \int_0^\infty \mathrm{d}\xi\, \frac{\xi^2\sqrt{\xi^2+x^2}}{\exp\left[\sqrt{\xi^2+x^2}\right]\pm 1}$$

Which in some cases can be evaluated analytically using the Riemann-Zeta, ζ, and Gama, Γ, functions. The following integral equalities involving, ζ and Γ, are particularly useful:

$$\int_0^\infty \mathrm{d}\xi\, \frac{\xi^n}{e^\xi - 1} = \zeta(n+1)\,\Gamma(n+1)\ ,$$

$$\int_0^\infty \mathrm{d}\xi\, \xi^n e^{-\xi^2} = \tfrac{1}{2}\Gamma\left(\tfrac{1}{2}(n+1)\right)\ ,$$

Thermal evolution at equilibrium:

Ultra-relativistic limit: $x \to 0$ ($m \ll T$ and $\mu = 0$)

Let us start with the $I_{\pm}(x \to 0)$ for ultra-relativistic particles. For **Bosons** one has:

$$I_{-}(0) = \int_0^{\infty} d\xi \frac{\xi^2}{\exp\sqrt{\xi^2 + 0^2} - 1} = \zeta(2+1)\Gamma(2+1) = 2\zeta(3) \simeq 2.4$$

For **Fermions,** the integral $I_{+}(x \to 0)$ is not directly related with the Riemann-Zeta and Gama integrals. However, one can use the mathematical equality,

$$\frac{1}{e^{\xi}+1} = \frac{1}{e^{\xi}-1} - \frac{2}{e^{2\xi}-1}$$

and then apply the Riemann-Zeta and Gama integral forms:

$$I_{+}(0) = \int_0^{\infty} d\xi \frac{\xi^2}{\exp\sqrt{\xi^2 + 0^2} + 1} = \int_0^{\infty} d\xi \frac{\xi^2}{\exp\xi - 1} - \int_0^{\infty} d\xi \frac{2\xi^2}{\exp 2\xi - 1}$$

Making a variable change, $y = 2\xi$, in the last integral, one obtains:

$$I_{+}(0) = I_{-}(0) - 2\int_0^{\infty} dy \left(\frac{1}{2}\right)\frac{(y/2)^2}{\exp(y) - 1} = I_{-}(0) - 2\left(\frac{1}{2}\right)^3 I_{-}(0) = \left(1 - \frac{1}{4}\right)2\zeta(3)$$
$$= \frac{3}{2}\zeta(3) \simeq 3.6$$

Thermal evolution at equilibrium:

Ultra-relativistic limit: $x \to 0$ ($m \ll T$ and $\mu = 0$)

So, one obtains the following expressions for the **number density**:

$$n = \frac{\zeta(3)}{\pi^2} g T^3 \begin{cases} 1 & \text{bosons} \\ \frac{3}{4} & \text{fermions} \end{cases}$$

Doing a similar computation for the $J_{\pm}(0)$, it is possible to derive (exercise) the following expression for the **energy density**:

$$\rho = \frac{\pi^2}{30} g T^4 \begin{cases} 1 & \text{bosons} \\ \frac{7}{8} & \text{fermions} \end{cases}$$

To compute the pressure for ultra-relativistic particles, $x \to 0$, with $\mu = 0$, it is straightforward to show (exercise) that:

$$P = \frac{1}{3}\rho$$

Thermal evolution at equilibrium:

Non-relativistic limit: $x \gg 1$ ($m \gg T$ and $\mu = 0$)

For $x \gg 1$ ($m \gg T$) the number density integral gives the **same expression for Fermions and Bosons**:

$$I_{\pm}(x) \approx \int_0^\infty \mathrm{d}\xi \, \frac{\xi^2}{e^{\sqrt{\xi^2+x^2}}}$$

Most of the contribution to this integral comes from $\boldsymbol{\xi} \ll \boldsymbol{x}$. Expanding the square root, $x(1+\xi^2/x^2)^{1/2}$, in a Taylor expansion to the lowest order in $\boldsymbol{\xi}$ one obtains:

$$I_{\pm}(x) \approx \int_0^\infty \mathrm{d}\xi \, \frac{\xi^2}{e^{x+\xi^2/(2x)}} = e^{-x} \int_0^\infty \mathrm{d}\xi \, \xi^2 e^{-\xi^2/(2x)} = (2x)^{3/2} e^{-x} \int_0^\infty dy \; y^2 e^{-y^2}$$

