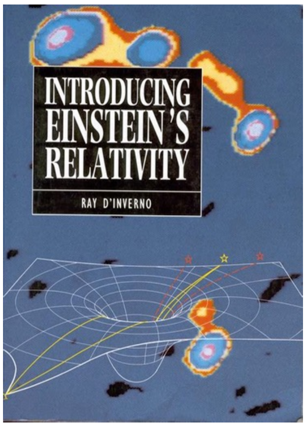
Universo Primordial

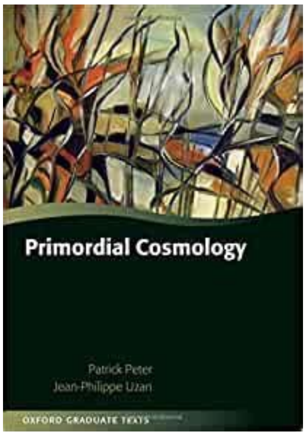
2024-2025 (1º Semestre)

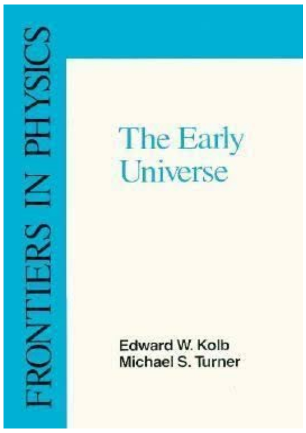
Mestrado em Física - Astronomia

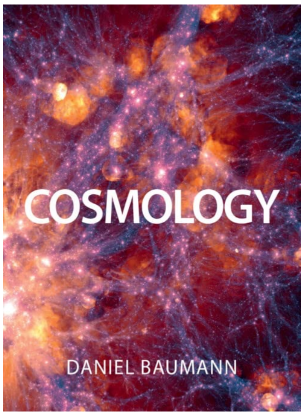
Chapter 2

2. The Standard Model of Cosmology (SMC)
   - Fundamental assumptions;
   - The GR equations and the Friedmann-Lemaitre-Robertson-Walker (FLRW) solution;
   - FLRW models:
     - Dynamic equations;
     - Energy-momentum conservation;
     - Fluid components and equations of state;
     - Cosmological parameters;
     - The Friedmann equation: the evolutionary phases of the Universe; exact solutions: age of the Universe;
     - Distances; horizons and volumes;
     - The accelerated expansion of the Universe;
   - Problems with the SMC: Horizon; Flatness; Relic particles; origin of perturbations; primordial Isotropy and homogeneity
   - The idea of Inflation

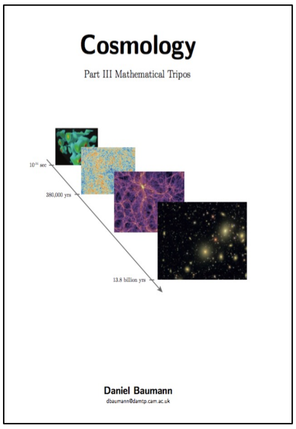




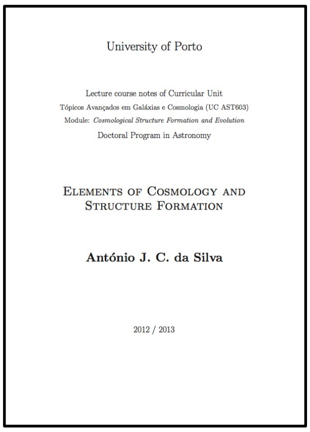




Ch. 2



Ch. 1



Ch. 1

Standard Model of Cosmology

SMC: Mathematical framework

Fundamental assumptions:

- The Universe is homogeneous and isotropic when observed on large scales and expands uniformly with respect to any position

- The dynamics of space-time is described by Einstein's theory of general relativity (GR).

$$G_{ab} = R_{ab} - \frac{1}{2} R g_{ab} = \frac{8\pi G}{c^4} T_{ab}$$

for the Universe to be homogeneous and isotropic the stress-energy tensor must be that of a perfect fluid

$$T_{ab} = (\rho + \frac{p}{c^2}) U_a U_b - \frac{p}{c^2} g_{ab}$$

SMC: Mathematical framework

The cosmological constant in the GR equation:

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu},$$ (Λ as "cosmological constant")

$$G_{\mu\nu} = 8\pi G \left(T_{\mu\nu} - \frac{\Lambda}{8\pi G} g_{\mu\nu} \right) = 8\pi G \tilde{T}_{\mu\nu}$$ (Λ as "vacuum energy")

The Einstein tensor, Ricci tensor and Ricci scalar are:

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R$$

$$R_{\mu\nu} = \Gamma^{\alpha}_{\mu\nu,\alpha} - \Gamma^{\alpha}_{\mu\alpha,\nu} + \Gamma^{\alpha}_{\mu\nu}\Gamma^{\beta}_{\alpha\beta} - \Gamma^{\alpha}_{\mu\beta}\Gamma^{\beta}_{\alpha\nu}$$

$$R = g^{\mu\nu} R_{\mu\nu}$$

$$\Gamma^{\mu}_{\nu\lambda} = \frac{1}{2} g^{\mu\alpha}(g_{\lambda\alpha,\nu} + g_{\nu\alpha,\lambda} - g_{\nu\lambda,\alpha}) \qquad g_{\mu\nu,\lambda} \equiv \frac{\partial g_{\mu\nu}}{\partial x^\lambda} = \partial_\lambda\, g_{\mu\nu} \qquad g^{\mu\lambda} g_{\lambda\nu} = \delta^{\mu}_{\nu}$$

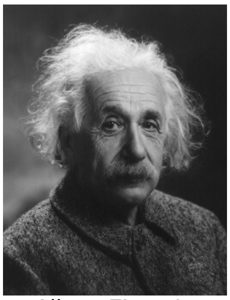
Metric tensor gives the line element:

$$\mathrm{d}s^2 = \sum_{\mu,\nu=0}^{3} g_{\mu\nu} \mathrm{d}X^\mu \mathrm{d}X^\nu \equiv g_{\mu\nu} \mathrm{d}X^\mu \mathrm{d}X^\nu$$

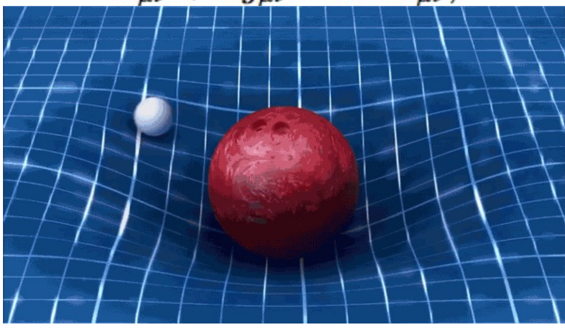
SMC: Mathematical framework

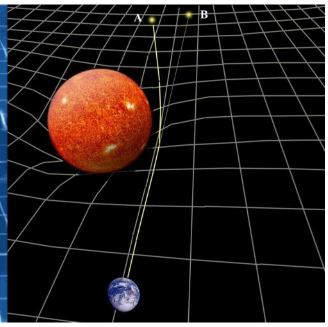
Einstein Equation:

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu},$$

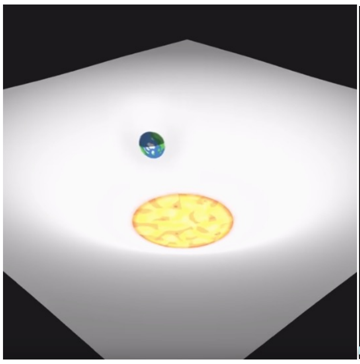


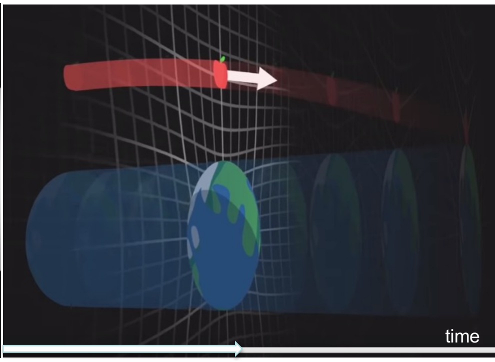
Albert Einstein
1879-1955

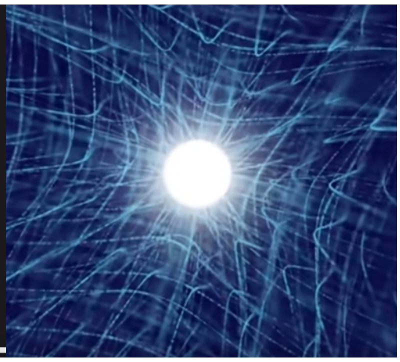




More accurate ways of visualizing GR: https://youtu.be/wrwgIjBUYVc







SMC: Mathematical framework

Geodesic Equation:

In the absence of non-gravitational forces, free falling particles move along "geodesics", described by the so-called Geodesic equation.

$$\boxed{\frac{dU^\mu}{ds} + \Gamma^\mu_{\alpha\beta} U^\alpha U^\beta = 0}$$

where,

$$U^\mu \equiv \frac{dX^\mu}{ds}$$

four-velocity of the particle along its free-falling path $X^\mu(s)$

$$S = -m \int_A^B \mathrm{d}s \; .$$

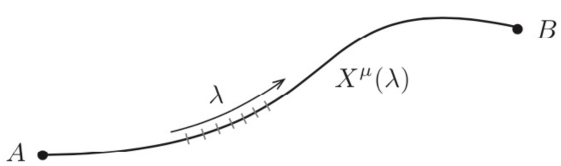


Figure 1.4: *Parameterisation of an arbitrary path in spacetime, $X^\mu(\lambda)$.*

SMC: Mathematical framework

Fundamental assumptions:

- The Universe is homogeneous and isotropic when observed on large scales and expands uniformly with respect to any position

- The dynamics of space-time is described by Einstein's theory of general relativity (GR).

$$G_{ab} = R_{ab} - \frac{1}{2} R g_{ab} = \frac{8\pi G}{c^4} T_{ab} - \Lambda g_{ab} \qquad T_{ab} = (\rho + \frac{p}{c^2}) U_a U_b - \frac{p}{c^2} g_{ab}$$

In these conditions **the solution of the Einstein equation** is the Friedmann-Lemaitre-Robertson-Walker (**FLRW**) metric:

$$ds^2 = c^2 dt^2 - a^2(t) \left[\frac{dr^2}{1 - kr^2} + r^2 \left(d\theta^2 + \sin^2\theta d\phi^2 \right) \right]$$

SMC: Mathematical framework

- Dynamical equations:
  (result from the Einstein equations and govern the time evolution of $a(t)$)

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho + \frac{\Lambda c^2}{3} - \frac{kc^2}{a^2}$$

Friedmann equation

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\left(\rho + 3\frac{p}{c^2}\right) + \frac{\Lambda c^2}{3}$$

Raychaudhuri (or acceleration) equation

- Energy momentum conservation: $\nabla_\mu T^\mu{}_\nu \equiv T^\mu{}_{\nu;\mu} = 0$

the covariant derivative reads: $\nabla_\mu T^\mu{}_\nu = \partial_\mu T^\mu{}_\nu + \Gamma^\mu_{\mu\lambda} T^\lambda{}_\nu - \Gamma^\lambda_{\mu\nu} T^\mu{}_\lambda = 0$

the $\nu = 0$ (time) component of this equation gives:

$$\dot{\rho} = -3\frac{\dot{a}}{a}\left(\rho + \frac{p}{c^2}\right) \quad \Rightarrow \quad d\left(\rho c^2 a^3\right) = -p\, d\left(a^3\right)$$

Energy conservation equation

$$p = w\rho c^2 \qquad -1 \le w \le 1$$

Equation of State (EoS)

for fluids with constant EoS parameter, w, the solution is:

$$\rho(t) = \rho_i \left(\frac{a(t)}{a_i}\right)^{-3(1+w)}$$

SMC: Mathematical framework

Covariant derivative:

Covariant derivative.—The covariant derivative is an important object in differential geometry and it is of fundamental importance in general relativity. The geometrical meaning of ∇_μ will be discussed in detail in the GR course. In this course, we will have to be satisfied with treating it as an operator that acts in a specific way on scalars, vectors and tensors:

- There is no difference between the covariant derivative and the partial derivative if it acts on a scalar

$$\nabla_\mu f = \partial_\mu f \ . \tag{1.3.83}$$

- Acting on a contravariant vector, V^ν, the covariant derivative is a partial derivative plus a correction that is linear in the vector:

$$\nabla_\mu V^\nu = \partial_\mu V^\nu + \Gamma^\nu_{\mu\lambda} V^\lambda \ . \tag{1.3.84}$$

  Look carefully at the index structure of the second term. A similar definition applies to the covariant derivative of covariant vectors, ω_ν,

$$\nabla_\mu \omega_\nu = \partial_\mu \omega_\nu - \Gamma^\lambda_{\mu\nu} \omega_\lambda \ . \tag{1.3.85}$$

  Notice the change of the sign of the second term and the placement of the dummy index.

- For tensors with many indices, you just repeat (1.3.84) and (1.3.85) for each index. For each upper index you introduce a term with a single $+\Gamma$, and for each lower index a term with a single $-\Gamma$:

$$\begin{aligned}\nabla_\sigma T^{\mu_1\mu_2\cdots\mu_k}{}_{\nu_1\nu_2\cdots\nu_l} = {} & \partial_\sigma T^{\mu_1\mu_2\cdots\mu_k}{}_{\nu_1\nu_2\cdots\nu_l} \\ & + \Gamma^{\mu_1}{}_{\sigma\lambda} T^{\lambda\mu_2\cdots\mu_k}{}_{\nu_1\nu_2\cdots\nu_l} + \Gamma^{\mu_2}{}_{\sigma\lambda} T^{\mu_1\lambda\cdots\mu_k}{}_{\nu_1\nu_2\cdots\nu_l} + \cdots \\ & - \Gamma^{\lambda}{}_{\sigma\nu_1} T^{\mu_1\mu_2\cdots\mu_k}{}_{\lambda\nu_2\cdots\nu_l} - \Gamma^{\lambda}{}_{\sigma\nu_2} T^{\mu_1\mu_2\cdots\mu_k}{}_{\nu_1\lambda\cdots\nu_l} - \cdots \ .\end{aligned} \tag{1.3.86}$$

  This is the general expression for the covariant derivative. Luckily, we will only be dealing with relatively simple tensors, so this monsterous expression will usually reduce to something managable.

SMC: Mathematical framework

- EoS for different energy density components:

$$\rho(t) = \rho_i \left(\frac{a(t)}{a_i}\right)^{-3(1+w)}$$

$$\left(\frac{\dot a}{a}\right)^2 = \frac{8\pi G}{3}\rho$$

  - w = 1/3 (radiation)

$$\rho_\gamma = \rho_{\gamma 0}\left(\frac{a_0}{a}\right)^4 \xrightarrow{(1)} \left(\frac{\dot a}{a}\right)^2 \propto \frac{1}{a^4} \longrightarrow a \propto t^{1/2}.$$

  - w = 0 (collisionless matter)

$$\rho_{\rm m} = \rho_{\rm m0}\left(\frac{a_0}{a}\right)^3 \xrightarrow{(2)} \left(\frac{\dot a}{a}\right)^2 \propto \frac{1}{a^3} \longrightarrow a \propto t^{2/3}.$$

  - w = -1 (cosmological constant)

$$\rho_\Lambda = \Lambda/8\pi G = -P_\Lambda \xrightarrow{(3)} a \propto e^{\sqrt{\Lambda/3}\,t}$$

(1) after integration of the Friedmann equation with $k = 0,\ \Lambda = 0,\ \rho = \rho_\gamma$.
(2) after integration of the Friedmann equation with $k = 0,\ \Lambda = 0,\ \rho = \rho_m$.
(3) after integration of the Friedmann equation with $k = 0,\ \Lambda = 8\pi G\rho_\Lambda,\ \rho = 0$

SMC: FLRW models

- Cosmological parameters:

The Friedmann equation is:

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho + \frac{\Lambda c^2}{3} - \frac{kc^2}{a^2}$$

Specifying each density component in $\rho = \rho_r + \rho_B + \rho_{DM}$, gives

$$H^2 = \frac{8\pi G}{3}(\rho_r + \rho_B + \rho_{DM}) + \frac{\Lambda c^2}{3} - \frac{kc^2}{a^2}$$

where,

$$H^2 = \frac{\dot{a}(t)}{a(t)}, \qquad \text{and} \qquad \rho = \rho_r + \rho_B + \rho_{DM}$$

Dark Matter energy density

Baryon energy density

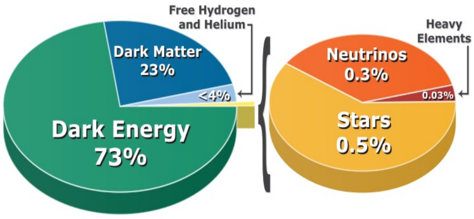
Radiation energy density

dividing by H^2 the Friedmann equation one gets

$$1 = \frac{8\pi G}{3H^2}\rho_r + \frac{8\pi G}{3H^2}\rho_B + \frac{8\pi G}{3H^2}\rho_{DM} + \frac{\Lambda c^2}{3H^2} - \frac{k c^2}{a^2 H^2}$$

or

$$1 = \Omega_r + \Omega_B + \Omega_{DM} + \Omega_\Lambda + \Omega_k$$

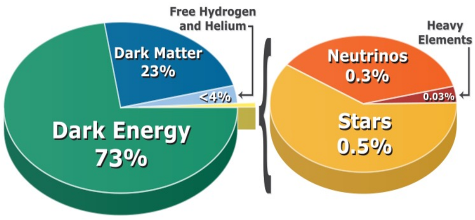


SMC: FLRW models

- Cosmological parameters:

$$1 = \frac{8\pi G}{3H^2}\rho_r + \frac{8\pi G}{3H^2}\rho_B + \frac{8\pi G}{3H^2}\rho_{DM} + \frac{\Lambda c^2}{3H^2} - \frac{k c^2}{a^2 H^2}$$

$$1 = \Omega_r + \Omega_B + \Omega_{DM} + \Omega_\Lambda + \Omega_k$$



The quantity:

$$\rho_{crit} = \frac{3H^2}{8\pi G}$$

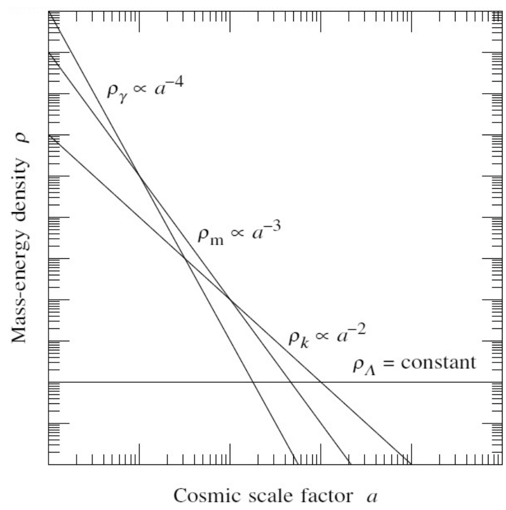
Is known as **critical density**. It can be used to express each energy density parameter as a ratio of densities:

$$\Omega_i \equiv \frac{\rho_i}{\rho_{\rm crit}}$$

In particular, the curvature and Λ terms give:

$$\Omega_k = \frac{\rho_k}{\rho_{\rm crit}} = -\frac{kc^2}{a^2H^2} \iff \rho_k = -\frac{3\,kc^2}{8\pi G\,a^2}$$

$$\Omega_\Lambda = \frac{\rho_\Lambda}{\rho_{\rm crit}} = \frac{\Lambda\, c^2}{3\,H^2} \iff \rho_\Lambda = \frac{\Lambda c^2}{8\pi G}$$



SMC: FLRW models

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho + \frac{\Lambda c^2}{3} - \frac{kc^2}{a^2}$$

- Cosmological parameters:

$$\frac{8\pi G}{3H^2}\rho + \frac{\Lambda c^2}{3H^2} - \frac{kc^2}{a^2H^2} = 1 \quad \Leftrightarrow \quad \sum_i \Omega_i + \Omega_\Lambda + \Omega_k = 1$$

$$H(t) = \frac{\dot{a}(t)}{a(t)}$$

Hubble parameter

$$\Omega_i \equiv \frac{\rho_i}{\rho_{\text{crit}}} \qquad \Omega_\Lambda = \frac{\Lambda c^2}{3H^2}, \qquad \Omega_k = -\frac{kc^2}{a^2H^2}, \qquad \rho_c = \frac{3H^2}{8\pi G}$$

Critical energy density

Mater-energy density parameters: $(\Omega = \frac{\rho}{\rho_c} = \sum \frac{\rho_i}{\rho_c} = \sum \Omega_i)$