The last integral is obtained after a change of variable $y^2 = \xi^2/2x$. It is related with the Gamma Function integral with $n = 2$ in slide 20. So, one gets:

$$I_{\pm}(x) \approx (2x)^{3/2} \, e^{-x} \left(\frac{1}{2} \, \Gamma\left(\frac{3}{2}\right)\right) = (2x)^{3/2} \, e^{-x} \left(\frac{1}{2}\frac{\sqrt{\pi}}{2}\right) = \sqrt{\pi} \, 2^{\left(\frac{3}{2}-2\right)} x^{3/2} \, e^{-x} = \sqrt{\frac{\pi}{2}} \, x^{3/2} \, e^{-x}$$

Which leads to (see next slide)

Thermal evolution at equilibrium:

Non-relativistic limit: $x \gg 1$ ($m \gg T$ and $\mu = 0$)

The **number density** ($n = g/(2\pi^2)T^3 I_{\pm}(x)$) of non-relativistic particles

$$\boxed{n = g \left(\frac{mT}{2\pi}\right)^{3/2} e^{-m/T}}$$

This translates to a Maxwell-Boltzmann like distribution. It tell us that *massive particles (in the plasma) are exponentially rare at low temperatures*.

For the **energy density**, at low temperature ($T \ll m$) one has:

$$E(p) = \sqrt{m^2 + p^2} \approx m + p^2/2m$$

The energy density integral can be obtained using this previous approximation, giving (exercise):

$$\boxed{\rho = mn + \frac{3}{2} nT}$$

The **pressure** can be also easily computed (exercise), giving

$$\boxed{P = nT}$$

Thermal evolution at equilibrium:

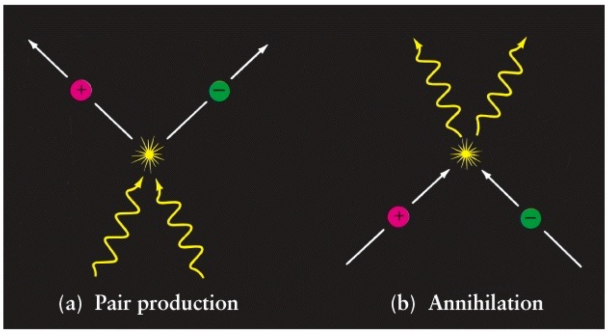
Non-relativistic limit: $x \gg 1$ ($m \gg T$ and $\mu = 0$)

From these expressions one concludes that:

- The densities and pressure of non-relativistic particles are strongly suppressed, by the exponential term $e^{-m/T}$, as temperature, T, drops below the particles mass, m. This is known as **Boltzmann suppression** is loosely referred to particle **'annihilations'**.

- These "annihilations" occur due to **changes in the interaction's cross sections** (that depend on temperature) involving the particle species. For example, for $X + \bar{X} \leftrightarrow \gamma + \gamma$ (**particle-antiparticle pair annihilation/production)** at low temperature (typically below $\sim m$), the thermal particle energies are not sufficient for pair production. So, with the expansion X and $\bar{X}$ will be supressed from the fluid.

- **Particle species suppressions** also occur due to other effects, such as **phase transitions** (as it happens to the less massive quarks in the QCD phase transition), particle decays,...



Thermal evolution at equilibrium:

Non-relativistic limit: $x \gg 1$ ($m \gg T$ and $\mu = 0$)

From the previous expressions one can concludes that (continuation):

- The transition from relativistic to non-relativistic behaviour is not instantaneous (in fact, about **80% of the "annihilations"** take place in the temperature range $\boldsymbol{T} \in [\boldsymbol{m/6}, \boldsymbol{m}]$).

- The **suppression of particles** from the fluid does not mean that all massive particles vanish from the universe. At the present temperature, all particle species with mass in the Universe are non-relativistic. The **suppression means that** these particle species (e.g. protons, electrons, atoms, dark matter) **are no longer "coupled"** to (or interacting with) the primordial fluid (that today is only made of CMB photons).

- For $m \gg T$ the energy density and pressure of non-relativistic particles (exercise)
  - $\rho = n\left(m + \frac{3}{2}T\right) \simeq nm$
  - $P = nT < \frac{3}{2}nT \ll nm \simeq \rho$

  This means that **non-relativistic particles have**, in general, **negligible pressure**. They behave as a "pressureless dust", (i.e. as $P = 0$ 'colissionless matter')

- Note also that $P = nT \Leftrightarrow PV = Nk_BT$ (in SI units) is the **ideal gas law**.

In a nutshell: **decoupled non-relativistic particles behave as a gas of pressureless matter.**

Thermal evolution at equilibrium:

Effective number of degrees of freedom of relativistic species

For a plasma of relativistic species, with bosons (labelled by *i*) and fermions (labelled by *j*) we have that:

$$\left.\begin{aligned} \rho_B^{(i)} &= \frac{\pi^2}{30} g_i T_i^4, \\ \rho_F^{(j)} &= \frac{7}{8}\frac{\pi^2}{30} g_j T_j^4 \end{aligned}\right\} \quad \rho_r = \sum_{i\ \mathrm{bosoes}} \frac{\pi^2}{30} g_i T_i^4 + \sum_{i\ \mathrm{fermioes}} \frac{7}{8}\frac{\pi^2}{30} g_i T_i^4$$

The total energy density of relativistic species can therefore be written as:

$$\rho_r = \sum_i \rho_i = \frac{\pi^2}{30} g_\star(T) T^4$$

where $T = T_\gamma$ is the photons temperature and g_* is the energy density ***effective number of degrees of freedom*** (in energy) of the fluid at temperature T:

$$g_* = \sum_{i\ \mathrm{bosões}} g_i \left(\frac{T_i}{T}\right)^4 + \frac{7}{8} \sum_{j\ \mathrm{fermiões}} g_j \left(\frac{T_j}{T}\right)^4$$

Thermal evolution at equilibrium:

Effective number of degrees of freedom of relativistic species

This expression allows that different species may not be in thermal equilibrium with the photon component. In fact, we can distinguish two situations:

- For **relativistic particles in thermal equilibrium with the photons** we have:

  $$g_\star^{th}(T) = \sum_{i=b} g_i + \frac{7}{8} \sum_{i=f} g_i$$

  when a species become non-relativistic, it is removed from the sums in g_*^{th}.
  So, when T is away from the "**mass thresholds**" of particles g_*^{th} is independent of temperature

- For **relativistic particles that are not in thermal equilibrium (or decoupling)** from the photon fluid, g_* varies with temperature:

$$g_\star^{dec}(T) = \sum_{i=b} g_i \left(\frac{T_i}{T}\right)^4 + \frac{7}{8} \sum_{i=f} g_i \left(\frac{T_i}{T}\right)^4$$

Thermal evolution at equilibrium:

Inventory of internal degrees of freedom of fundamental particles

type		mass	spin	g
quarks	$t, \bar{t}$	173 GeV	$\frac{1}{2}$	$2 \cdot 2 \cdot 3 = 12$
	$b, \bar{b}$	4 GeV		
	$c, \bar{c}$	1 GeV		
	$s, \bar{s}$	100 MeV		
	$d, \bar{s}$	5 MeV		
	$u, \bar{u}$	2 MeV		
gluons	g_i	0	1	$8 \cdot 2 = 16$
leptons	$\tau^\pm$	1777 MeV	$\frac{1}{2}$	$2 \cdot 2 = 4$
	$\mu^\pm$	106 MeV		
	$e^\pm$	511 keV		
	$\nu_\tau, \bar{\nu}_\tau$	< 0.6 eV	$\frac{1}{2}$	$2 \cdot 1 = 2$
	$\nu_\mu, \bar{\nu}_\mu$	< 0.6 eV		
	$\nu_e, \bar{\nu}_e$	< 0.6 eV		
gauge bosons	W^+	80 GeV	1	3
	W^-	80 GeV		
	Z^0	91 GeV		
	γ	0		2
Higgs boson	H^0	125 GeV	0	1

Internal degrees of freedom of fundamental particles in the Standard Model of Particle Physics:

- Massless spin-1 (photons and gluons): 2 polarizations
- Massive spin-1 ($W^\pm, Z^0$): 3 "polarizations"
- Massive spin-1/2 leptons ($e^\pm, \mu^\pm, \tau^\pm$): 2 spins
- Massive spin-1/2 quarks: 2 spin and 3 colour states
- Neutrinos/anti-neutrinos: 1 helicity state