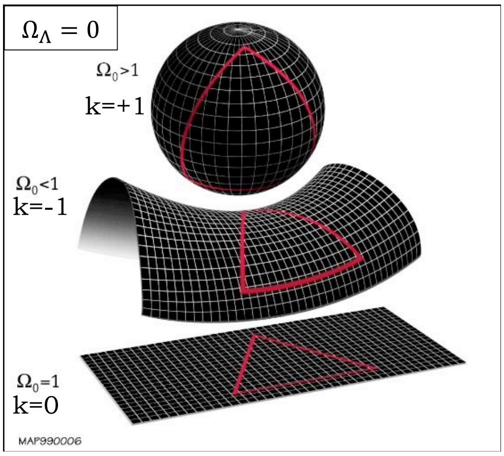
Vacuum or dark energy density parameter

Curvature density parameter

$$\rho \equiv \sum_i \rho_i$$

$$\Omega_i \equiv \frac{\rho_i}{\rho_{\text{crit}}}$$

includes all matter and radiation components (baryons, dark mater, radiation, …)

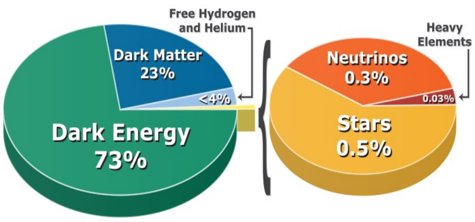


SMC: FLRW models

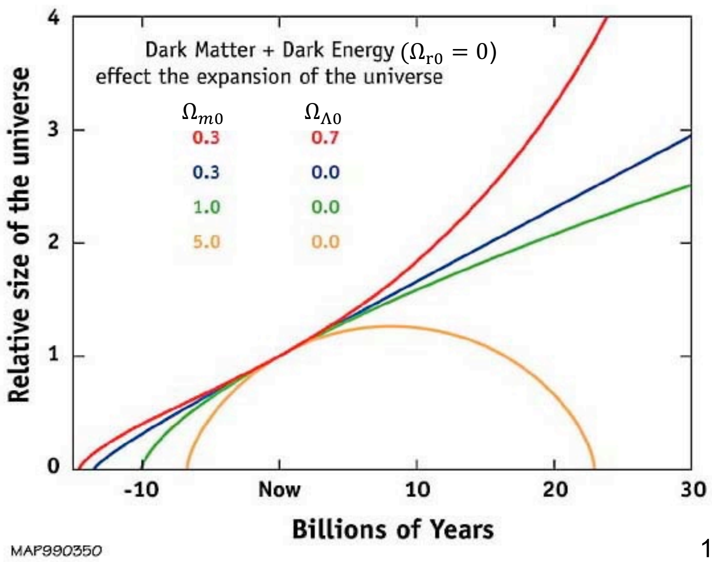
$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho + \frac{\Lambda c^2}{3} - \frac{kc^2}{a^2}$$

- Friedmann equation revisited

$$H^2(t) = \frac{8\pi G}{3}(\rho_r + \rho_m) - \frac{kc^2}{a^2} + \frac{\Lambda c^2}{3}$$
$$= H_0^2\left[\Omega_{r0}\left(\frac{a_0}{a}\right)^4 + \Omega_{m0}\left(\frac{a_0}{a}\right)^3 + \Omega_{k0}\left(\frac{a_0}{a}\right)^2 + \Omega_{\Lambda 0}\right]$$



The evolutionary fate of the Universe is determined by cosmological parameters



SMC: Exact solutions of the Friedmann equation

- Scale factor: $$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho + \frac{\Lambda c^2}{3} - \frac{kc^2}{a^2}$$

$$\frac{d}{dt}\frac{a(t)}{a_0} = H_0\sqrt{1-\Omega_0+\Omega_{m0}\left(\frac{a}{a_0}\right)^{-1}+\Omega_{r0}\left(\frac{a}{a_0}\right)^{-2}-\Omega_{\Lambda 0}\left[1-\left(\frac{a}{a_0}\right)^2\right]}$$

with $\Omega_0 = 1 - \Omega_{\Lambda 0} - \Omega_{k0}$. For a critical density ($\Omega_k = \Omega_\Lambda = 0$) universe, gives:

$$\frac{a(t)}{a_0} = \left(\frac{3(1+w)}{2}H_0 t\right)^{2/(3(1+w))}$$
$$H(t) = \frac{\dot{a}}{a} = \frac{2}{3(w+1)t}$$

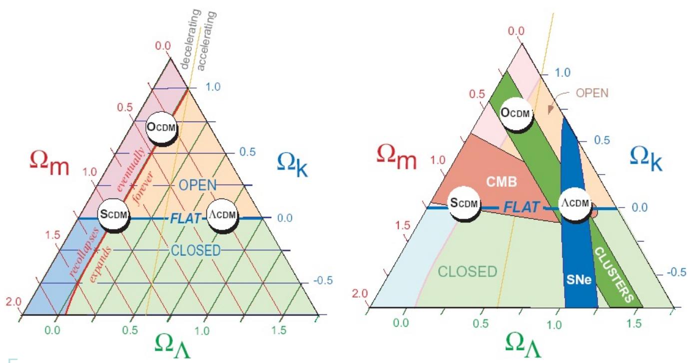
Redshift:
$$z = \frac{E-E_0}{E_0} = \frac{\nu}{\nu_0} - 1 = \frac{\lambda_0}{\lambda} - 1 = \frac{a_0}{a} - 1$$

- Age of the Universe:

$$t = H_0^{-1}\int_0^{\frac{a(t)}{a_0}=(1+z)^{-1}} \frac{1}{\sqrt{1-\Omega_0+\Omega_{m0}x^{-1}+\Omega_{r0}x^{-2}-\Omega_\Lambda(1-x^2)}}dx$$

SMC: Concordance Cosmology

Combination of different observational datasets…

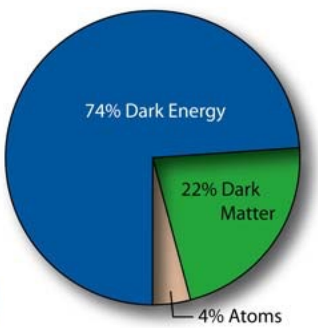


From:

WMAP3 parameters

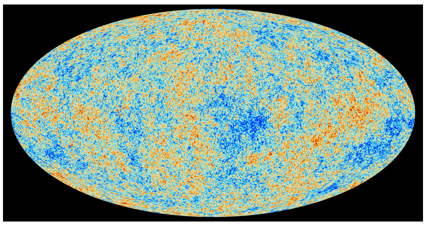
Parameter	Value	Description
	Basic parameters	
H_0	$70.9^{+2.4}_{-3.2}$ km s^{-1} Mpc^{-1}	Hubble parameter
Ω_b	$0.0444^{+0.0042}_{-0.0035}$	Baryon density
Ω_m	$0.266^{+0.025}_{-0.040}$	Total matter density (baryons + dark matter)
τ	$0.079^{+0.029}_{-0.032}$	Optical depth to reionization
A_s	$0.813^{+0.042}_{-0.052}$	Scalar fluctuation amplitude
n_s	$0.948^{+0.015}_{-0.018}$	Scalar spectral index
	Derived parameters	
ρ_0	$0.94^{+0.06}_{-0.09} \times 10^{-26}$ kg/m^3	Critical density
Ω_Λ	$0.732^{+0.040}_{-0.025}$	Dark energy density
z_{ion}	$10.5^{+2.6}_{-2.9}$	Reionization red-shift
σ_8	$0.772^{+0.036}_{-0.048}$	Galaxy fluctuation amplitude
t_0	$13.73^{+0.13}_{-0.17} \times 10^9$ years	Age of the universe

… allow us to impose constraints on cosmological parameters



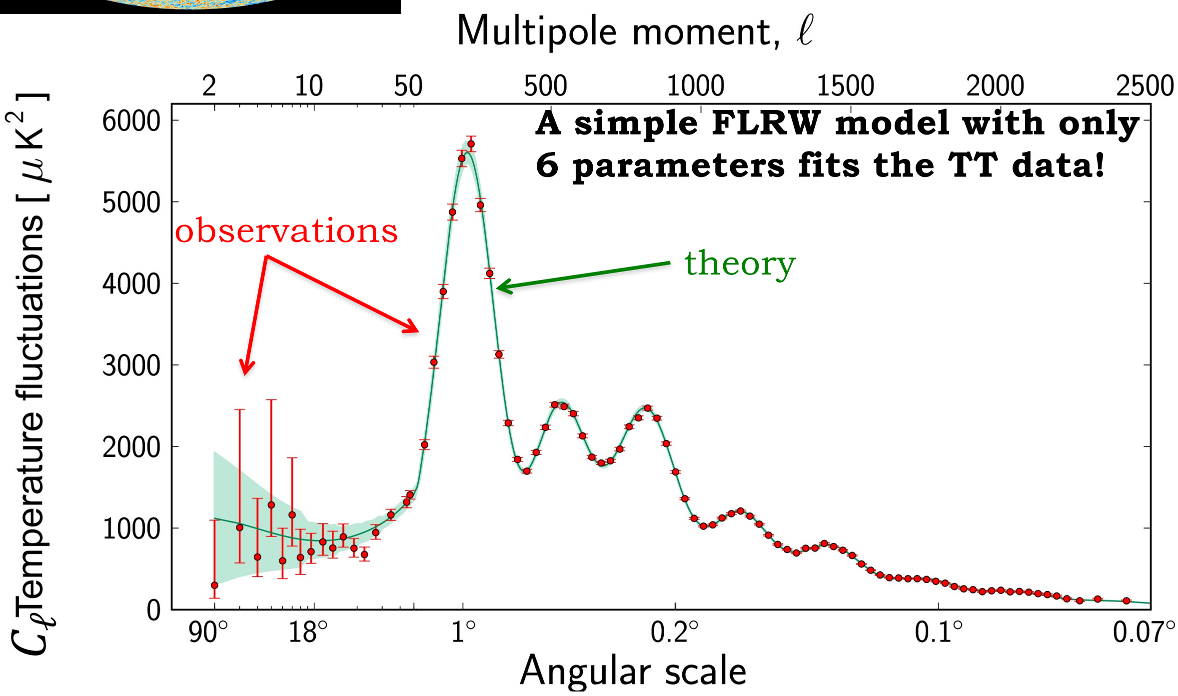
$$\sum_i \Omega_i + \Omega_\Lambda + \Omega_k = 1$$

SMC: Cosmological parameters after Planck



What's C_ℓ in this plot?

Short Answer: The $C_\ell = \ell(\ell+1)C_\ell^{TT}/2\pi$, where $C_\ell^{TT} = \sum_l |a_{lm}^2|$.
The $a_{\ell m}$ are the *projection coefficients* of the temperature fluctuation function $\Delta T/T_0\ (\theta,\phi)$ *over a basis of functions of the celestial sphere: the spherical harmonics* $Y_{\ell m}$. See Appendix I.



SMC: Cosmological parameters after Planck

From: Planck collaboration. XVI. arXiv:1303.5076

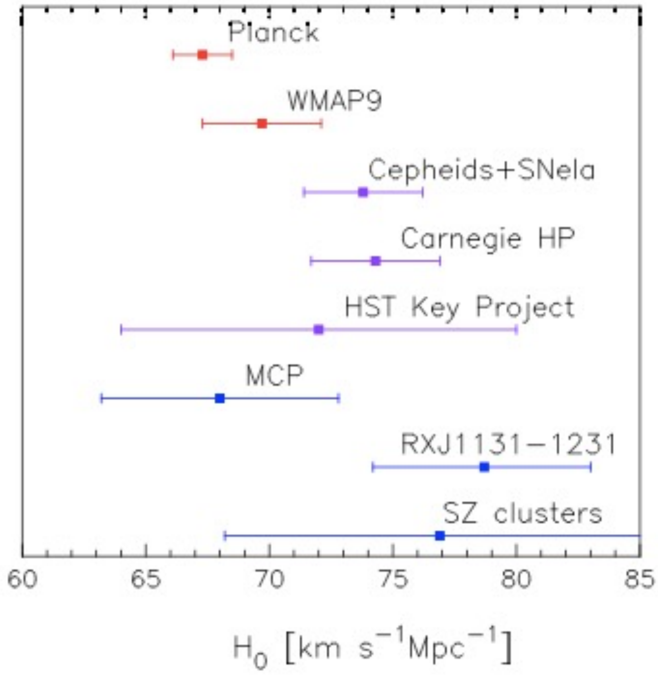
Table 2. Cosmological parameter values for the six-parameter base ΛCDM model. Columns 2 and 3 give results for the *Planck* temperature power spectrum data alone. Columns 4 and 5 combine the *Planck* temperature data with *Planck* lensing, and columns 6 and 7 include *WMAP* polarization at low multipoles. We give best fit parameters (i.e. the parameters that maximise the overall likelihood for each data combination) as well as 68% confidence limits for constrained parameters. The first six parameters have flat priors. The remainder are derived parameters as discussed in Sect. 2. Beam, calibration parameters, and foreground parameters (see Sect. 4) are not listed for brevity. Constraints on foreground parameters for *Planck*+WP are given later in Table 5.

	Planck		*Planck*+lensing		*Planck*+WP	
Parameter	Best fit	68% limits	Best fit	68% limits	Best fit	68% limits
$\Omega_b h^2$	0.022068	0.02207 ± 0.00033	0.022242	0.02217 ± 0.00033	0.022032	0.02205 ± 0.00028
$\Omega_c h^2$	0.12029	0.1196 ± 0.0031	0.11805	0.1186 ± 0.0031	0.12038	0.1199 ± 0.0027
$100\theta_{MC}$	1.04122	1.04132 ± 0.00068	1.04150	1.04141 ± 0.00067	1.04119	1.04131 ± 0.00063
τ	0.0925	0.097 ± 0.038	0.0949	0.089 ± 0.032	0.0925	$0.089^{+0.012}_{-0.014}$
n_s	0.9624	0.9616 ± 0.0094	0.9675	0.9635 ± 0.0094	0.9619	0.9603 ± 0.0073
$\ln(10^{10}A_s)$	3.098	3.103 ± 0.072	3.098	3.085 ± 0.057	3.0980	$3.089^{+0.024}_{-0.027}$
Ω_Λ	0.6825	0.686 ± 0.020	0.6964	0.693 ± 0.019	0.6817	$0.685^{+0.018}_{-0.016}$
Ω_m	0.3175	0.314 ± 0.020	0.3036	0.307 ± 0.019	0.3183	$0.315^{+0.016}_{-0.018}$
σ_8	0.8344	0.834 ± 0.027	0.8285	0.823 ± 0.018	0.8347	0.829 ± 0.012
z_{re}	11.35	$11.4^{+4.0}_{-2.8}$	11.45	$10.8^{+3.1}_{-2.5}$	11.37	11.1 ± 1.1
H_0	67.11	67.4 ± 1.4	68.14	67.9 ± 1.5	67.04	67.3 ± 1.2
$10^9 A_s$	2.215	2.23 ± 0.16	2.215	$2.19^{+0.12}_{-0.14}$	2.215	$2.196^{+0.051}_{-0.060}$
$\Omega_m h^2$	0.14300	0.1423 ± 0.0029	0.14094	0.1414 ± 0.0029	0.14305	0.1426 ± 0.0025
$\Omega_m h^3$	0.09597	0.09590 ± 0.00059	0.09603	0.09593 ± 0.00058	0.09591	0.09589 ± 0.00057
Y_P	0.247710	0.24771 ± 0.00014	0.247785	0.24775 ± 0.00014	0.247695	0.24770 ± 0.00012
Age/Gyr	13.819	13.813 ± 0.058	13.784	13.796 ± 0.058	13.8242	13.817 ± 0.048
z_*	1090.43	1090.37 ± 0.65	1090.01	1090.16 ± 0.65	1090.48	1090.43 ± 0.54
r_*	144.58	144.75 ± 0.66	145.02	144.96 ± 0.66	144.58	144.71 ± 0.60
$100\theta_*$	1.04139	1.04148 ± 0.00066	1.04164	1.04156 ± 0.00066	1.04136	1.04147 ± 0.00062
z_{drag}	1059.32	1059.29 ± 0.65	1059.59	1059.43 ± 0.64	1059.25	1059.25 ± 0.58
r_{drag}	147.34	147.53 ± 0.64	147.74	147.70 ± 0.63	147.36	147.49 ± 0.59
k_D	0.14026	0.14007 ± 0.00064	0.13998	0.13996 ± 0.00062	0.14022	0.14009 ± 0.00063
$100\theta_D$	0.161332	0.16137 ± 0.00037	0.161196	0.16129 ± 0.00036	0.161375	0.16140 ± 0.00034
z_{eq}	3402	3386 ± 69	3352	3362 ± 69	3403	3391 ± 60
$100\theta_{eq}$	0.8128	0.816 ± 0.013	0.8224	0.821 ± 0.013	0.8125	0.815 ± 0.011
$r_{drag}/D_V(0.57)$	0.07130	0.0716 ± 0.0011	0.07207	0.0719 ± 0.0011	0.07126	0.07147 ± 0.00091

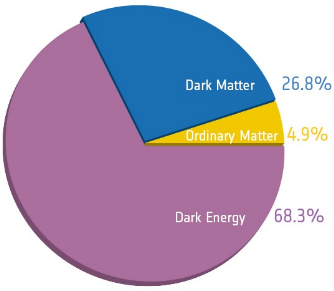
SMC: Cosmological parameters after Planck

Constraints after Planck ("background" parameters): $\Omega_{\Lambda 0}, \Omega_{B0}, \Omega_{DM0}, \Omega_{r0}, \Omega_{k0}, H_0$

From: Planck collaboration. XVI. arXiv:1303.5076



Unprecedented precision from CMB observations!

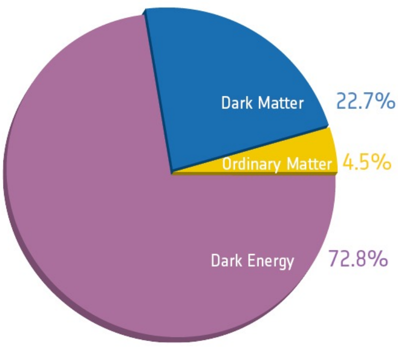


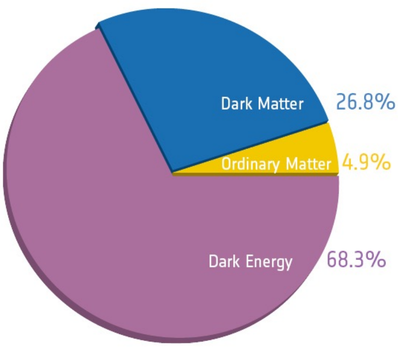
After Planck
(2013)

SMC: Cosmological parameters after Planck

Constraints before and after Planck ("background" parameters): $\Omega_{\Lambda 0}, \Omega_{B0}, \Omega_{DM0}, \Omega_{r0}, \Omega_{k0}$

$$\Omega_{r0} + \Omega_{B0} + \Omega_{DM0} + \Omega_{\Lambda 0} + \Omega_{k0} = 1$$





Before Planck

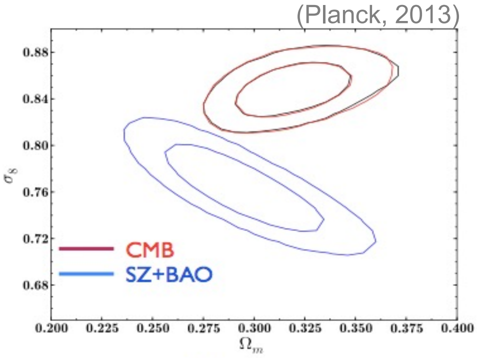
After Planck

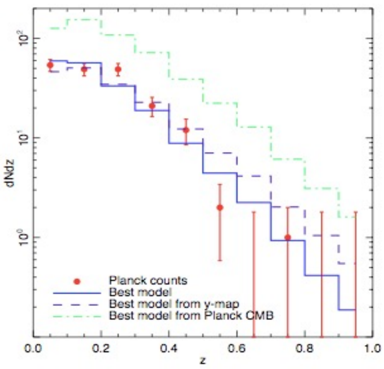
Fig. credits: ESA / PLANCK collaboration

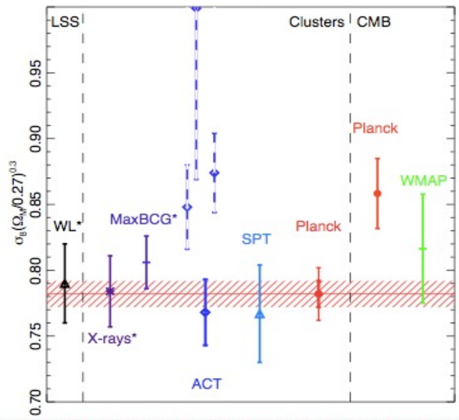
SMC: Limitations of a 6-parameter model…

Constraints after Planck ("perturbations" related parameters): σ_8

- Higher values of Ω_m, σ_8 in *Planck* CMB analysis
- 3σ tension
- More general tension between clusters and CMB ?