So, **the internal degrees of freedom for relativistic bosons and fermions in equilibrium** are:

$g_b = 28$ photons (2), $W^\pm$ and Z^0 $(3 \cdot 3)$, gluons $(8 \cdot 2)$, and Higgs (1)

$g_f = 90$ quarks $(6 \cdot 12)$, charged leptons $(3 \cdot 4)$, and neutrinos $(3 \cdot 2)$

This gives:

$$g_\star = g_b + \frac{7}{8} g_f = 106.75$$

Thermal evolution at equilibrium:

Evolution of relativistic degrees of freedom (SMPP)

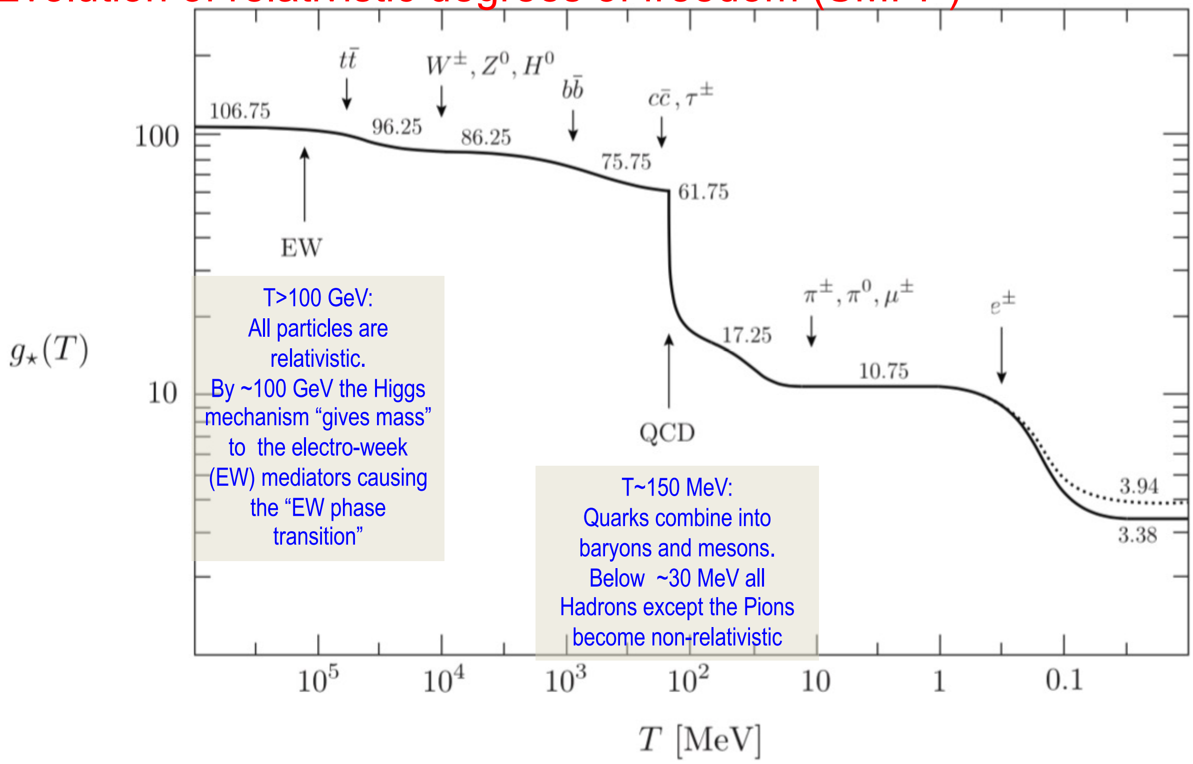


Figure 3.4: *Evolution of relativistic degrees of freedom $g_\star(T)$ assuming the Standard Model particle content. The dotted line stands for the number of effective degrees of freedom in entropy $g_{\star S}(T)$.*

Thermal evolution at equilibrium:

Entropy at equilibrium

From to the 1st law of thermodynamics for a particle species (and setting $\mu = 0$) one has:

$$\mathrm{d}U = TdS - PdV \qquad \text{and} \qquad dS = \frac{1}{T}dU + \frac{P}{T}dV$$