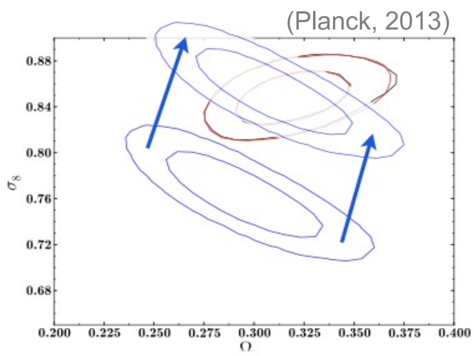
IAS

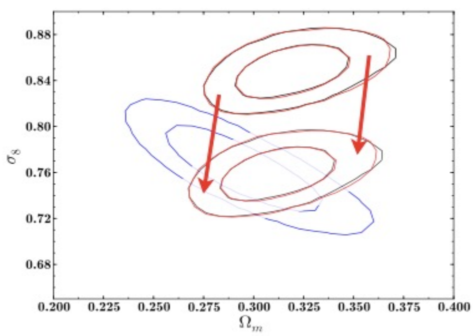
M. Douspis, 03/04/2013, Cosmology from Planck SZ cluster counts

SMC: Limitations of a 6 parameter model…

Constraints after Planck ("perturbations" related parameters): σ_8

- Getting higher σ_8 from clusters
  - Change scaling
  - Change bias
  - Account for missing clusters
- Getting lower σ_8 from CMB
  - Change initial power spectrum
  - Change transfert function





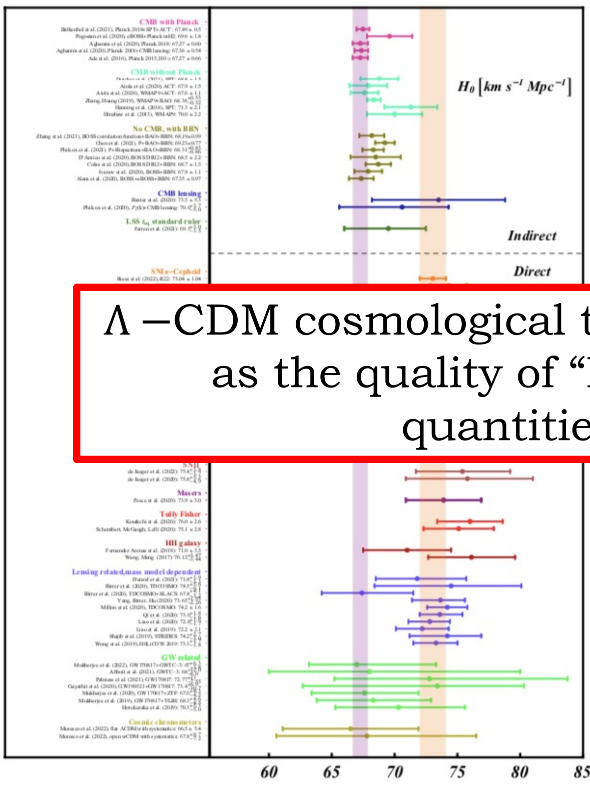
The "era" of Cosmological tensions starts!!!

SMC: Limitations of a 6-parameter model?...

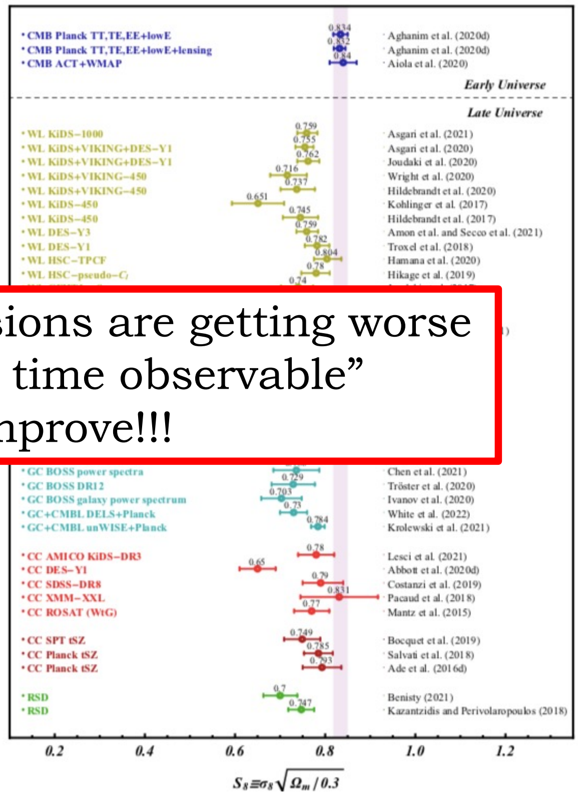
Review paper (2021): Cosmology Intertwined: A Review of the Particle Physics, Astrophysics, and Cosmology Associated with the Cosmological Tensions and Anomalies, Abdala et al. , Snowmass 2021.

https://arxiv.org/abs/2203.06142

The H_0 tension



The σ_8 tension



Λ −CDM cosmological tensions are getting worse as the quality of "late time observable" quantities improve!!!

Planck Legacy: A new baseline cosmological model for the Euclid era

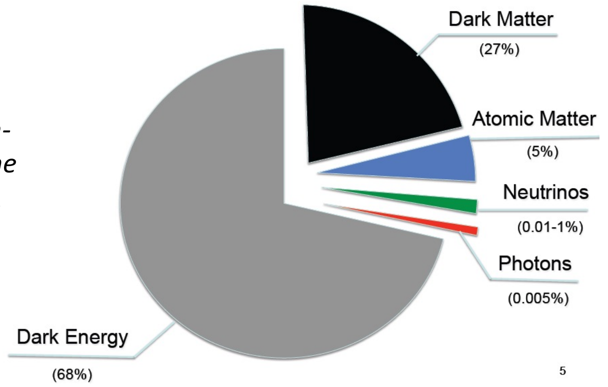
The (new) concordance model: ΛCDM + massive neutrinos

From: Planck collaboration. XIII (2015)

Parameter	TT	TT+lensing	TT+lensing+ext	TT, TE, EE	TT, TE, EE+lensing	TT, TE, EE+lensing+ext
Ω_K	$-0.052^{+0.049}_{-0.055}$	$-0.005^{+0.016}_{-0.017}$	$-0.0001^{+0.0054}_{-0.0052}$	$-0.040^{+0.038}_{-0.041}$	$-0.004^{+0.015}_{-0.015}$	$0.0008^{+0.0040}_{-0.0039}$
Σm_ν [eV]	< 0.715	< 0.675	< 0.234	< 0.492	< 0.589	< 0.194
$N_{\rm eff}$	$3.13^{+0.64}_{-0.63}$	$3.13^{+0.62}_{-0.61}$	$3.15^{+0.41}_{-0.40}$	$2.99^{+0.41}_{-0.39}$	$2.94^{+0.38}_{-0.38}$	$3.04^{+0.33}_{-0.33}$
Y_P	$0.252^{+0.041}_{-0.042}$	$0.251^{+0.040}_{-0.039}$	$0.251^{+0.035}_{-0.036}$	$0.250^{+0.026}_{-0.027}$	$0.247^{+0.026}_{-0.027}$	$0.249^{+0.025}_{-0.026}$
$dn_s/d\ln k$	$-0.008^{+0.016}_{-0.016}$	$-0.003^{+0.015}_{-0.015}$	$-0.003^{+0.015}_{-0.014}$	$-0.006^{+0.014}_{-0.014}$	$-0.002^{+0.013}_{-0.013}$	$-0.002^{+0.013}_{-0.013}$
$r_{0.002}$	< 0.103	< 0.114	< 0.114	< 0.0987	< 0.112	< 0.113
w	$-1.54^{+0.62}_{-0.50}$	$-1.41^{+0.64}_{-0.56}$	$-1.006^{+0.085}_{-0.091}$	$-1.55^{+0.58}_{-0.48}$	$-1.42^{+0.62}_{-0.56}$	$-1.019^{+0.075}_{-0.080}$

$$\sum m_\nu = 0.16^{+0.08}_{-0.11}\,\text{eV} \quad (\textit{Planck}\ \text{TT+lowP+aggressive lensing + BAO; 68\%})$$

If neutrinos are considered to have non-negligible mass, then the σ_8 tension is reduced…



Horizons, Distances and Volumes

SMC: Particle and Event horizons

Consider light travelling along radial ($d\theta = d\phi = 0$) geodesics in a FLRW metric ($c = 1$):

$$ds^2 = dt^2 - a^2(t)\left[\frac{dr^2}{1-kr^2} + r^2(d\theta^2 + \sin^2\theta d\phi^2)\right],$$
$$= dt^2 - a^2(t)\left[d\chi^2 + f_k(\chi)(d\theta^2 + \sin^2\theta d\phi^2)\right],$$

($d\chi = dr$ for flat geometries, see e.g. Sec. 2.1.3 Baumann). Let's set $d\theta = d\phi = 0$ and define **conformal time** as $d\tau = dt/a$. This allows us to write:

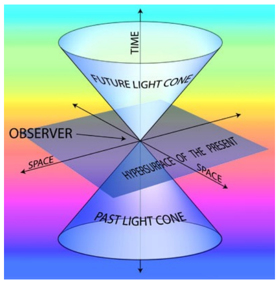
$$\mathrm{d}s^2 = a^2(\tau)\left[\mathrm{d}\tau^2 - \mathrm{d}\chi^2\right]$$

Since **light rays travel along null** ($ds^2 = 0$) **geodesics:** $d\chi = \pm d\tau$

Integrating from the **past** (t_i) to **present** (t) or from **the present to the future** (t_f) one can define:

- Particle horizon: $\chi_{\mathrm{ph}}(\tau) = \tau - \tau_i = \int_{t_i}^{t} \frac{\mathrm{d}t}{a(t)}$ with $t_i = 0$
- Event horizon: $\chi_{\mathrm{eh}}(\tau) = \tau_f - \tau = \int_{t}^{t_f} \frac{\mathrm{d}t}{a(t)}$ with $t_f = \infty$

SMC: Particle and Event horizons



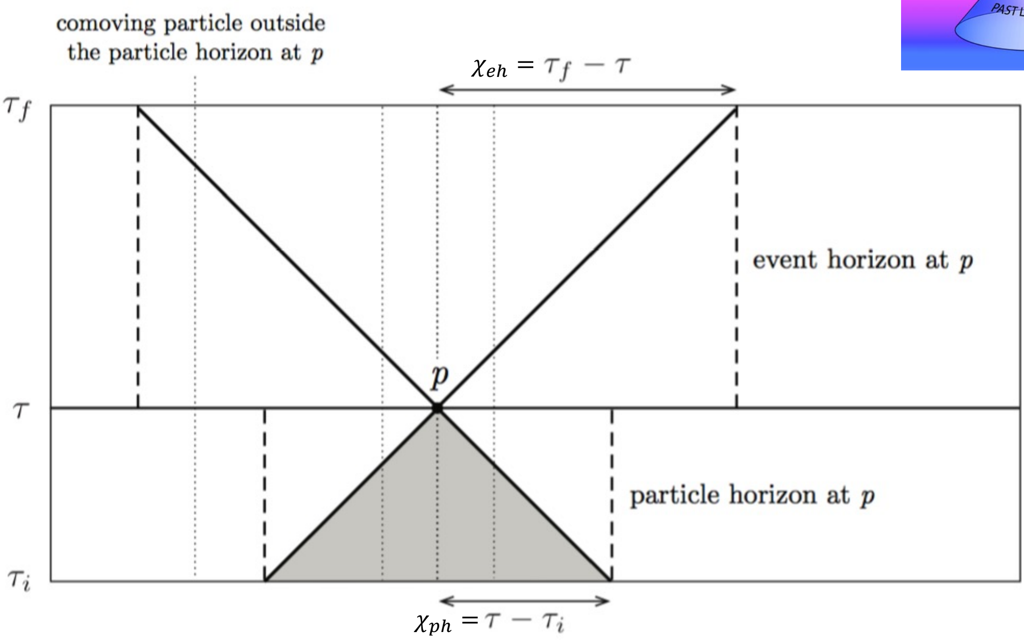


Figure 2.1: *Spacetime diagram illustrating the concept of horizons. Dotted lines show the worldlines of comoving objects. The event horizon is the maximal distance to which we can send signal. The particle horizon is the maximal distance from which we can receive signals.*

SMC: distances, angular sizes and volumes

Comoving distance (line of sight):

(*also computed using photons that travel along null geodesics*, $ds^2 = 0$, *with* $d\theta = d\phi = 0$)

$$ds^2 = c^2dt^2 - a(t)^2 \frac{dr^2}{1-kr^2} \longrightarrow \chi = \int_r^{r_0} \frac{dr'}{\sqrt{1-kr'^2}} = c\int_t^{t_0} \frac{dt'}{a} = \frac{c}{a_0}\int_0^z \frac{dz}{H(z)}$$

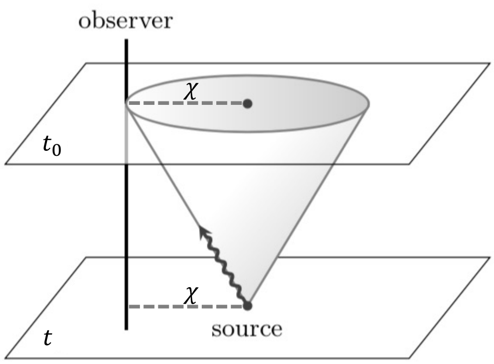
The last equality arises from the fact that ($z+1 = a_0/a$):

$$\frac{dz}{dt} = -\frac{a_0}{a^2}\frac{da}{dt} = -\frac{a_0}{a}\frac{\dot{a}}{a} = -(z+1)H \longrightarrow \frac{dt}{dz} = -\frac{1}{(z+1)H}$$

So:

$$\chi = c\int_t^{t_0} \frac{dt}{a} = \int_t^{t_0} \left(\frac{z+1}{a_0}\right) \frac{dt}{dz} dz = \frac{c}{a_0}\int_z^0 \frac{-1}{H(z)} dz \Leftrightarrow$$

$$\chi = \frac{c}{a_0}\int_0^z \frac{dz}{H(z)} = \frac{c}{a_0 H_0}\int_0^z \frac{dz}{E(z)}$$



where $H(z) = H_0 E(z)$. Form the Friedmann Eq. (slide 16) one has:

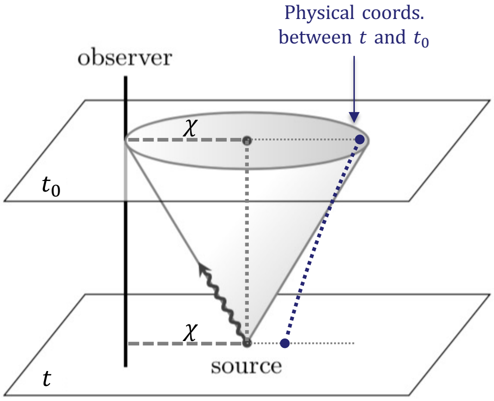
$$E(z) = \sqrt{\Omega_{r0}(z+1)^4 + \Omega_{m0}(z+1)^3 + \Omega_{k0}(z+1)^2 + \Omega_\Lambda}$$

SMC: distances, angular sizes and volumes

Physical (proper) distance (line of sight):

Is the **physical distance to the object at a given cosmic time *t*.** It is simply obtained from the comoving coordinate distance, by multiplying by $a(t)$:

$$d(t) = a(t)\chi = a(t)\int_r^{r_0} \frac{dr'}{\sqrt{1-kr'^2}} = a(t)\, c\int_t^{t_0} \frac{dt'}{a(t')} =$$
$$= \frac{a(t)}{a_0}\frac{c}{H_0}\int_0^z \frac{dz}{E(z)} = \frac{1}{z+1}\frac{c}{H_0}\int_0^z \frac{dz}{E(z)}$$



Note that for $t \rightarrow t_0$ one has: $r \rightarrow r_0$, and $z \rightarrow 0$. So:

$$d(t_0) = a_0\int_{r_0}^{r_0} \frac{dr}{\sqrt{1-kr^2}} = = \frac{c}{H_0}\int_0^0 \frac{dz}{E(z)} = 0$$

This is not the physical coordinates of the source $z = 0$.

Hubble radius and Hubble length:

Are defined as the length scale obtained when one sets $v_H = c$ in the Hubble law:

Hubble radius: $d_H = \frac{c}{H_0}$ **Hubble length**: $R_H(t) = \frac{c}{H(t)}$

SMC: distances, angular sizes and volumes

Luminosity distance:

It is defined by the relation between the source Luminosity, L, (energy emitted by the source per unit of time) at t_1 and the Flux, F, (energy per unit of time per area) arriving at t_0:

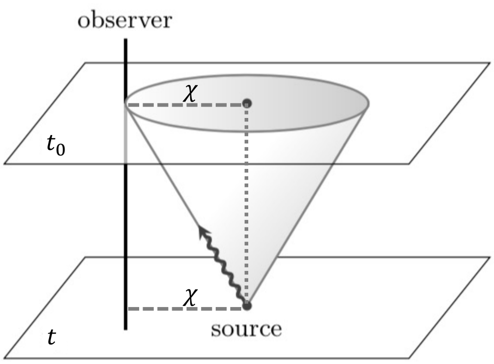
$$d_L^2 = \frac{L}{4\pi\, F(z+1)^2} \Leftrightarrow F = \frac{L}{4\pi\,(z+1)^2 d_L^2}$$

The $(z+1)^2$ appears because between emission and reception the energy of photons decreases by $a(t_1)/a_0 = z+1$ and the rate of photons arriving at t_0 also decreases by $a(t_1)/a_0 = z+1$. It can be shown (e.g. D'Inverno) that:

$$d_L = (z+1)\, d_M$$

Where $\boldsymbol{d_M}$ is **Metric distance** or **Transverse comoving distance** given by ($d_H = c/H_0$):

$$d_M = d_H \begin{cases} \Omega_k^{-1/2}\sinh\left(\Omega_k^{1/2}\int_0^z \frac{dz'}{E(z')}\right) & (\Omega_k > 0) \\ \int_0^z \frac{dz'}{E(z')} & (\Omega_k = 0) \\ |\Omega_k|^{-1/2}\sinh\left(|\Omega_k|^{1/2}\int_0^z \frac{dz'}{E(z')}\right) & (\Omega_k < 0) \end{cases}$$



So, for $\Omega_k = 0$: $d_L = (z+1)\, a_0\chi = (z+1)\frac{c}{H_0}\int_0^z \frac{dz'}{E(z')}$

SMC: distances, angular sizes and volumes

Luminosity distance:

It is defined by the relation between the source Luminosity, L, (energy emitted by the source per unit of time) at t_1 and the Flux, F, (energy per unit of time per area) arriving at t_0

$$d_L = (z+1)\, d_M$$

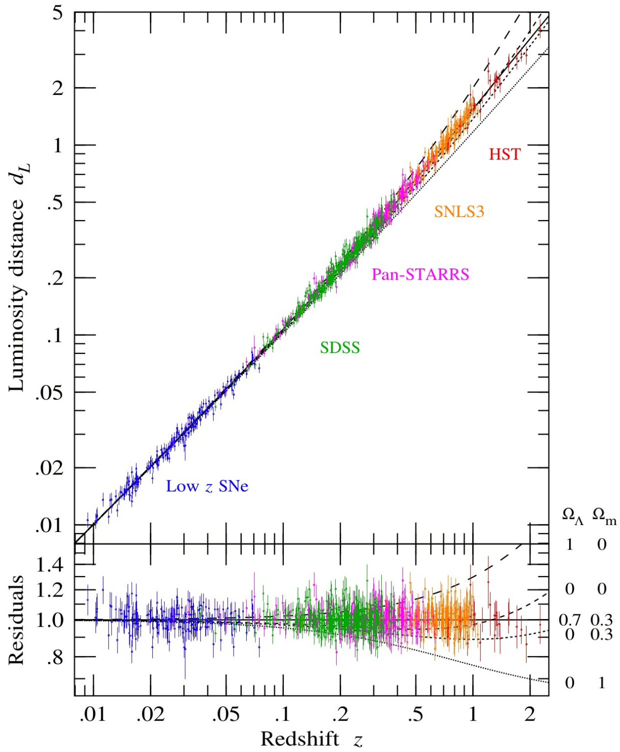
with

$$d_M = d_H \begin{cases} \Omega_k^{-1/2} \sinh\left(\Omega_k^{1/2} \int_0^z \frac{dz'}{E(z')}\right) & (\Omega_k > 0) \\ \int_0^z \frac{dz'}{E(z')} & (\Omega_k = 0) \\ |\Omega_k|^{-1/2} \sinh\left(|\Omega_k|^{1/2} \int_0^z \frac{dz'}{E(z')}\right) & (\Omega_k < 0) \end{cases}$$

For $\Omega_k = 0$:

$$d_L = (z+1)\, a_0 \chi = (z+1) \frac{c}{H_0} \int_0^z \frac{dz'}{E(z')}$$

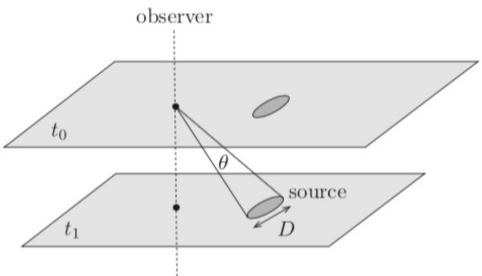
$d_L = d_L(z, H_0, \Omega_r, \Omega_M, \Omega_\Lambda, \ldots)$ is a powerful probe of cosmological parameters