Since $U = \rho\, V$ one has:

$$dS = \frac{1}{T}d(\rho V) + \frac{P}{T}dV \;=\; \frac{1}{T}(Vd\rho + \rho dV) + \frac{P}{T}dV =$$

$$= \frac{V}{T}\,d\rho + \frac{\rho}{T}dV + \frac{P}{T}dV = \frac{V}{T}d\rho + \left(\frac{\rho+P}{T}\right)dV$$

So, one concludes that:

$$\boxed{\left(\frac{\partial S}{\partial \rho}\right)_V = \frac{V}{T}} \qquad \text{and} \qquad \boxed{\left(\frac{\partial S}{\partial V}\right)_\rho = \frac{\rho + P}{T}}$$

The entropy differential dS can also be written as:

$$dS = \frac{1}{T}(\,d(\rho V) + PdV\,) \;=\; \frac{1}{T}(\,d(\rho V) + PdV + VdP - VdP\,) =$$

$$= \frac{1}{T}(\,d(\rho V) + d(PV) - VdP\,) \;=\; \frac{1}{T}(d[(\rho + P)V] - VdP)$$

Thermal evolution at equilibrium:

Entropy at equilibrium

If one uses the Schawrtz theorem for the free Energy $dF = -SdT - PdV$ one has:

$$\frac{\partial^2 F}{\partial V \partial T} = \frac{\partial^2 F}{\partial T \partial V} \;\Leftrightarrow\; \frac{\partial}{\partial V}(-S) = \frac{\partial}{\partial T}(-P)$$

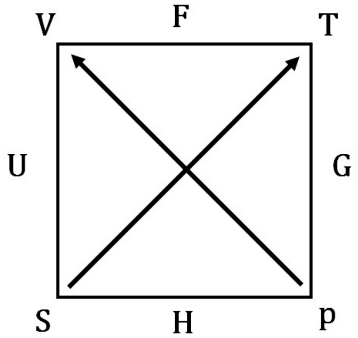
So, one concludes that

$$\boxed{\frac{\partial P}{\partial T} = \frac{\partial S}{\partial V} = \frac{\rho + P}{T}}$$

Where the last equality was established in the previous slide. This allows to go back to dS (and using $dP = (\partial P/\partial T)\, dT$) conclude that:

$$\begin{aligned} \mathrm{d}S &= \frac{1}{T}\Big(\mathrm{d}[(\rho + P)V] - V\mathrm{d}P\Big) \\ &= \frac{1}{T}\mathrm{d}[(\rho + P)V] - \frac{V}{T^2}(\rho + P)\,\mathrm{d}T \\ &= \mathrm{d}\left[\frac{\rho + P}{T}V\right] \end{aligned}$$

Max Born – Tisza diagram
of thermodynamic functions



$$dU = TdS - pdV$$
$$dF = -SdT - pdV$$
$$dH = TdS + Vdp$$
$$dG = -SdT + Vdp$$

$$\left(\frac{\partial T}{\partial V}\right)_S = -\left(\frac{\partial p}{\partial S}\right)_V \qquad \left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial p}{\partial T}\right)_V$$

$$-\left(\frac{\partial S}{\partial p}\right)_T = \left(\frac{\partial V}{\partial T}\right)_p \qquad \left(\frac{\partial T}{\partial p}\right)_s = \left(\frac{\partial V}{\partial S}\right)_p$$

Thermal evolution at equilibrium:

Entropy at equilibrium

This expression allows defining **entropy** and **entropy density** (or specific entropy), up to a constant, as:

$$\boxed{S = \frac{\rho + P}{T} V}$$

$$\boxed{s \equiv \frac{S}{V} = \frac{\rho + P}{T}}$$

The **specific entropy** of a **relativistic boson species** *i* can then be computed as (using the expressions of ρ_i, $P_i = \rho_i/3$, obtained earlier):

$$s_i = \frac{\pi^2}{30} g_i \left(1 + \frac{1}{3}\right) \frac{T_i^4}{T_i} = \frac{2\pi^2}{45} g_i T_i^3$$

Relativistic Bosons

A similar result holds for **relativistic fermion species,** i:

$$s_i = \frac{7}{8}\frac{\pi^2}{30}\left(1 + \frac{1}{3}\right) g_i T_i^3 = \frac{7}{8}\frac{2\pi^2}{45} g_i T_i^3$$