SMC: distances angular sizes and volumes

Angular diameter distance:

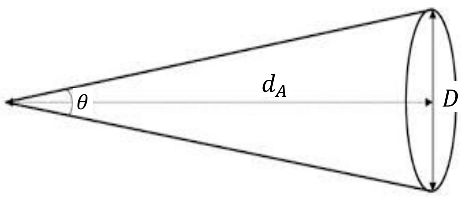
The angular size of an object at a given time is given by the ratio of 2 physical (proper) distances:



$$\tan\frac{\theta}{2} = \frac{D/2}{d_A} \iff \theta \simeq \frac{D}{d_A}$$

where D is the physical size of the object and

$$d_A(z) = \frac{d_M}{z+1}$$



is the physical proper distance to the object. For $\Omega_k = 0$

$$d_A = a(t) \int_r^{r_0} \frac{dr'}{\sqrt{1-kr'^2}} = a(t) \int_t^{t_0} \frac{dt'}{a(t')} =$$

$$= \frac{a(t)}{a_0} \frac{c}{H_0} \int_0^z \frac{dz}{E(z)} = \frac{1}{z+1} \frac{c}{H_0} \int_0^z \frac{dz}{E(z)}$$

For matter dominated models with $\Omega_\Lambda = 0$ one can obtain simple analytical expressions for $\theta(z)$, which becomes also a function of the Universe's cosmological parameters:

$$\theta(z) \simeq \frac{[D/(2cH_0^{-1})]\,(1+z)^2 \Omega_{m0}^2}{\Omega_{m0} z + (\Omega_{m_0} - 2)\left(\sqrt{1+\Omega_{m0} z} - 1\right)}$$

SMC: distances angular sizes and volumes

Angular diameter distance:

Comparing the expressions of $d_A = d_M/(z+1)$ and $d_L = (z+1)d_M$ one concludes that:

$$d_A(z) = \frac{d_M(z)}{z+1} = \frac{d_L(z)}{(z+1)^2} \quad \Longleftrightarrow \quad \boxed{d_L(z) = (z+1)^2\, d_A(z)}$$

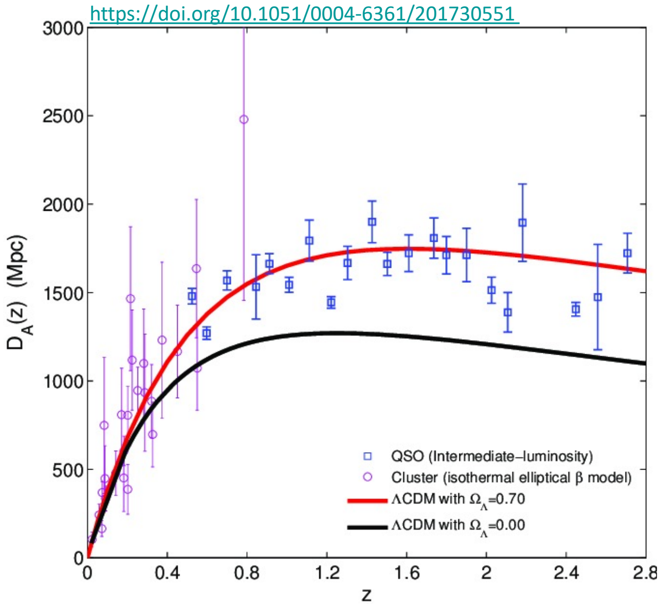
This equality is known as the **Etherington's reciprocity theorem**.

Computing d_A for matter dominated models with $\Omega_\Lambda = 0$, the angular size of a collapsed object of physical size D is:

$$\theta(z) \simeq \frac{[D/(2cH_0^{-1})]\,(1+z)^2\Omega_{m0}^2}{\Omega_{m0}z + (\Omega_{m_0} - 2)\left(\sqrt{1+\Omega_{m0}z} - 1\right)}$$

But, collapsed objects (e.g galaxies) should not be affected by cosmic expansion so, if the physics that intervenes in their formation is also independent of redshift, their **physical sizes are the same at all redshifts**. Measuring their angular sizes can therefore be used to compute their d_A distances from us and constrain cosmological parameters!

https://doi.org/10.1051/0004-6361/201730551



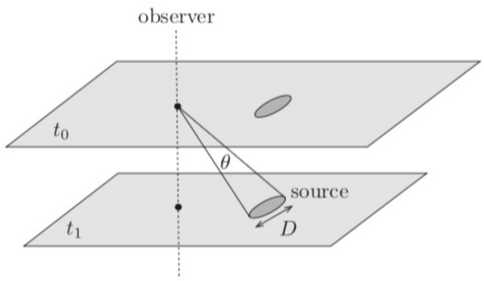
SMC: distances angular sizes and volumes

- Angular size of an object at a given time is given by the ratio of 2 proper (physical) distances:

$$\boxed{\tan\frac{\theta}{2} = \frac{D/2}{d_A}} \Longleftrightarrow \theta \simeq \frac{D}{d_A}$$

For matter dominated models with $\Omega_\Lambda = 0$

$$\theta(z) \simeq \frac{[D/(2cH_0^{-1})]\,(1+z)^2\Omega_{m0}^2}{\Omega_{m0}z + (\Omega_{m_0} - 2)\left(\sqrt{1+\Omega_{m0}z} - 1\right)}$$



- **An interesting question:** is to ask what is the angular size of the **particle horizon's light cone** at a given time/redshift, $\theta_H(z)$. To compute this one must compute $\boldsymbol{\chi_{ph}(z)}$.

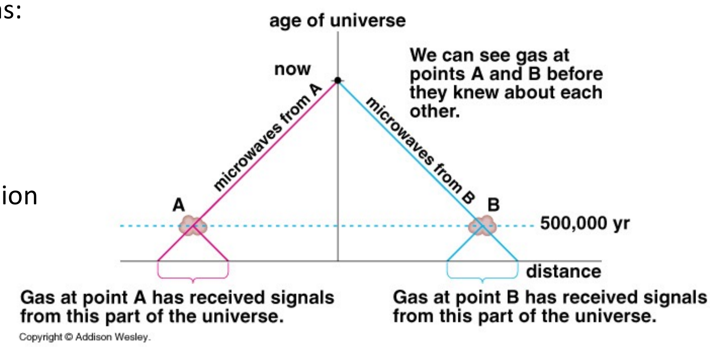
At CMB decoupling (z_{CMB}~1100) the 2 main fluid components are **matter and radiation**. One can compute $\theta_H(z_{CMB})$ (see Baumann Sec 4.1.1) as:

$$\boxed{\theta_H(z_{CMB}) = \frac{2\,\eta_{CMB}}{\eta_0 - \eta_{CMB}} = 0.036\,rad \approx 2^0}$$

Note: η_0 and η_{CMB} are conformal times at present and z_{CMB}. Their values (after integration $\eta = \int da/a(t)$) are:

$$\eta_0 \simeq \frac{2\,cH_0^{-1}}{\sqrt{\Omega_{m0}}} \qquad \eta_{CMB} = 0.0175\,\eta_0$$

From:



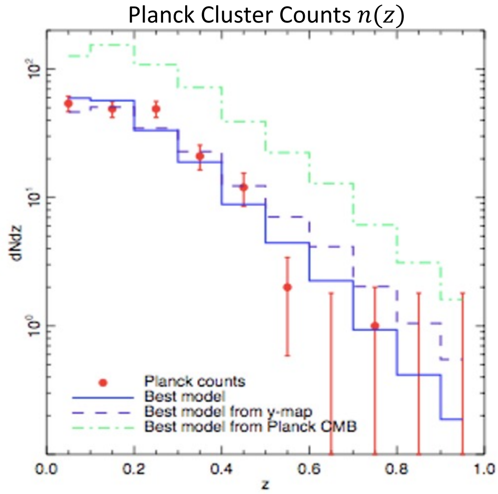
SMC: distances, angular sizes and volumes

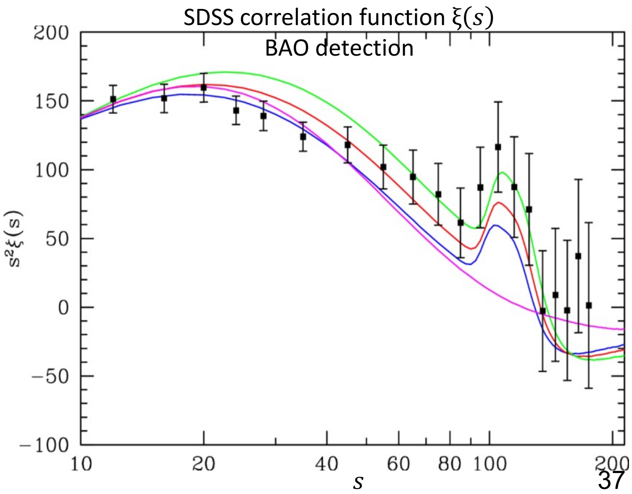
Physical volume element:

It is defined in the usual way "$dV = dx\, dy\, dz$". In spherical coordinates is $(d\Omega = d\theta\, d\phi)$:

$$dV = \sqrt{|g|}\, dr d\theta d\phi \quad \longrightarrow \quad \frac{dV}{d\Omega dz} = \frac{c\, d_A^2}{H(z)(z+1)} = \frac{c}{H_0} \frac{d_A^2}{E(z)(z+1)}$$

It intervenes in any observable that requires averaging over space, such as quantities that involve **volume densities**. For example, the number density of galaxies or galaxy clusters, $n(z) = N(z)/V(z)$, or the evaluation galaxy correlation functions, $\xi(s) = \langle \rho(r)\rho(r+s) \rangle$





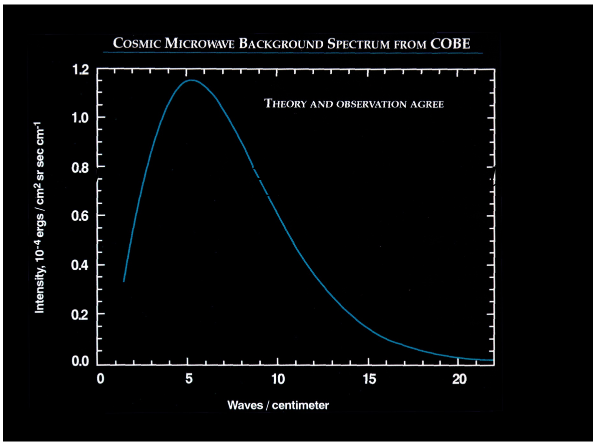
"Foudational" problems
of the FLRW models as the sole
ingredient of the SMC

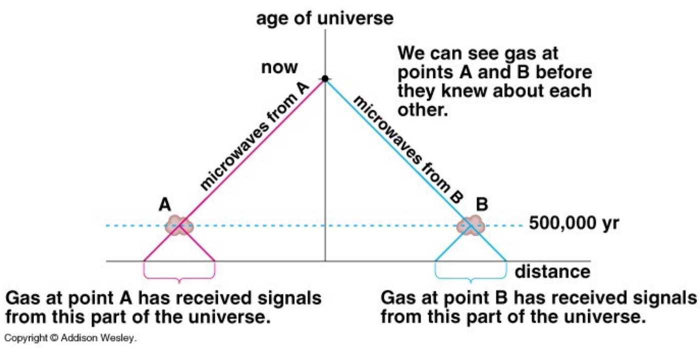
The Horizon Problem

As we have seen the angular size of the particle horizon at CMB decoupling z_{cmb}~1100 is just:

$$\theta_H(z_{CMB}) = \frac{2\,\eta_{CMB}}{\eta_0 - \eta_{CMB}} = 0.036\,rad \approx 2^0$$

So, there are about ~27000 **causal disconnected** angular areas in the CMB sky. So, *why the CMB has a thermal spectrum with a so uniform temperature in all directions (2.725 °K)?*





The Flatness Problem

From the Friedmann Equation, with Λ=0, one has

$$|\Omega(t) - 1| = \frac{|k|}{a^2(t)H^2(t)} = \frac{|k|}{\dot{a}^2(t)}$$

is a decreasing function of time: So as t → 0 , Ω → 1

decreases tremendously as time approaches the big bang instant.

This means that as we go back in time the **energy density of universe must be extremely close to the critical density** (t → 0 ⇒ Ω → 1). For t=1e-43 s (Planck time) Ω should deviate no more than 1e-60 from the unity!

Why the universe has to "start" with Ω(t) so close to 1?

The Monopoles & other relics Problem

Particle physics predicts that a variety of **"exotic" stable particles**, such as the magnetic monopoles, should be produced in the early phase of the Universe and remain in measurable amounts until the present.

No such particles have yet been observed. Why?

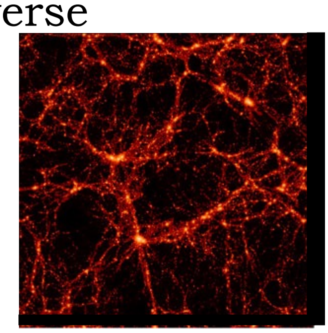
This either implies that the predictions from particle physics are wrong, or their densities are very small and therefore there is something missing from this evolutionary picture of the Big Bang.

The Origin of Perturbations Problem

Locally the universe is not homogeneous. It displays a complex hierarchical pattern of galaxies, clusters and super clusters.

What's the origin of cosmological structure?
Does it grow from gravitational instability?
What is the origin of the initial perturbations?

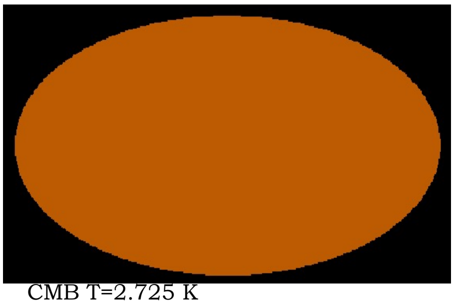
Without a mechanism to explain their existence one must assume that they "were born" with the universe already showing the correct amplitudes on all scales, so that gravity can accurately reproduce the present-day structures?

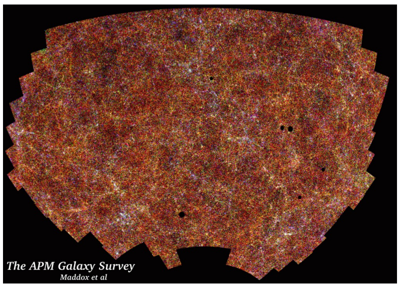


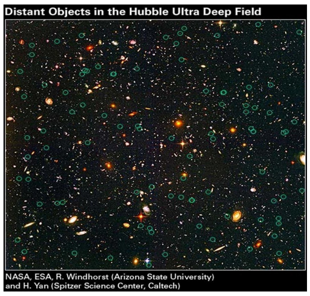
The homogeneity and isotropy Problem

Why is the universe homogeneous on large scales? **At early times homogeneity had to be even more "perfect"**.

The **FLRW** universes form a **very special subset of solutions** of the GR equations. So, *why nature "prefers" homogeneity and isotropy from the beginning as opposed to having evolved into that stage?*


CMB T=2.725 K





Theory of Inflation: solves the problems?

Inflation can be defined as

$$\text{Inflation} \Leftrightarrow \ddot{a} > 0 \Leftrightarrow \frac{d}{dt}\left(cH^{-1}/a\right) < 0.$$

This happens when

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\left(\rho + 3\frac{p}{c^2}\right)$$

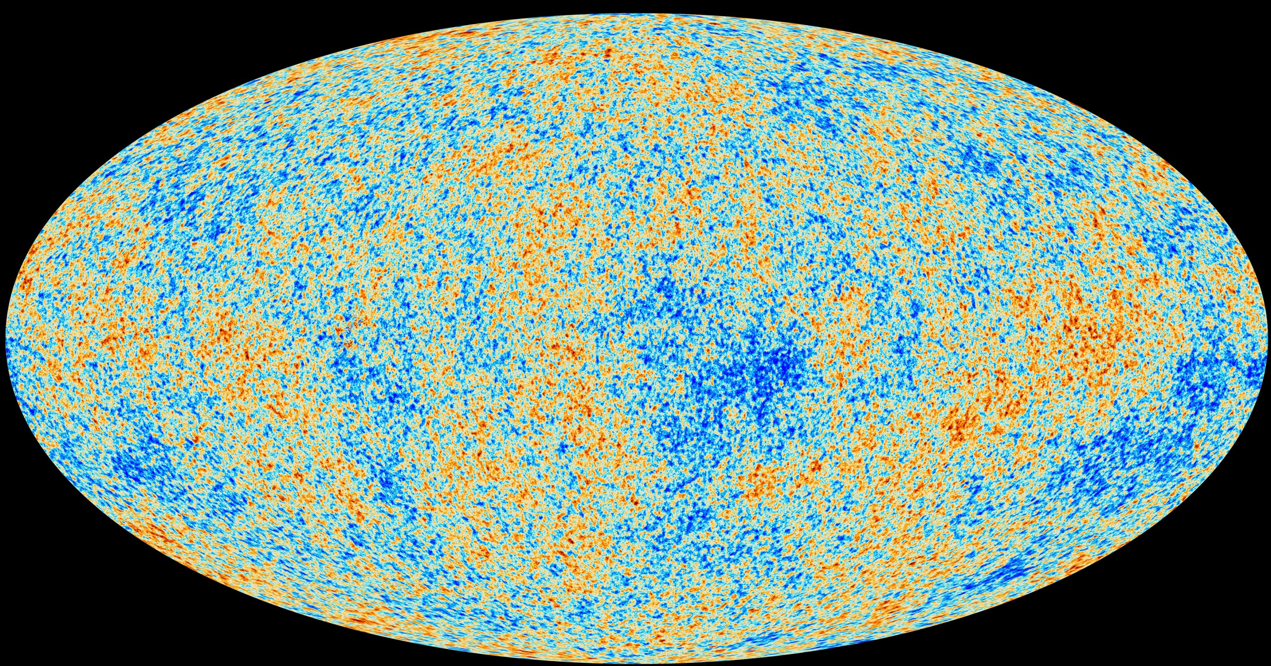
$$\Rightarrow \quad \ddot{a} > 0 \;\Leftrightarrow\; \rho + \frac{3p}{c^2} < 0 \Leftrightarrow$$

$$p < -\rho c^2/3$$

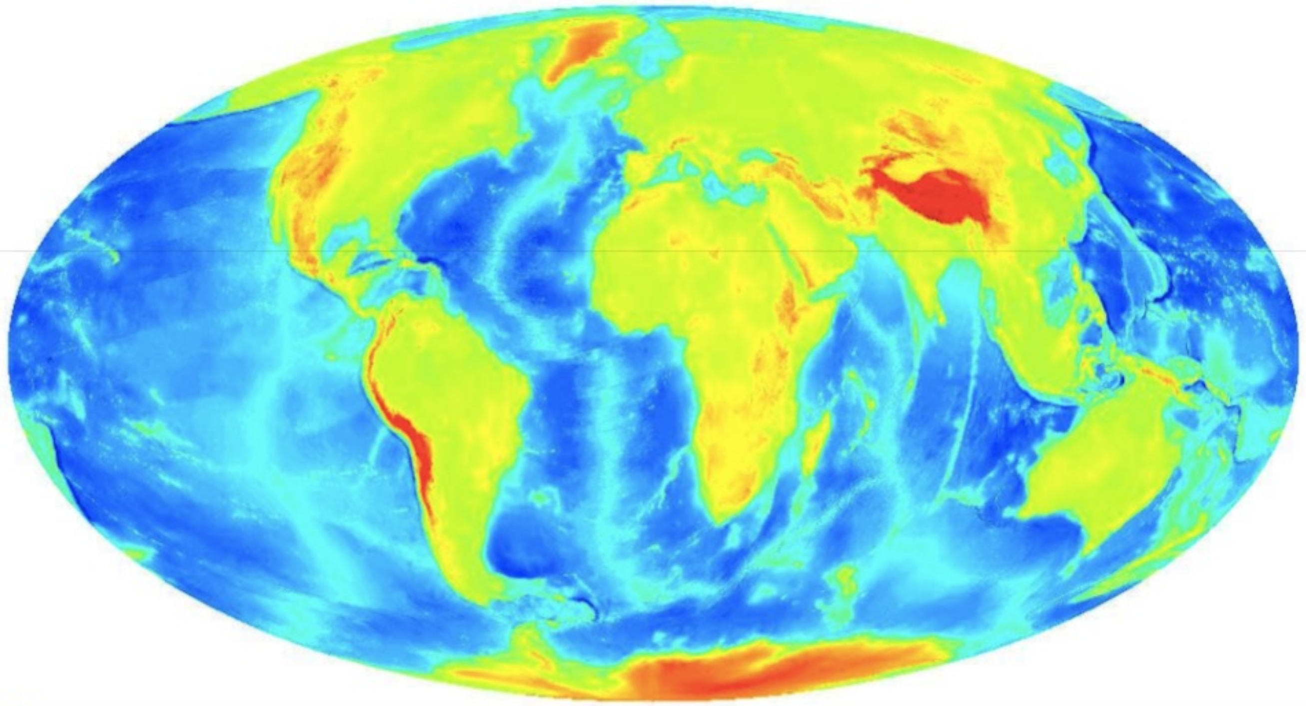
Riddle: no known matter / energy component has an equation of state parameter $w = \rho c^2/p <$ -1/3... (continues in Chapter 8)

Appendix 1
CMB angular power spectrum: a primer

Why an “elliptical” map?

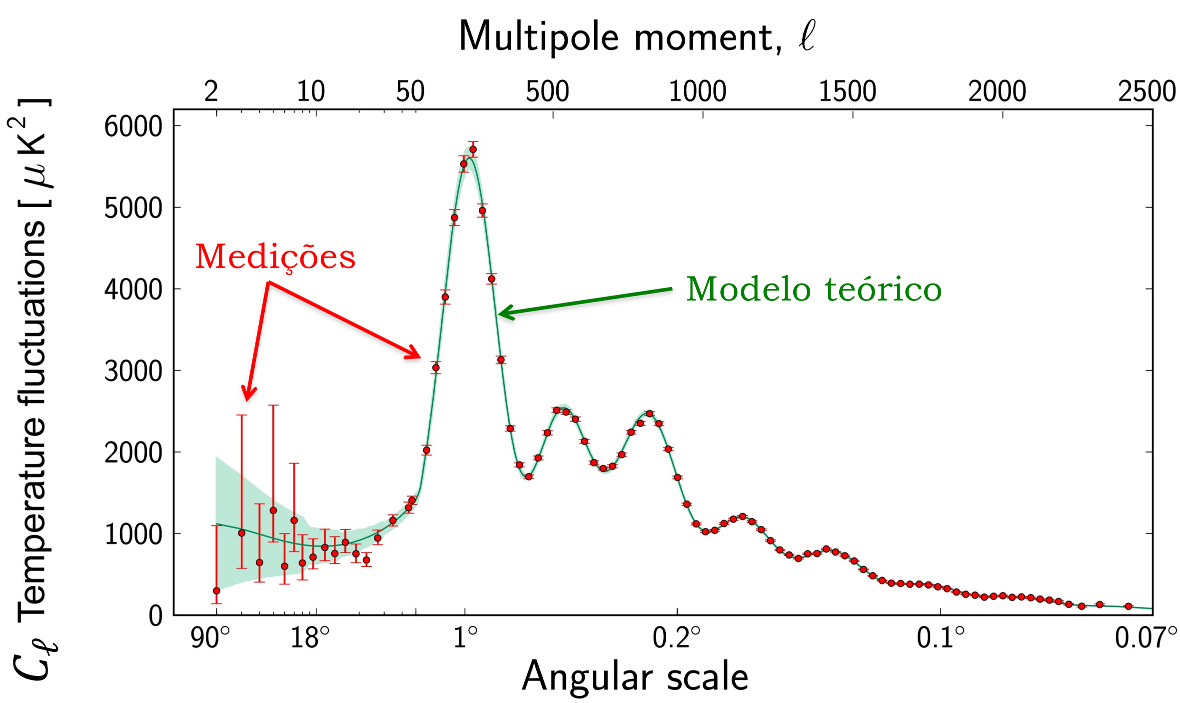


Earth's "elliptical" map (Mollweide projection)



CMB angular power spectrum

Planck



CMB: temperature fluctuations on the sphere

- Can be expanded as a sum of functions, the spherical harmonics Y_{lm}, that are a basis on the surface of a sphere:

$$\Theta(\hat{n}) = \Delta T/T(\theta,\phi) = \sum_{\ell=0}^{\infty}\sum_{m=-\ell}^{\ell} a_{\ell m} Y_{\ell m}(\theta,\phi)$$

- The coefficients a_{lm} are the projection of the temperature fluctuation function onto the basis function Y_{lm} (it measures the contribution of a given Y_{lm} function to the temperature fluctuation):

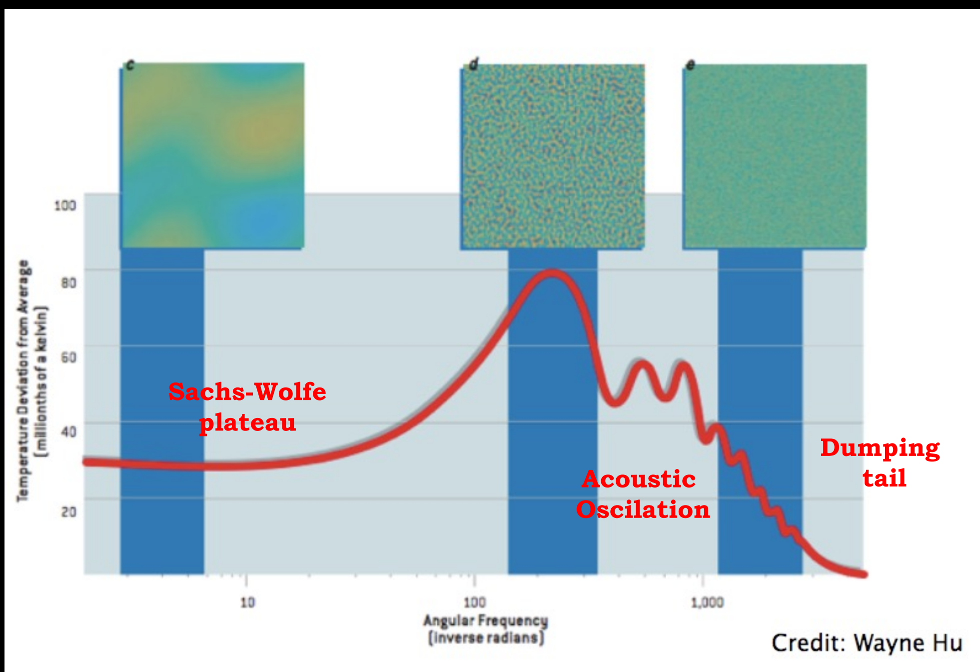
$$a_{\ell m} = \int Y^{*}_{\ell m}(\theta',\phi') \frac{\Delta T}{T}(\theta',\phi') d\Omega'$$

- The angular power spectrum is defined as an angular correlation function in the celestial sphere:

Cl

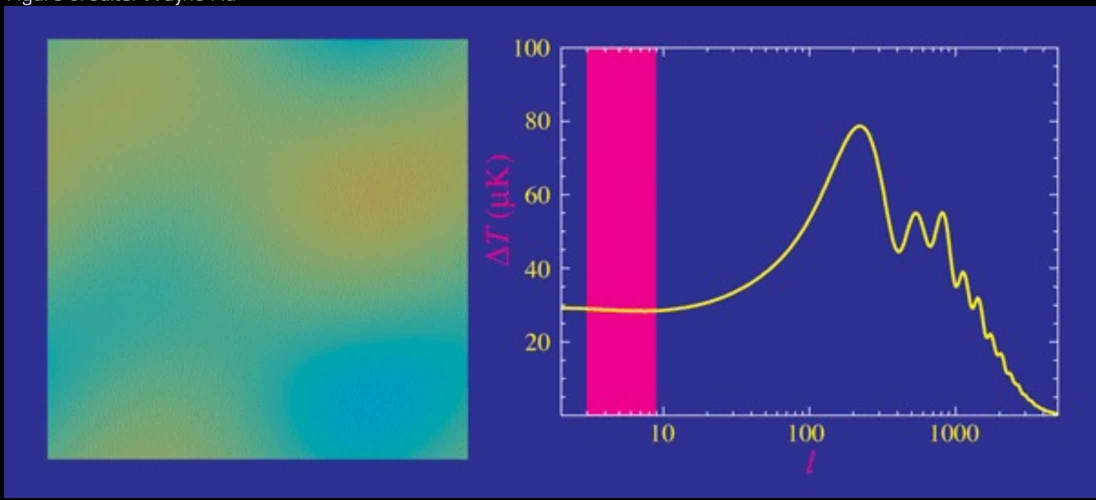
$$C(\hat{n},\hat{n}') \equiv \left\langle \frac{\Delta T}{T}(\hat{n}) \frac{\Delta T}{T}(\hat{n}') \right\rangle = \sum_{\ell\,\ell'}\sum_{m\,m'} \langle a^{*}_{\ell m} a_{\ell' m'} \rangle Y^{*}_{\ell m}(\hat{n}) Y_{\ell' m'}(\hat{n}')$$

CMB angular power spectrum

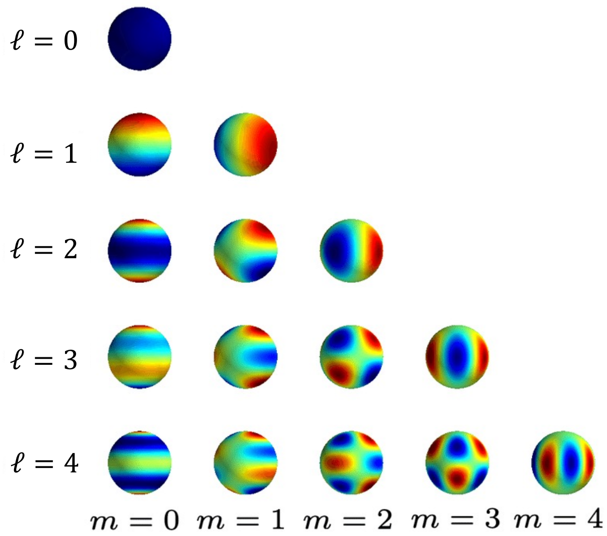


CMB angular power spectrum

Figure credits: Wayne Hu



Spherical harmonics

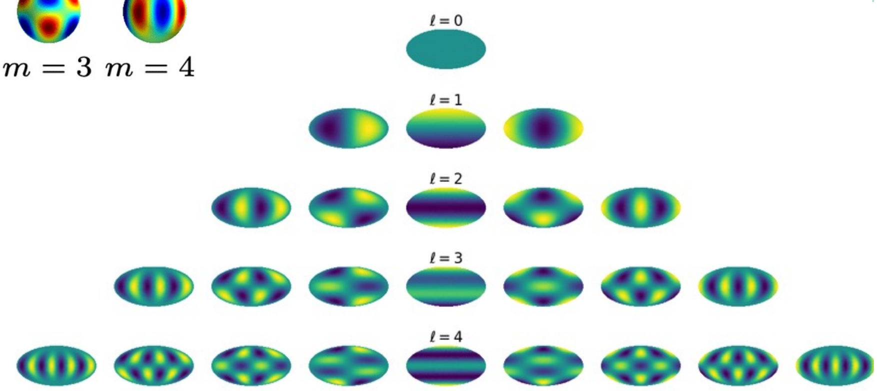


The $Y_{\ell m}$ are a basis in S^2 (celestial sphere)

$$Y_\ell^m(\theta,\varphi) = \sqrt{\frac{(2\ell+1)}{4\pi}\frac{(\ell-m)!}{(\ell+m)!}} \cdot e^{im\varphi} \cdot P_\ell^m(\cos\theta)$$

$$\int_{\theta=0}^{\pi}\int_{\varphi=0}^{2\pi} Y_\ell^m Y_{\ell'}^{m'*}\, d\Omega = \delta_{\ell\ell'}\delta_{mm'} \qquad d\Omega = \sin\theta\, d\varphi\, d\theta$$

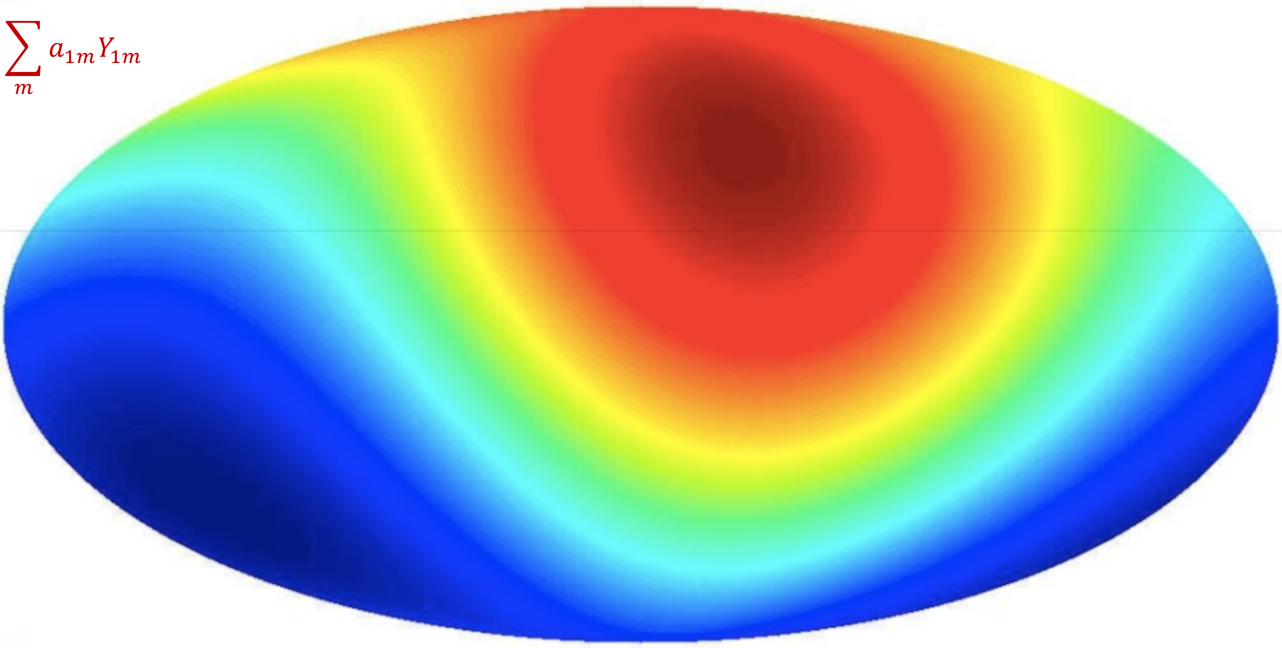
Right: Mollweide projection, of the real part of $Y_{\ell m}$ for $l = 1 \ldots 4$. For each $-l \le m \le l$. In this projection, the z–axis is oriented upwards, with colatitudinal angle θ = 0 at the north pole and θ = π at the south pole. The azimuthal angle runs from [π, −π] from left to right. The zonal spherical harmonics ($m = 0$), which have rotational symmetry about the z–axis, are plotted in the central column. The sectorial harmonics ($l = |m|$) are shown in the outermost panels of each row. The only difference between modes with ±m is a phase shift of 90^0.



Spherical harminic components of a well know map...

$\ell=1$

$\sum_m a_{1m} Y_{1m}$



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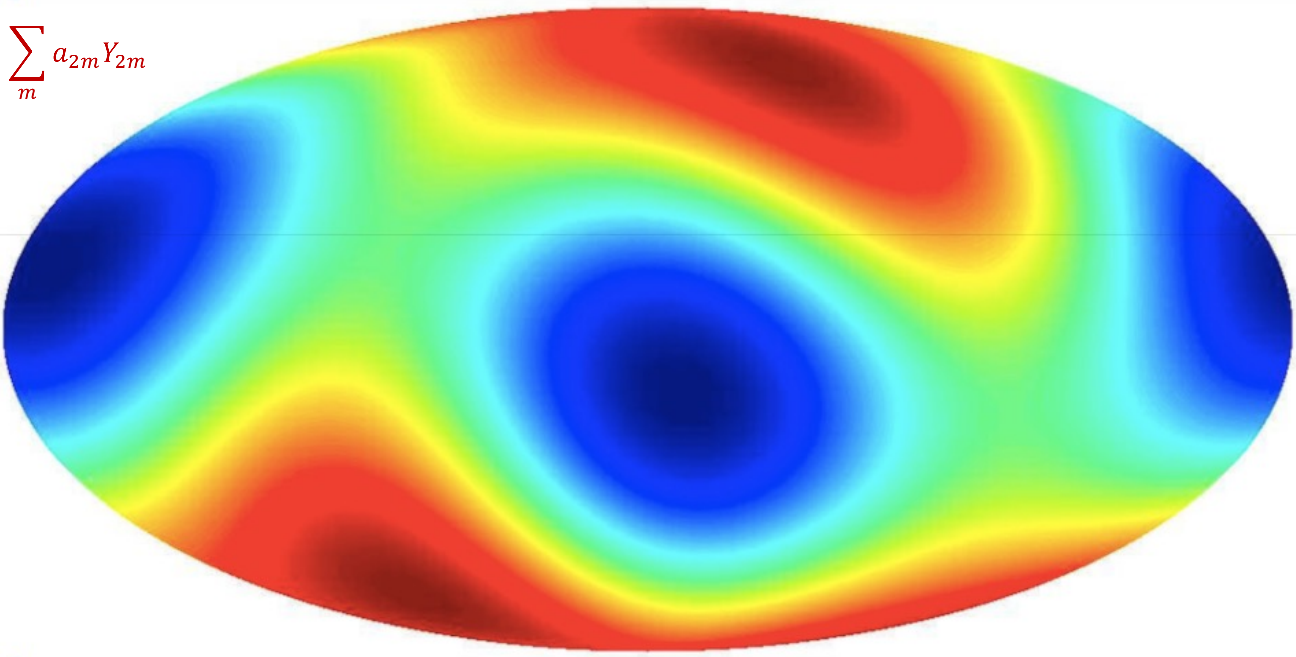
$\ell=2$