Relativistic Fermions

Thermal evolution at equilibrium:

Entropy at equilibrium

For a plasma of relativistic species, with bosons (labelled by *i*) and fermions (labelled by *j*) we have that:

$$s_B = \frac{2\pi^2}{45} g_i T_i^3$$

$$s_F = \frac{7}{8}\frac{2\pi^2}{45} g_i T_i^3$$

$$s = \sum_{i \text{ bosoes}} \frac{2\pi^2}{45} g_i T_i^3 + \sum_{i \text{ fermioes}} \frac{7}{8}\frac{2\pi^2}{45} g_i T_i^3$$

The total specific entropy of relativistic species can therefore be written as:

$$\boxed{s = \frac{2\pi^2}{45} g_{*s} T^3}$$

where $T = T_\gamma$ is the photons temperature and g_{*s} is the ***effective number of degrees of freedom in entropy*** of the fluid at temperature T:

$$\boxed{g_{*s} = \sum_{i \text{ bosões}} g_i \left(\frac{T_i}{T}\right)^3 + \frac{7}{8} \sum_{j \text{ fermiões}} g_j \left(\frac{T_j}{T}\right)^3}$$

Thermal evolution at equilibrium:

Entropy at equilibrium

One should note that g_{*s} is a function of $(T_i/T)^3$ whereas g_*, varies as $(T_i/T)^4$
This means that:

- Relativistic species in thermal equilibrium ($T_i = T$): $g_{*s} = g_*$
- Non-relativistic decoupling species ($T_i \neq T$): $g_{*s} \neq g_*$

In other words, if one writes

$$g_{\star S}(T) = g_{\star S}^{th}(T) + g_{\star S}^{dec}(T)$$

One has $g_{*S}^{th}(t) = g_*^{th}(T)$ for relativistic species in thermal equilibrium, and $g_{*S}^{dec}(T) \neq g_*^{dec}(T)$ for non-relativistic species in the process of decoupling from fluid.

Slide 30 shows both g_{*s} (dotted line) and g_* (solid line).

At high values of the degrees of freedom (i.e. higher temperatures) the curves appear on top of each other because the differences are small and only more visible at low T.

Thermal evolution at equilibrium:

Conservation of Entropy

A most important result about the evolution of the fluid in thermal equilibrium is that its **entropy remains constant with the expansion of the Universe** (as opposed to its energy density that decreases with time).
This can be proved by taking the **time derivative** of S:

$$\frac{\mathrm{d}S}{\mathrm{d}t} = \frac{\mathrm{d}}{\mathrm{d}t}\left[\frac{\rho + P}{T}V\right]$$
$$= \frac{V}{T}\left[\frac{\mathrm{d}\rho}{\mathrm{d}t} + \frac{1}{V}\frac{\mathrm{d}V}{\mathrm{d}t}(\rho + P)\right] + \frac{V}{T}\left[\frac{\mathrm{d}P}{\mathrm{d}t} - \frac{\rho + P}{T}\frac{\mathrm{d}T}{\mathrm{d}t}\right] = 0$$

- The <u>first term</u> vanishes, because

  $$\dot{\rho} + 3\frac{\dot{a}}{a}(\rho + P) = 0$$

  (FLRW continuity equation) and $V = L^3 a^3$.

- The <u>second term</u> also vanishes, because

  $$\frac{\partial P}{\partial T} = \frac{\partial S}{\partial V} = \frac{(\rho + P)}{T}$$

Thermal evolution at equilibrium:
Conservation of Entropy: two important implications
Entropy conservation has **two important consequences**:

1. It allows to identify a way to define **Number of particles**, N_i**,** of a give species.
   Let S be the total entropy and s the specific entropy of relativistic particles in equilibrium. One has (see Quiz):

$$S = sV = const. \Rightarrow s \propto a^{-3}$$

One also expects that **in equilibrium and away from mass thresholds** the (net) number of particles of a given species, i, should remain constant so the number density n_i is:

$$n_i = N_i/V \propto a^{-3}$$

Combining these expressions, one concludes that the ratio $n_i\ /\ s \propto a^0$ does not change with time **away from mass thresholds or as long as equilibrium persists**:

$$\frac{n_i}{s} = \frac{N_i/V}{S/V} = \frac{N_i}{S} = const. \qquad \Rightarrow \qquad N_i = \frac{n_i}{s} S$$