$\sum_m a_{2m} Y_{2m}$



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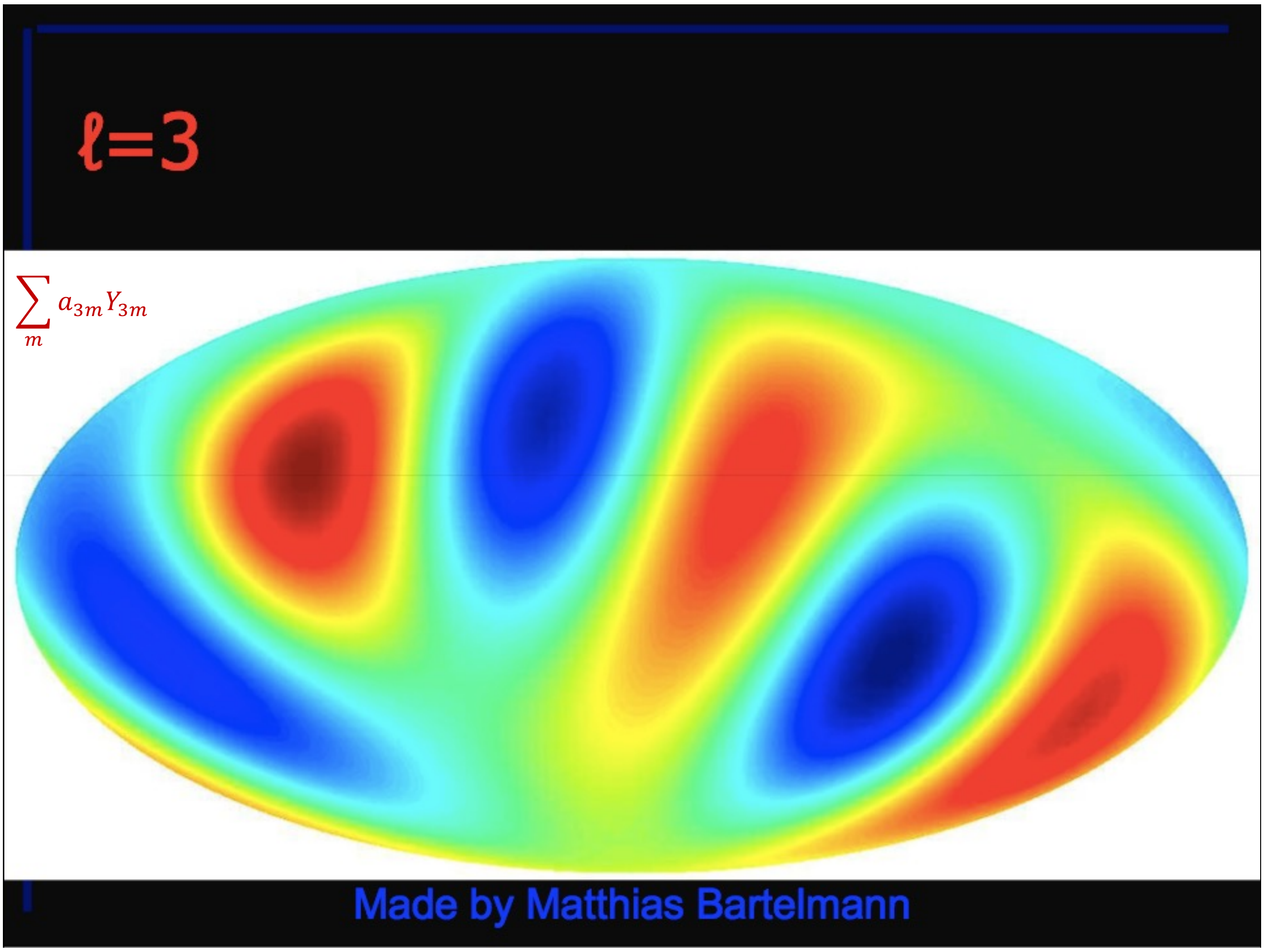
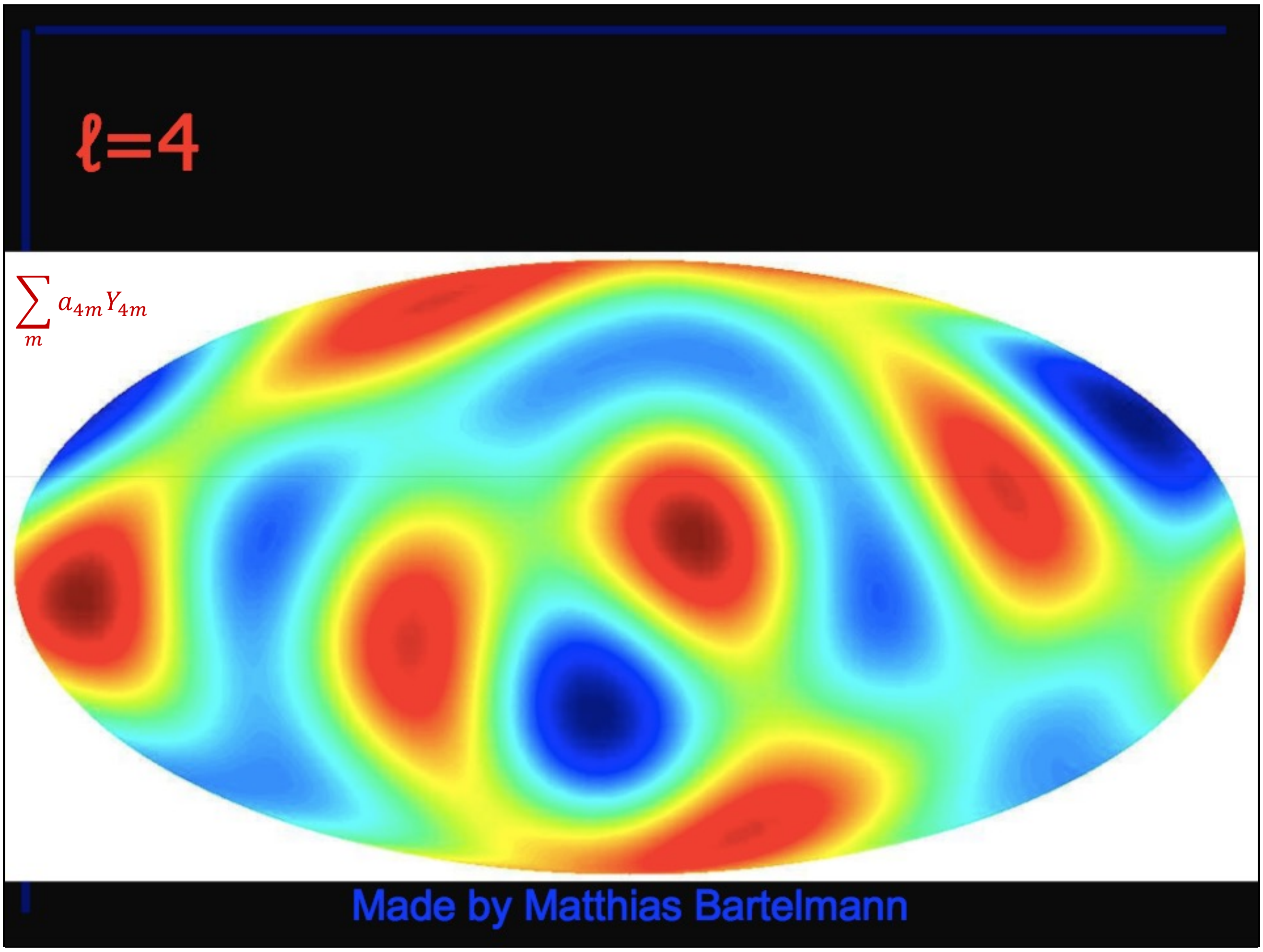
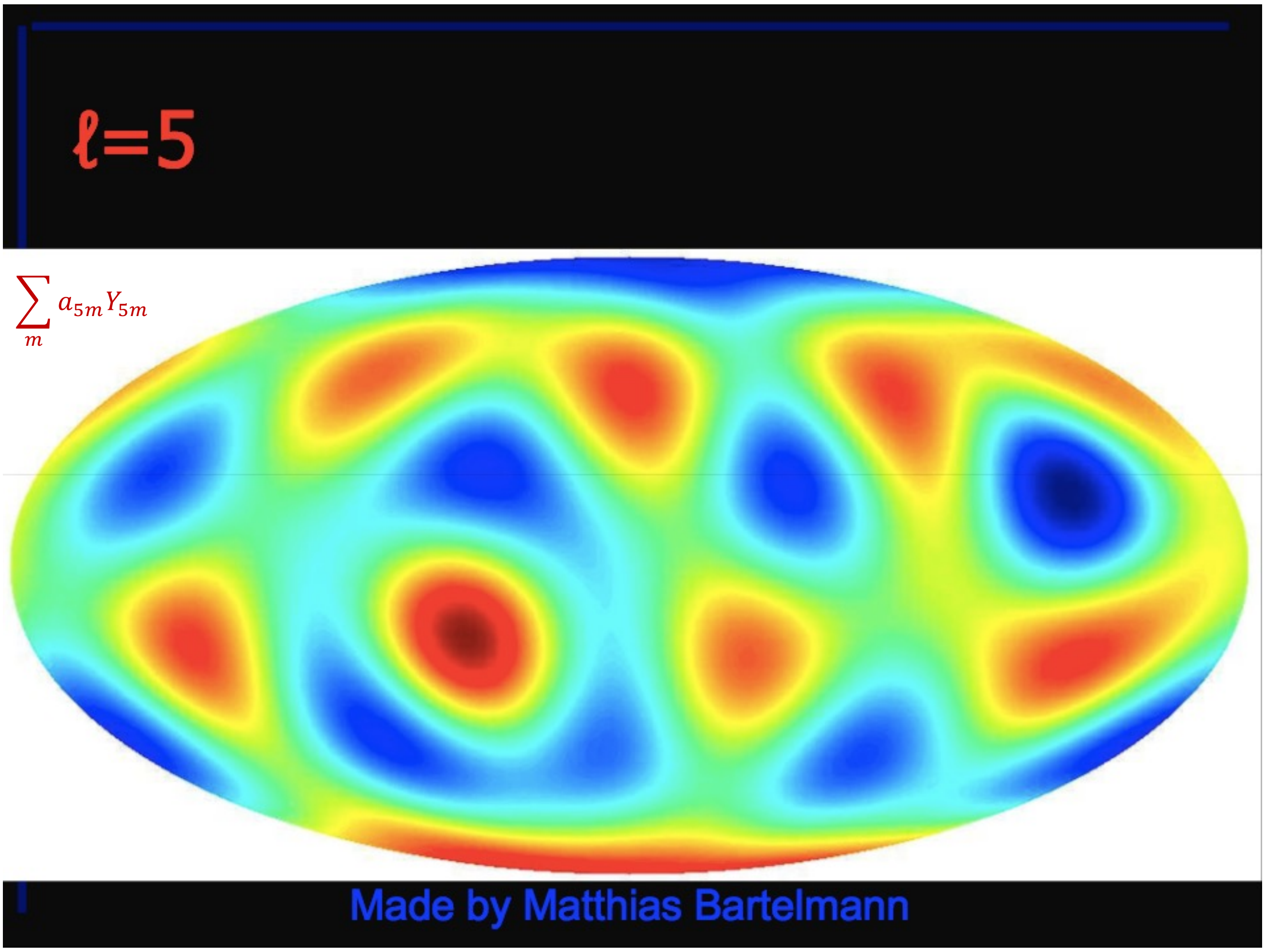
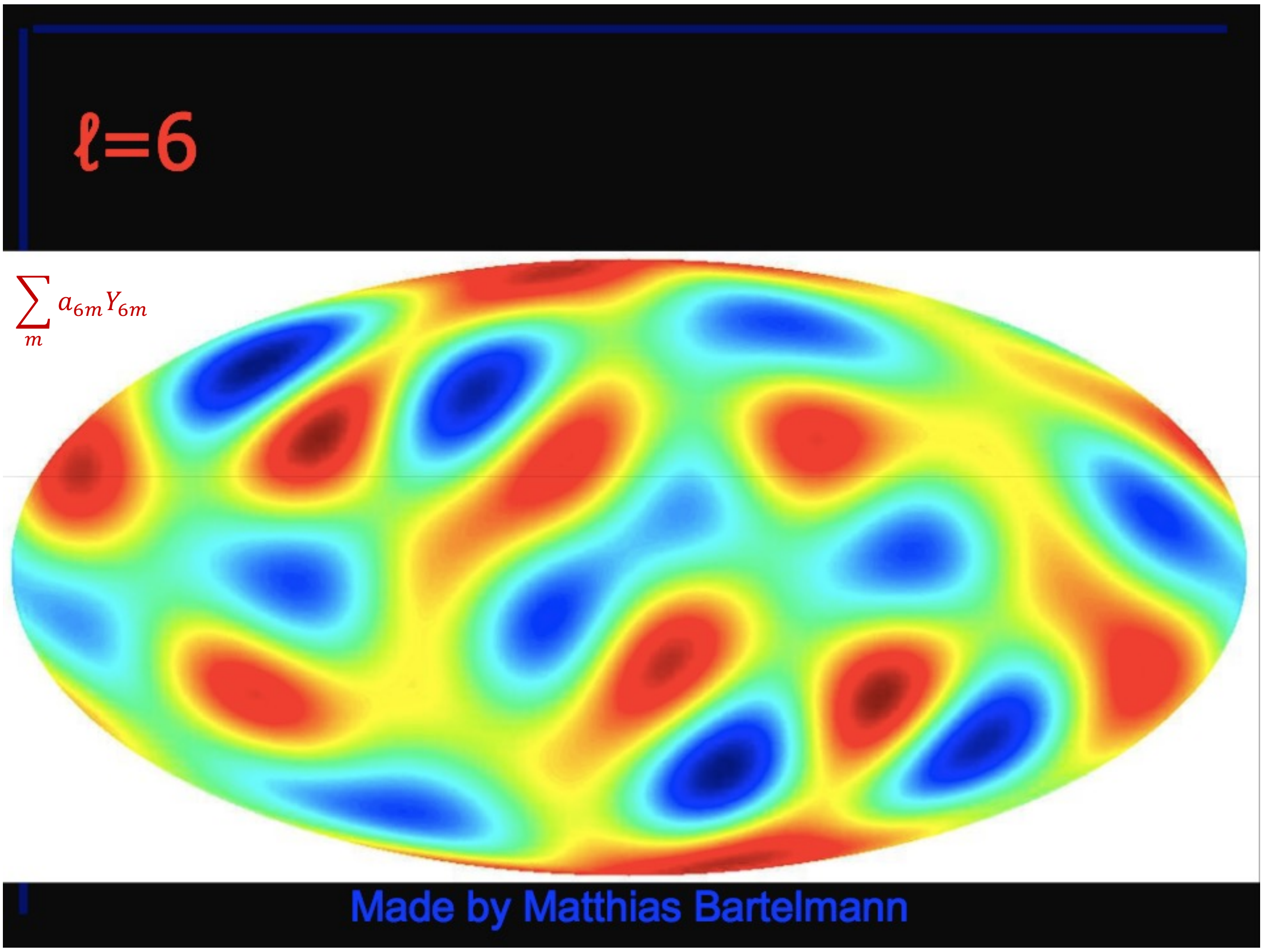
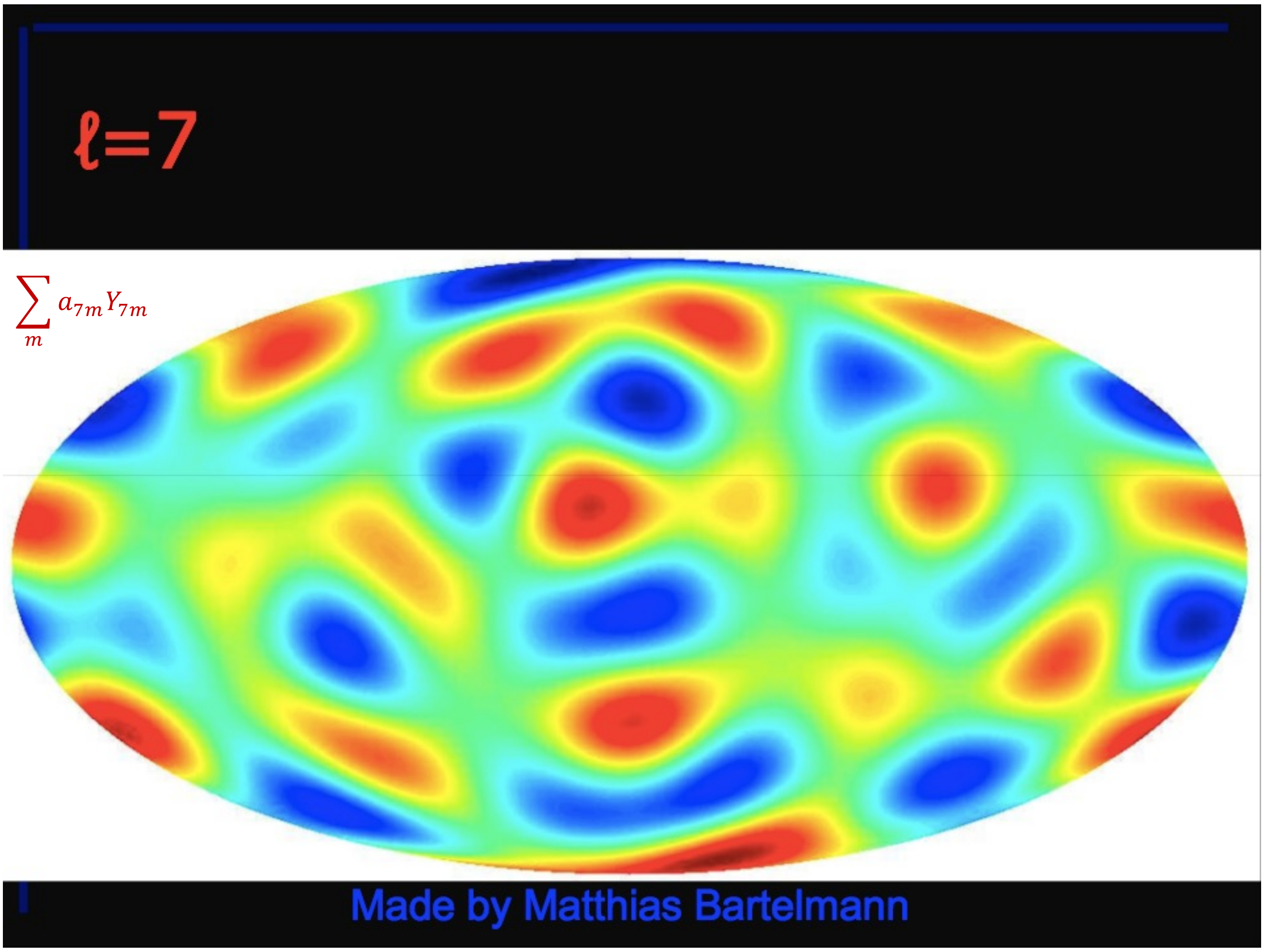
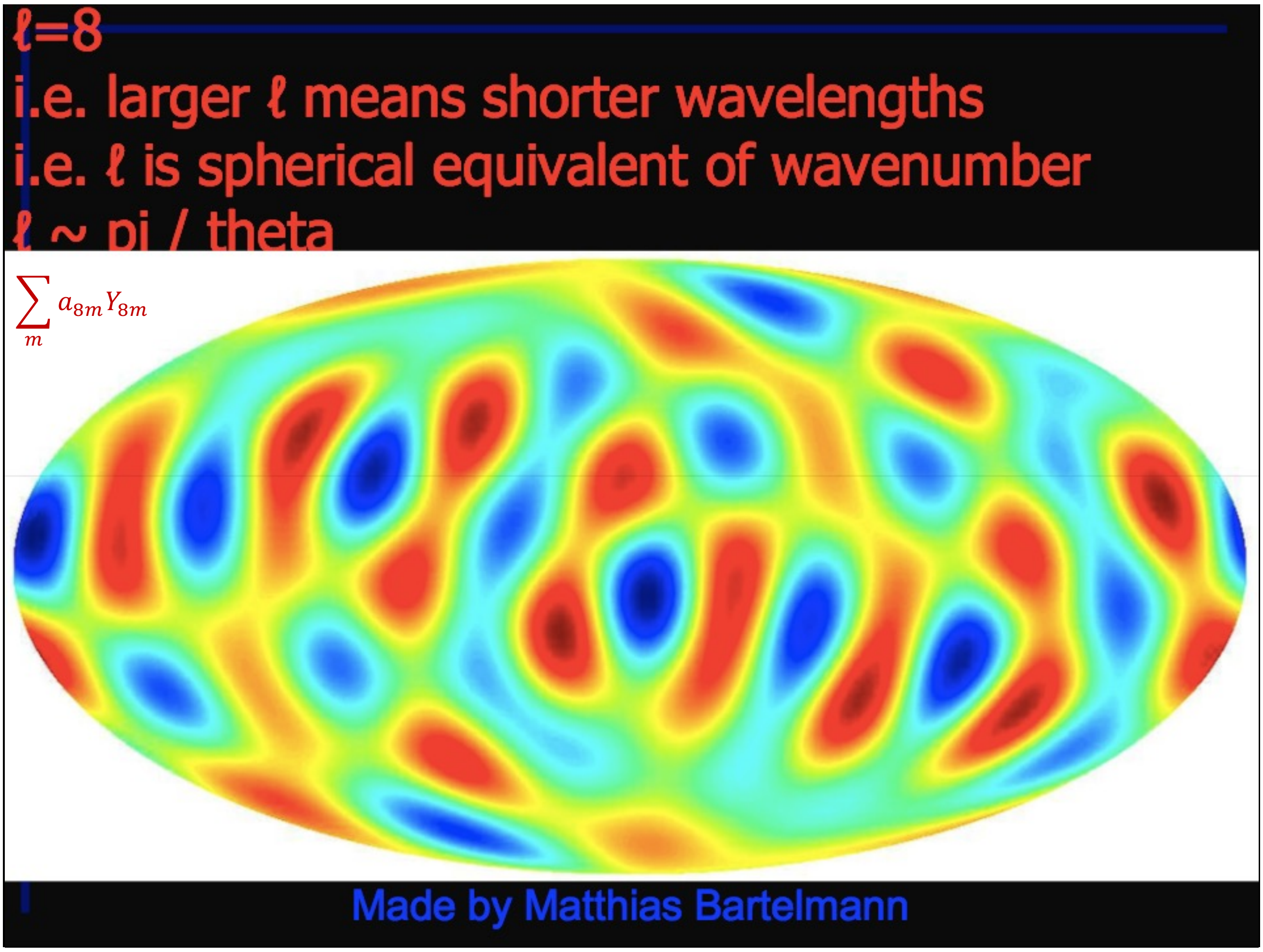

ℓ=3
$\sum_m a_{3m} Y_{3m}$
Made by Matthias Bartelmann

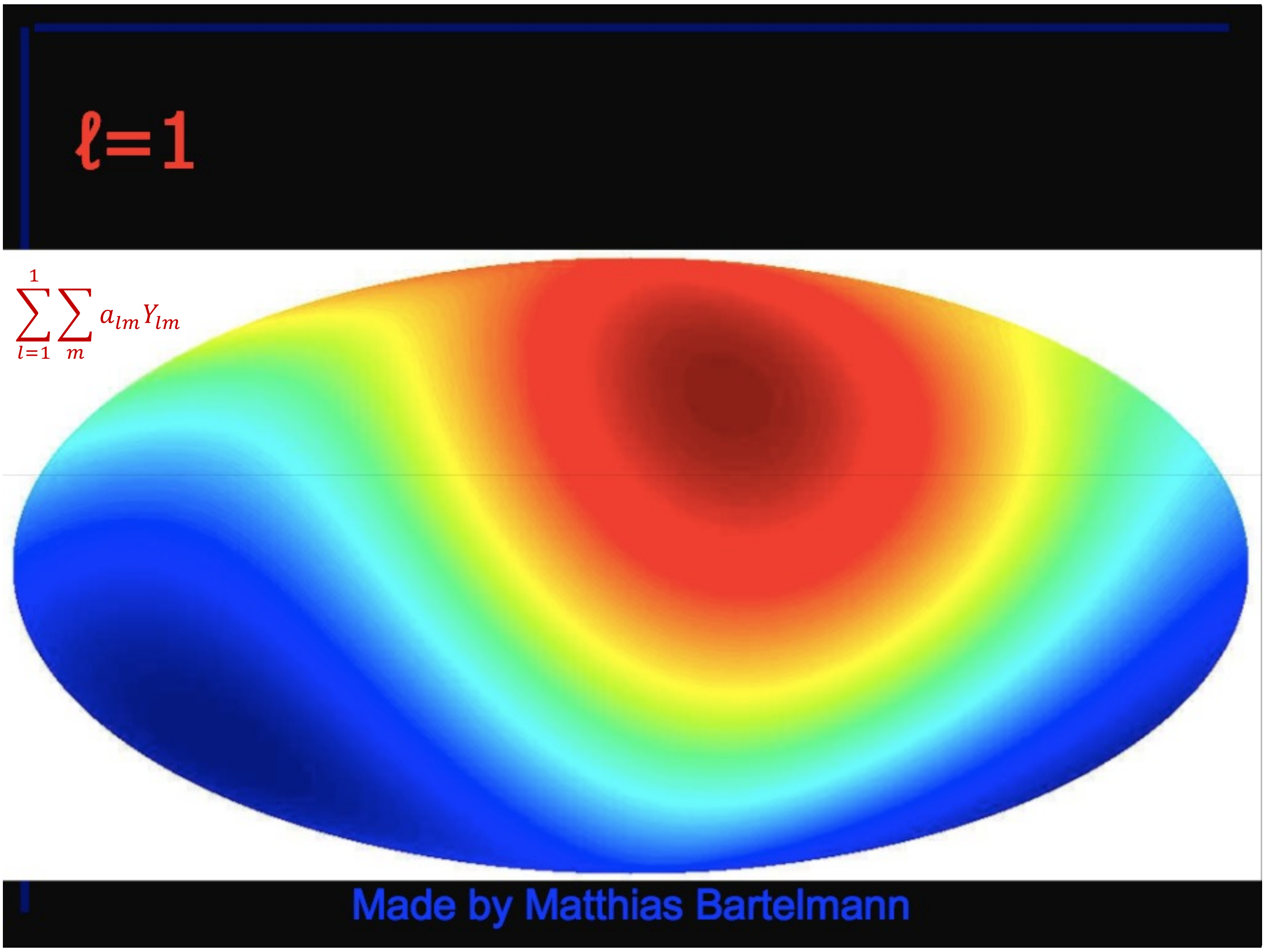

ℓ=4
$\sum_m a_{4m} Y_{4m}$
Made by Matthias Bartelmann

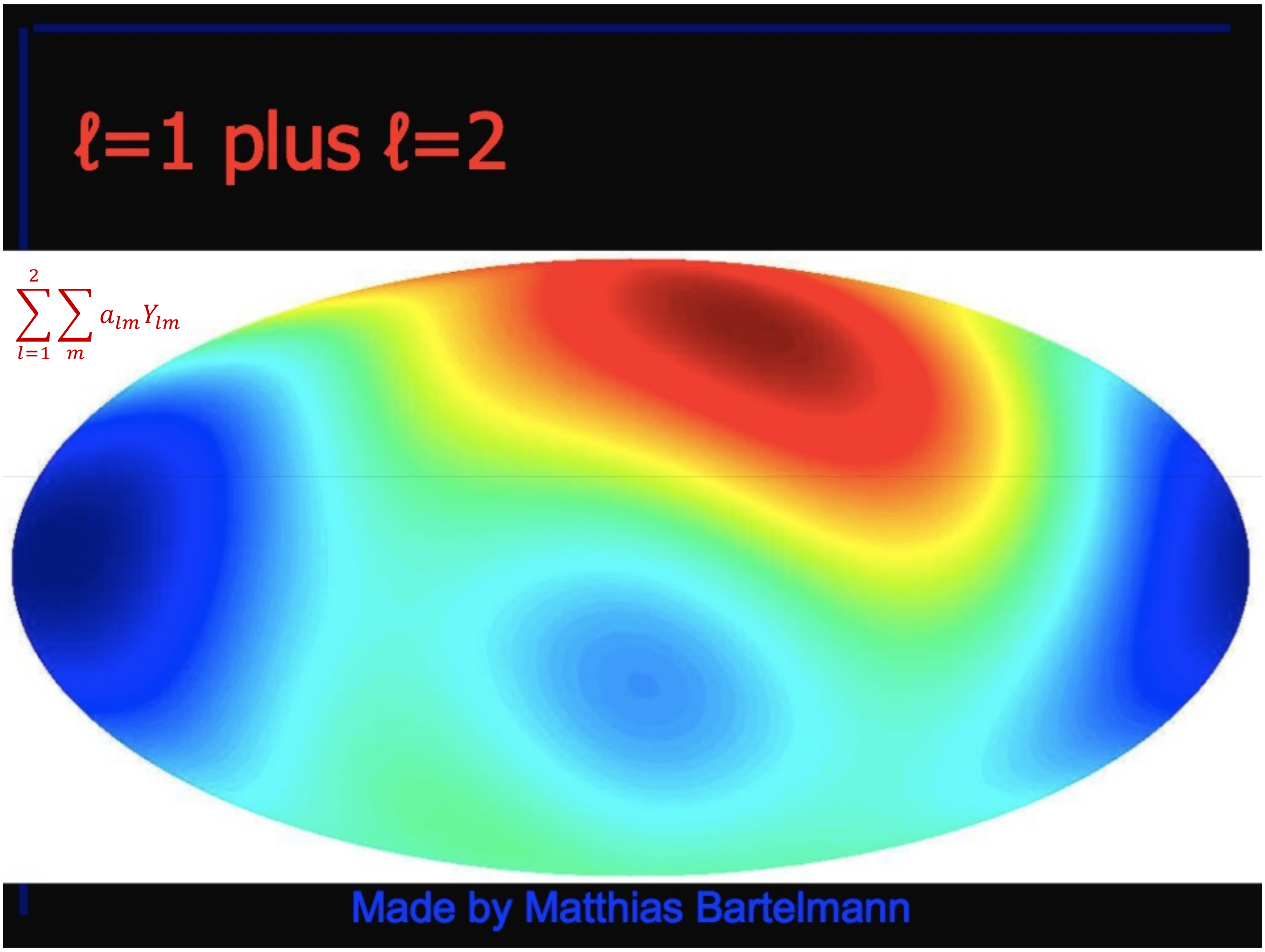

ℓ=5
$\sum_m a_{5m} Y_{5m}$
Made by Matthias Bartelmann


ℓ=6
$\sum_m a_{6m} Y_{6m}$
Made by Matthias Bartelmann


ℓ=7
$\sum_m a_{7m} Y_{7m}$
Made by Matthias Bartelmann

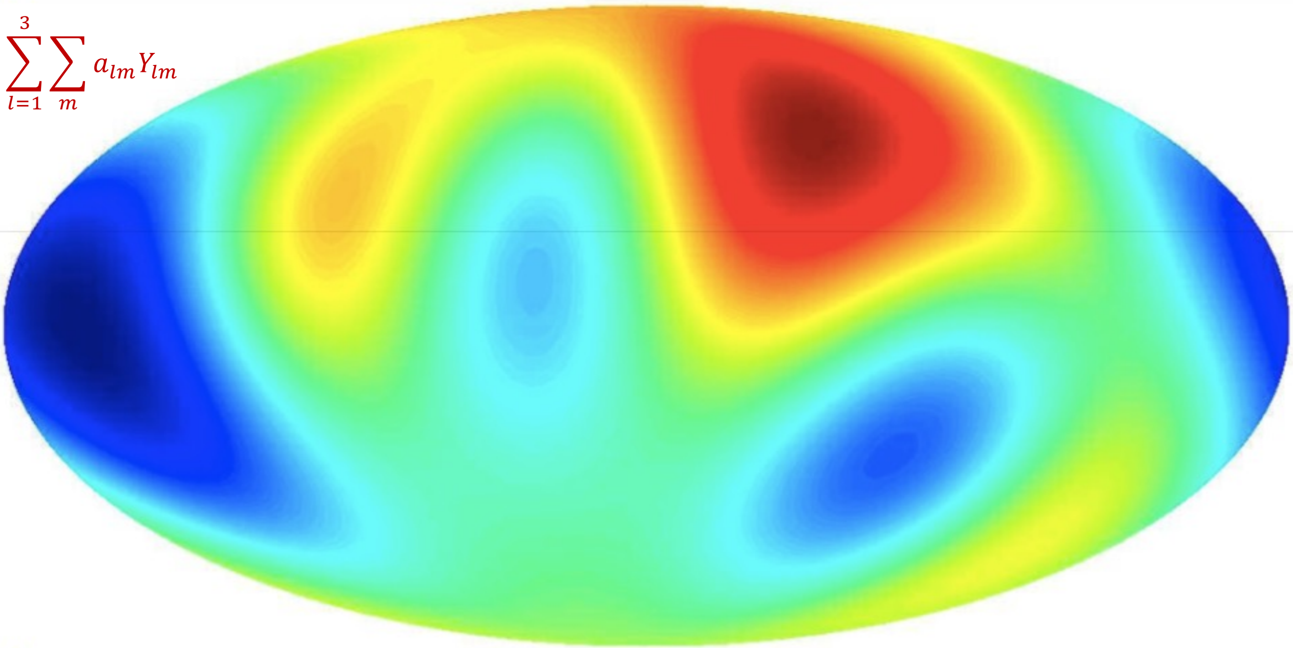

ℓ=8
i.e. larger ℓ means shorter wavelengths
i.e. ℓ is spherical equivalent of wavenumber
ℓ ~ pi / theta
$\sum_m a_{8m} Y_{8m}$
Made by Matthias Bartelmann


ℓ=1
$\sum_{l=1}^{1}\sum_{m} a_{lm}Y_{lm}$
Made by Matthias Bartelmann


ℓ=1 plus ℓ=2
$\sum_{l=1}^{2}\sum_{m} a_{lm}Y_{lm}$
Made by Matthias Bartelmann

ℓ=1 plus ℓ=2 plus ℓ=3

$$\sum_{l=1}^{3}\sum_{m} a_{lm}Y_{lm}$$



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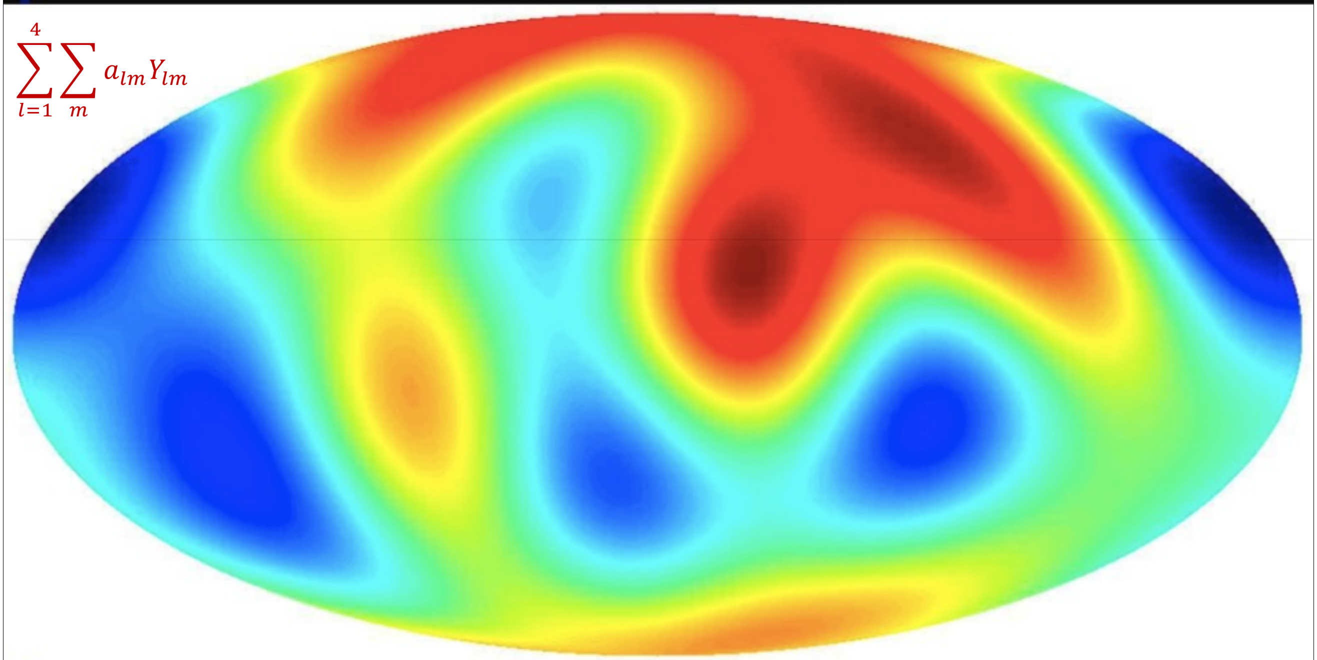
Sum ℓ=1 to 4