This expression allows us to define the Number of particles of the particle species i. Since $S\ =\ const.$ (and its value is only determined up to an arbitrary constant) S is usually reset to 1. In that case it is common to define the **number of particles N_i** simply as:

$$\boxed{N_i \equiv \frac{n_i}{s}}$$

Thermal evolution at equilibrium:
Conservation of Entropy: Two important implications
Entropy conservation has **two important consequences**:

2. It allows to establish a **Temperature – time** ($T\ - a(t)$) **relation.**

Considering two temperature epochs T and T_i (where "i" stands for "initial") the entropy conservation $S = s\ V = const.$ can be expressed as:

$$s(T)V = s(T_i)V_i \qquad \Leftrightarrow \qquad s(T)(L\ a)^3 = s(T_i)(L\ a_i)^3$$

where L is a comoving scale ($V = L^3 a(t)^3 = V_c\ a(t)^3$). Using the total specific entropy expression, $s = 2\pi^2\ g_{*S} T^3/45$, one has:

$$\frac{2\pi}{45} g_{*S}(T)\ T^3 a^3 = \frac{2\pi}{45} g_{*S}(T_i)\ T^3 a_i^3 \qquad \Leftrightarrow \qquad g_{*S}(T)\ T^3\ a^3 = g_{*S}(T_i)\ T_i^3 a_i^3$$

So:

$$\boxed{T = T_i \left(\frac{g_{*S}(T_i)}{g_{*S}(T)}\right)^{1/3} \frac{a_i}{a} \quad = \ A_i\ g_{*S}(T)^{-1/3}\ a^{-1}} \longrightarrow \boxed{T \propto g_{\star S}^{-1/3} a^{-1}}$$

where $A_i = T_i\ g_{*S}^{1/3}(T_i)\ a_i$. So, **away from mass thresholds** ($g_{*s} = const.$) the temperature of the relativistic fluid scales as:

$$\boxed{T \propto g_{*S}(T)^{-1/3}\ a^{-1} \propto a(t)^{-1} \propto (z+1)}$$

Thermal evolution at equilibrium:

Conservation of Entropy: Temperature – time dependence

Combining this equation in the **energy density equation** of relativistic particles one obtains:

$$\rho_r = \frac{\pi^2}{30} g_* T^4 = \frac{\pi^2}{30} g_* \left(A_i\, g_{*S}^{-1/3} a^{-1}\right)^4 = \frac{\pi^2}{30} A_i^4 \left(g_* g_{*S}^{-4/3}\right) a^{-4}$$

which is a well know result for radiation, if $g_*(T)$, $g_{*S}\,(T)$ are constants.
Plugging this result in the Friedman Equation (accounting only for relativistic particles) gives:

$$H^2 = \frac{8\pi G}{3} \rho_r = \frac{8\pi G}{3} \frac{\pi^2}{30} A_i^4 \left(g_* g_{*S}^{-4/3}\right) a^{-4}$$

These results show that, whenever $g_*(T)$ and $g_{*s}(T)$ are constants (i.e. **away from particle mass thresholds or while equilibrium persists)** one obtains:

- the well know scaling for radiation $\boldsymbol{\rho_r \propto a^{-4}}$
- the solution of the Friedman equation with $\boldsymbol{\rho = \rho_r \propto a^{-4}}$ is (Exer. Sheet 1): $\boldsymbol{a \propto t^{1/2}}$
- the temperature scaling is therefore $\boldsymbol{T \propto g_{*s} a^{-1} \propto t^{-1/2}}$

At particle mass thresholds $g_*(T)$ and $g_{*s}(T)$ are a function of temperature. The solution of the Friedmann equation is **numerical** and generally **leads to deviations** to the $\boldsymbol{a \propto t^{1/2}}$ scaling.

Thermal evolution at equilibrium:

Conservation of Entropy: Temperature – time dependence

Doing the maths, one can obtain the exact time dependence of the temperature of the relativistic fluid. Typically, one obtains:

$$\frac{T}{1\,\mathrm{MeV}} \simeq 1.5\, g_\star^{-1/4} \left(\frac{1\,\mathrm{sec}}{t}\right)^{1/2}$$

(which allows to write the rule of thumb: $T \sim 1$ MeV at about 1 second after the Big-Bang)

The temperature-time relation allows one to establish a direct **correspondence between a given energy scale of the relativistic fluid and time until the end of the radiation domination period** (see next slide).