$$\sum_{l=1}^{4}\sum_{m} a_{lm}Y_{lm}$$



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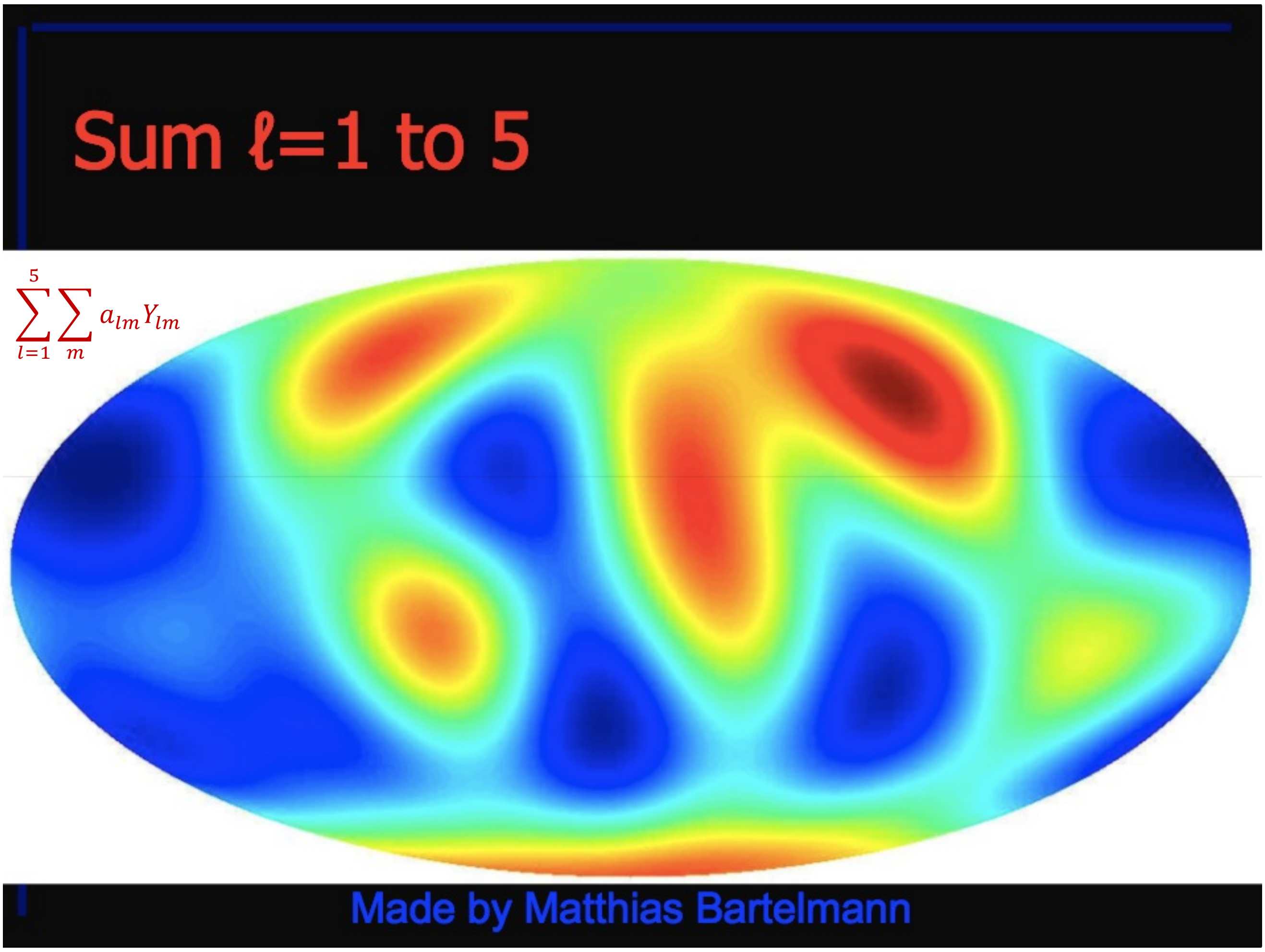
Sum ℓ=1 to 5
$\sum_{l=1}^{5}\sum_{m} a_{lm}Y_{lm}$
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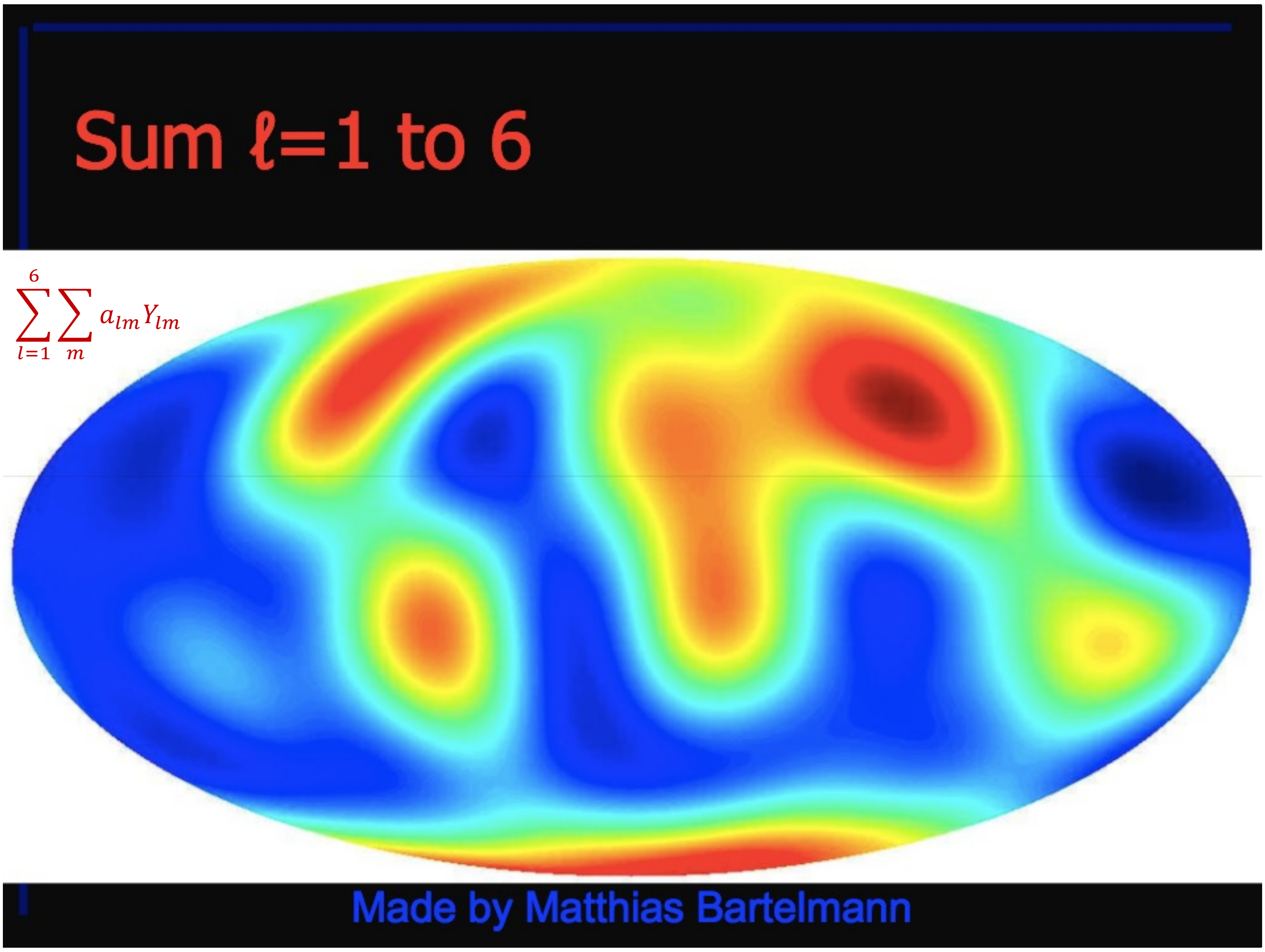
Sum ℓ=1 to 6
$\sum_{l=1}^{6}\sum_{m} a_{lm}Y_{lm}$
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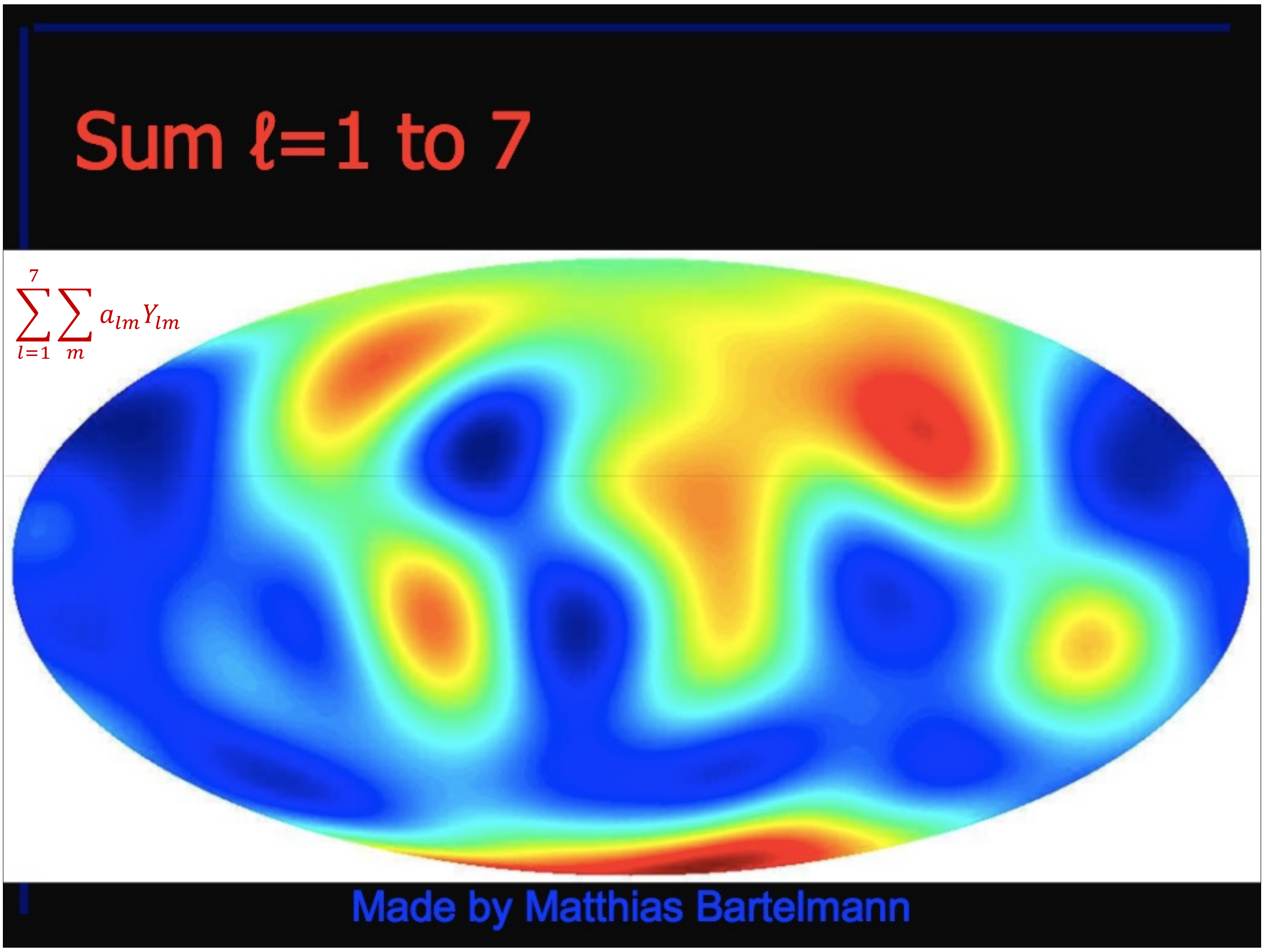
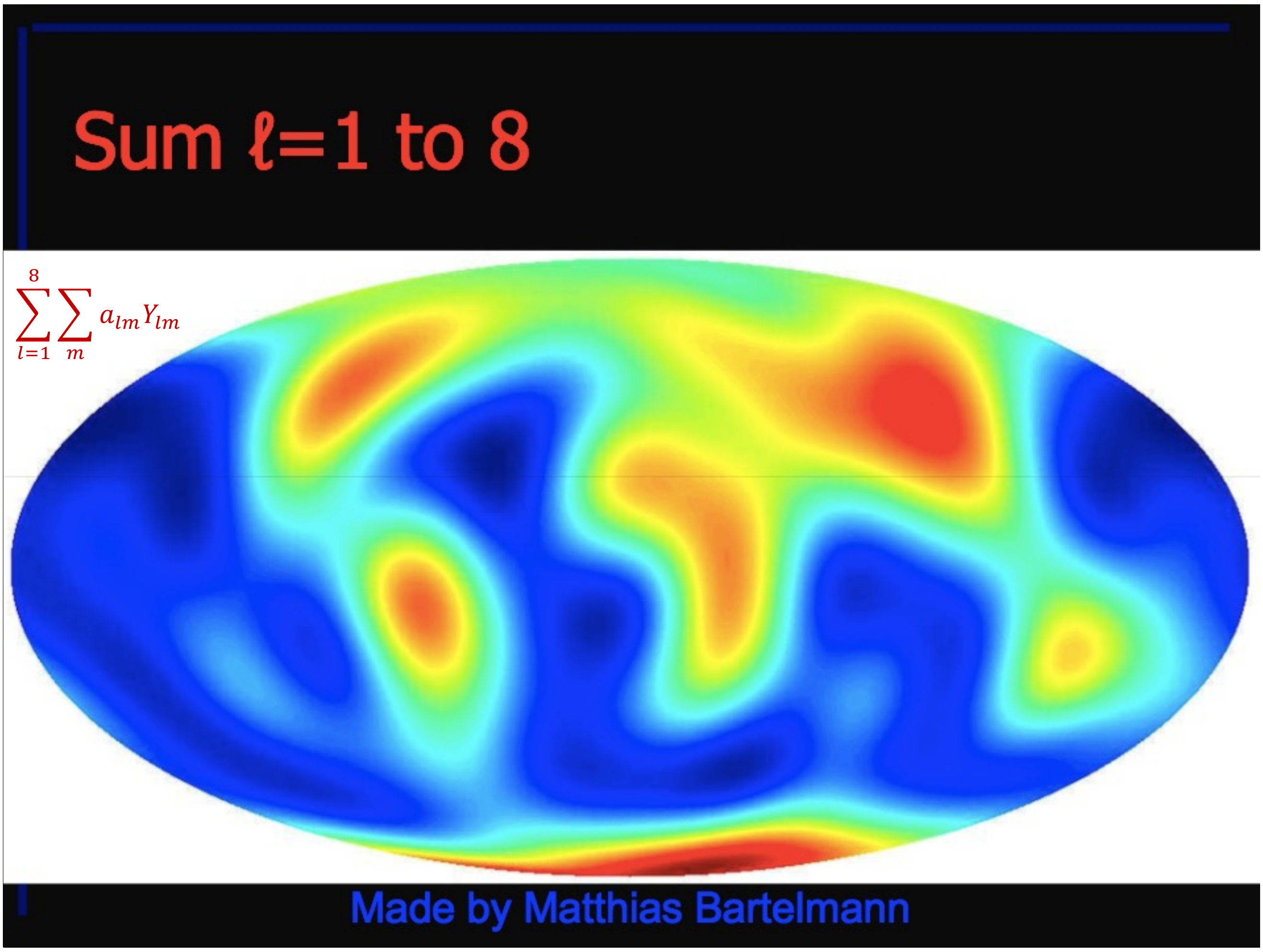
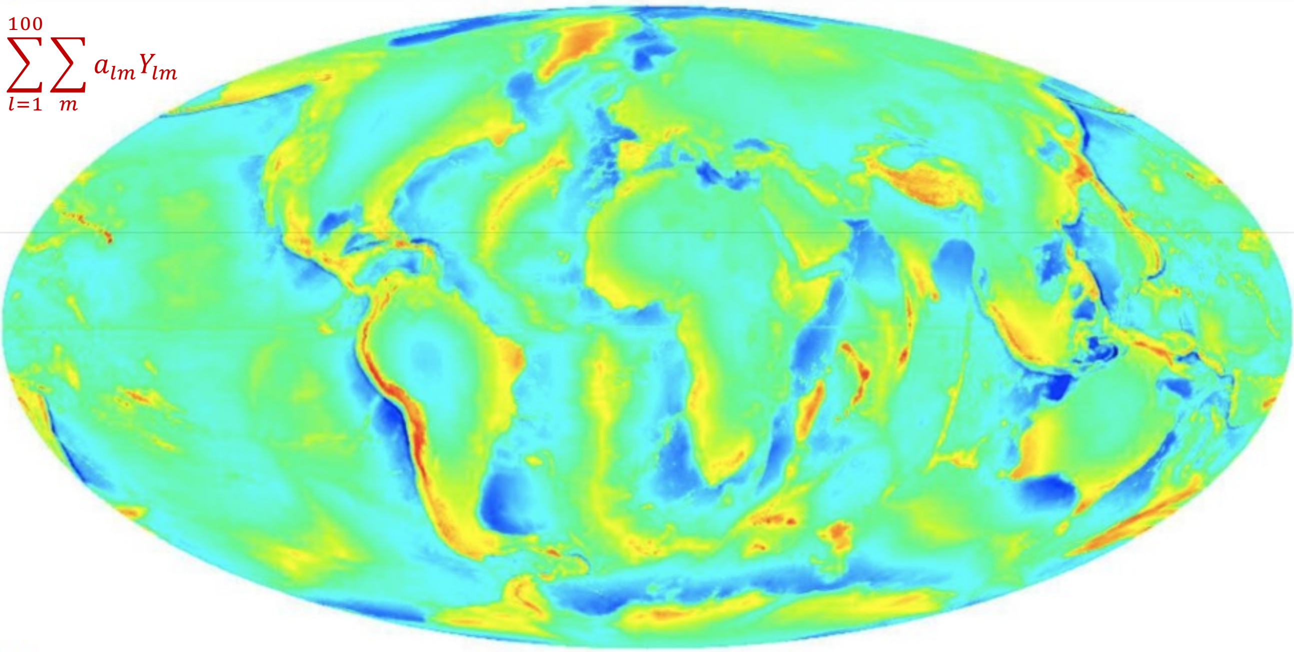

Sum ℓ=1 to 7
$\sum_{l=1}^{7}\sum_{m} a_{lm}Y_{lm}$
Made by Matthias Bartelmann


Sum ℓ=1 to 8
$\sum_{l=1}^{8}\sum_{m} a_{lm}Y_{lm}$
Made by Matthias Bartelmann

Sum up to some high ℓ

$$\sum_{l=1}^{100}\sum_{m} a_{lm} Y_{lm}$$



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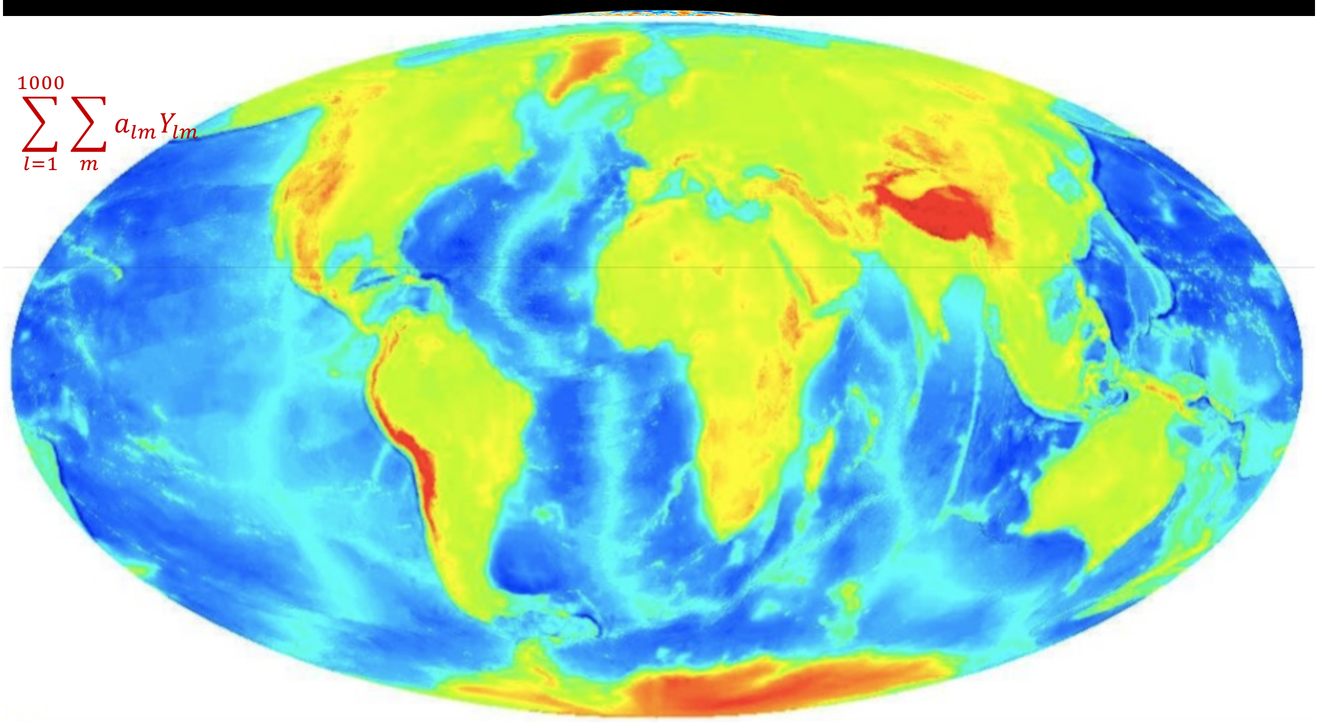
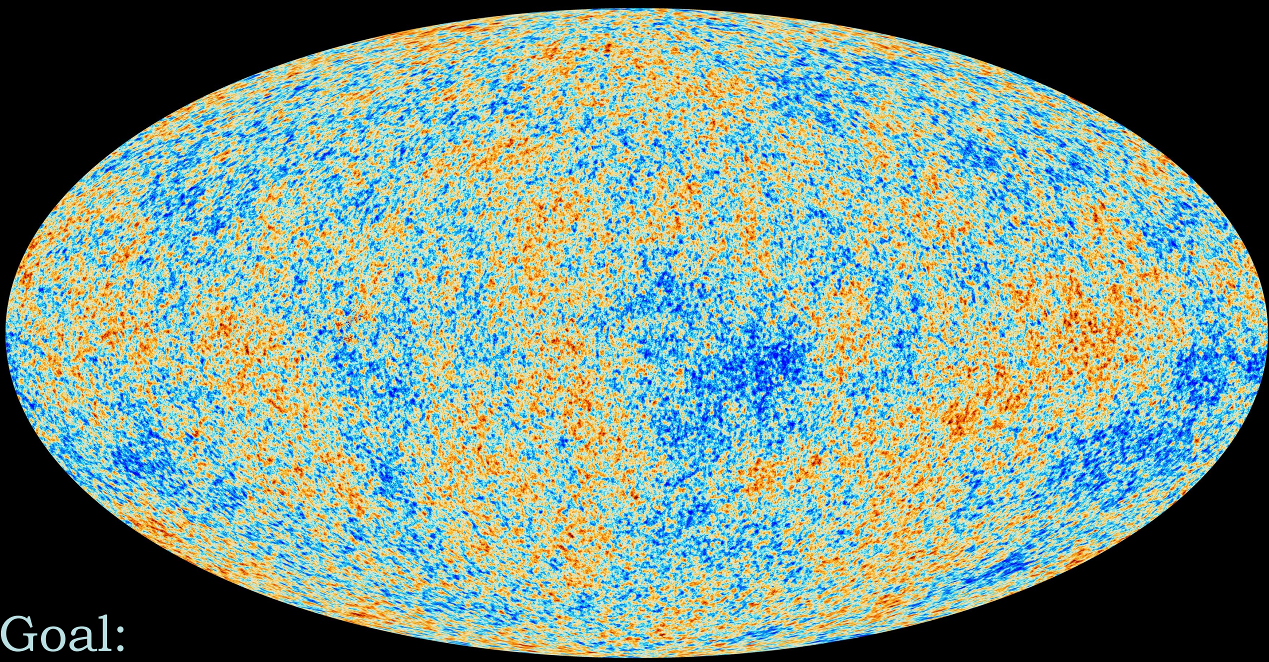
Earth's map with all contributions up to Planck's CMB map resolution

$$\sum_{l=1}^{1000}\sum_{m} a_{lm} Y_{lm}$$

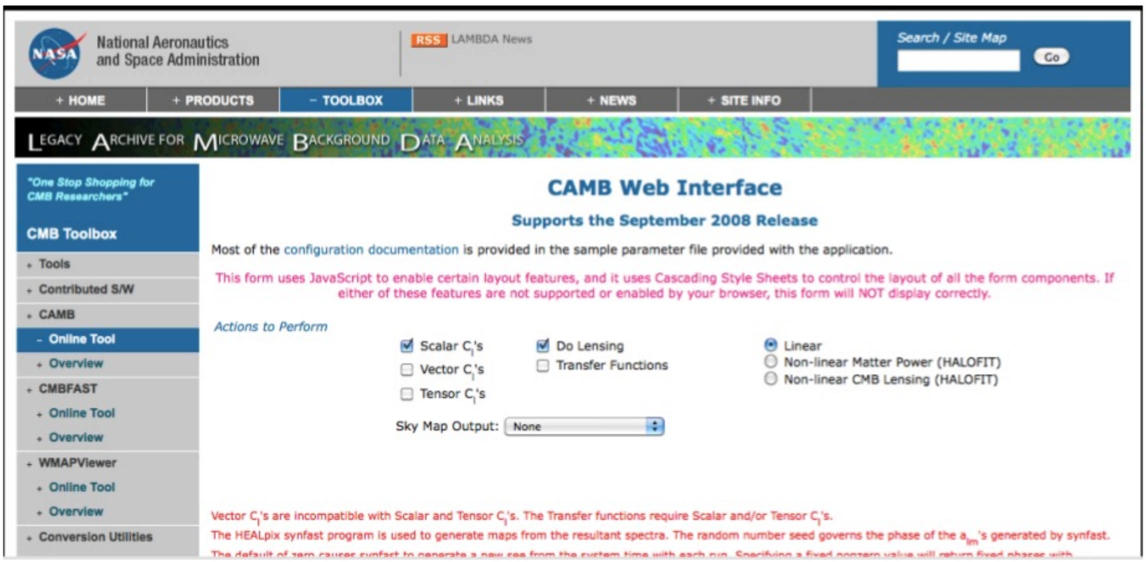


CMB temperature fluctuations map



Goal:
Use CMB map information to constrain theoretical Cosmological models

Online C_l calculators

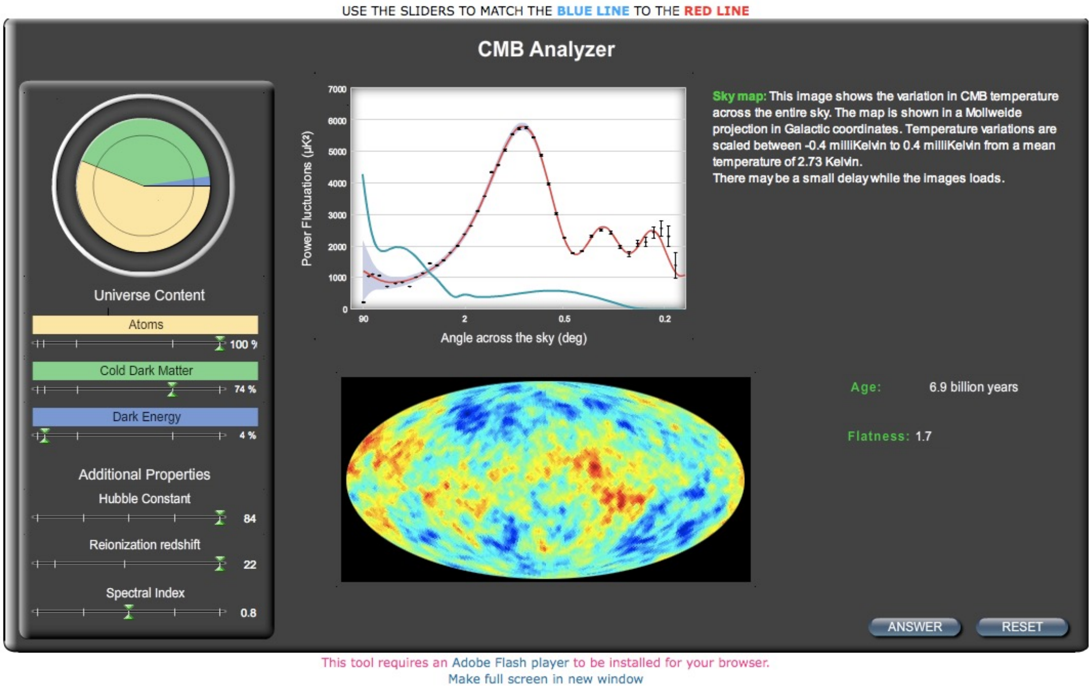


CMB Toolbox: http://lambda.gsfc.nasa.gov/toolbox/

CAMB website: http://camb.info/
CMBFast website: http://www.cmbfast.org/

CMB analyzer

http://lambda.gsfc.nasa.gov/education/cmb_plotter/



CMB parameter cheat sheet