Beyond the radiation domination phase one needs to account for the other terms accounting for non-relativists matter, curvature and dark energy in the Friedmann equation to accurately compute the age of the universe (Exer. Sheet 3…).

Thermal evolution at equilibrium:

Key events in the thermal history of the universe

Event	time t	redshift z	temperature T
Inflation	10^{-34} s (?)	–	–
Baryogenesis	?	?	?
EW phase transition	20 ps	10^{15}	100 GeV
QCD phase transition	20 μs	10^{12}	150 MeV
Dark matter freeze-out	?	?	?
Neutrino decoupling	1 s	6×10^{9}	1 MeV
Electron-positron annihilation	6 s	2×10^{9}	500 keV
Big Bang nucleosynthesis	3 min	4×10^{8}	100 keV
Matter-radiation equality	60 kyr	3400	0.75 eV
Recombination	260–380 kyr	1100–1400	0.26–0.33 eV
Photon decoupling	380 kyr	1000–1200	0.23–0.28 eV
Reionization	100–400 Myr	11–30	2.6–7.0 meV
Dark energy-matter equality	9 Gyr	0.4	0.33 meV
Present	13.8 Gyr	0	0.24 meV

The previous sets of equations allows to compute all thermodynamic properties of the primordial relativistic fluid and establish their dependence with time and redshifts.

All one needs to know is what matter/energy components exist in the universe and the physics of each of these components!

Thermal evolution at equilibrium:

Key events in the thermal history of the universe

Baryogenesis (& Leptogenesis):

Quantum field theory requires the existence of anti-particles. This poses a problem: *particle-antiparticle creation and annihilation (allowed by the Heisenberg principle) creates/destroys* ***equal amounts of particle and anti-particles***. However, we do observe an excess of matter (mostly baryons and leptons) over anti-matter (that can be produced, e.g., in accelerators)! Models of ***baryogenesis*** attempt to describe this observational evidence using some ***dynamical mechanism*** (instead of assuming this particle-anti-particle asymmetry *ab initio*).

Electroweak phase transition:

At $\sim 100\ GeV$ particles acquire mass through the Higgs mechanism. This leads to a drastic change of the weak interaction. The gauge bosons Z^0, $W^{\pm}$ become massive and soon after decouple from thermal equilibrium.

QCD phase transition:

Above $\sim 150\ MeV$ quarks are asymptotically free (i.e. weakly interacting). Below this energy/mass threshold the strong force (mediated by the gluons) becomes more intense; the more massive quarks start to decouple from the fluid. The less massive become confined (with the gluons) inside the baryons (3 quarks + gluons) and mesons (quarks + anti-quark + gluons).

Dark Matter freeze-out:

Present observations indicates that dark matter is **very-weakly** or **non-interacting**. Depending on the mass of the dark matter candidates one should expect that they should decouple from the fluid early on. For example, if dark matter is made of WIMPs (weakly interactive massive particles), one should expect their abundance should freeze around $1\ MeV$.

Thermal evolution at equilibrium:

Key events in the thermal history of the universe

Neutrino decoupling:
Neutrinos only interact with the rest of the plasma through the weak force. They are expected to decouple from the fluid at $\sim 0.8\ MeV$.

Electron-positron annihilation:
Electrons and positron annihilate soon after the neutrinos. Positrons vanish, because electron-positron pair production is strongly suppressed below $\sim 1 MeV$

Big Bang Nucleosynthesis:
At $\sim 0.1 MeV$ (~ 3 minutes after the Big-Bang) protons and neutrons combine to form the first light nuclear elements.

Recombination:
At $\sim 0.3\ eV$ $(260 - 380$ kyr) free electrons combine with nuclei to form atoms. Predominantly Hydrogen: $e^- + p^+ \rightarrow H + \gamma$. Below this range of energies, this chemical reaction can no longer occur in the reverse order.

Photon CMB decoupling:
By $\sim 0.23\ eV$ (380 kyr) the primordial fluid is reduced to photons, that no longer interact with matter (free electrons). The Cosmic Microwave Background radiation propagates freely in the Universe.

Thermal evolution at equilibrium:

Brief history of the Universe